Nucleosynthesis in heavy-ion collisions at the LHC via the Saha equation

Volodymyr Vovchenko

Goethe University Frankfurt & Frankfurt Institute for Advanced Studies

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in collaboration with K. Gallmeister, J. Schaffner-Bielich, C. Greiner

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Heavy-ion collisions



Heavy-ion collision experiments study properties of strongly interacting matter at extreme temperatures and densities, recreate conditions present in Early Universe

Loosely-bound objects in heavy-ion collisions



binding energies: ²H, ³He, ⁴He, ³_AH: 2.22, 7.72, 28.3, 0.130 MeV $\ll T \sim 150$ MeV "snowballs in hell"

The production mechanism is not established. Common approaches include **thermal** nuclei emission together with hadrons [Andronic et al., PLB '11;...] or final-state **coalescence** of nucleons close in phase-space [Butler, Pearson, PRL '61; Scheibl, Heinz, PRC '99;...]

Two experimental observations at the LHC

1. Measured yields are described by thermal model at $T_{ch} \approx 155 \text{ MeV}^*$



2. Spectra described by blast-wave model at $T_{kin} \approx 100 - 120 \text{ MeV}^*$



[ALICE collaboration, PRC 93, 024917 (2016)]

What happens between T_{ch} and T_{kin} ?

Big Bang nucleosynthesis





• Early stage of Big Bang nucleosynthesis described by Nuclear Statistical Equilibrium (Saha equation)

$$X_{A} = d_{A} \left[\zeta(3)^{A-1} \pi^{\frac{1-A}{2}} 2^{\frac{3A-5}{2}} \right] A^{\frac{5}{2}} \left(\frac{T}{m_{N}} \right)^{\frac{3}{2}(A-1)} \eta^{A-1} X_{p}^{Z} X_{n}^{A-Z} \exp \left(\frac{1}{2} \left(\frac{T}{m_{N}} \right)^{\frac{3}{2}(A-1)} \right)^{\frac{3}{2}(A-1)} \eta^{A-1} X_{p}^{Z} X_{n}^{A-Z} \exp \left(\frac{1}{2} \left(\frac{T}{m_{N}} \right)^{\frac{3}{2}(A-1)} \right)^{\frac{3}{2}(A-1)} \eta^{A-1} X_{p}^{Z} X_{n}^{A-Z} \exp \left(\frac{1}{2} \left(\frac{T}{m_{N}} \right)^{\frac{3}{2}(A-1)} \right)^{\frac{3}{2}(A-1)} \eta^{A-1} X_{p}^{Z} X_{n}^{A-Z} \exp \left(\frac{1}{2} \left(\frac{T}{m_{N}} \right)^{\frac{3}{2}(A-1)} \right)^{\frac{3}{2}(A-1)} \eta^{A-1} X_{p}^{Z} X_{n}^{A-Z} \exp \left(\frac{1}{2} \left(\frac{T}{m_{N}} \right)^{\frac{3}{2}(A-1)} \right)^{\frac{3}{2}(A-1)} \eta^{A-1} X_{p}^{Z} X_{n}^{A-Z} \exp \left(\frac{1}{2} \left(\frac{T}{m_{N}} \right)^{\frac{3}{2}(A-1)} \right)^{\frac{3}{2}(A-1)} \eta^{A-1} X_{p}^{Z} X_{n}^{A-Z} \exp \left(\frac{1}{2} \left(\frac{T}{m_{N}} \right)^{\frac{3}{2}(A-1)} \right)^{\frac{3}{2}(A-1)} \eta^{A-1} X_{p}^{Z} X_{n}^{A-Z} \exp \left(\frac{1}{2} \left(\frac{T}{m_{N}} \right)^{\frac{3}{2}(A-1)} \right)^{\frac{3}{2}(A-1)} \eta^{A-1} X_{p}^{Z} X_{n}^{A-Z} \exp \left(\frac{1}{2} \left(\frac{T}{m_{N}} \right)^{\frac{3}{2}(A-1)} \right)^{\frac{3}{2}(A-1)} \eta^{A-1} X_{p}^{Z} X_{n}^{A-Z} \exp \left(\frac{1}{2} \left(\frac{T}{m_{N}} \right)^{\frac{3}{2}(A-1)} \right)^{\frac{3}{2}(A-1)} \eta^{A-1} X_{p}^{Z} X_{n}^{A-Z} \exp \left(\frac{1}{2} \left(\frac{T}{m_{N}} \right)^{\frac{3}{2}(A-1)} \right)^{\frac{3}{2}(A-1)} \eta^{A-1} X_{p}^{Z} X_{n}^{A-Z} \exp \left(\frac{1}{2} \left(\frac{T}{m_{N}} \right)^{\frac{3}{2}(A-1)} \eta^{A-1} X_{p}^{Z} X_{n}^{A-Z} \exp \left(\frac{T}{m_{N}} \right)^{\frac{3}{2}(A-1)} \eta^{A-1} X_{p}^{Z} X_{n}^{A-1} X_{p}^{Z} X_{n}^{A-1} X_{p}^{2} X_{n}^{A-1} X_{p}^{2} X_{n}^{A-1} X_{p}^{2} X_{n}^{A-1} X_{p}^{2} X_{n}^{A-1} X_{p}^{2} X_{n}^{A-1} X_{p}^{A-1} X_{p}^$$

 $\eta \sim 10^{-10}$ – baryon-to-photon ratio

[E. Kolb, M. Turner, "The Early Universe" (1990)] 5

¹H

Ήe

(d,n)

(p, y)

 $\frac{B_A}{T}$

(d,p)

He,2p)

He

(d,n)

Big Bang vs "Little Bangs"



- Hadrons (nucleons) form and "freeze-out" chemically before nuclei
- Bosons (photons or pions) catalyse nucleosynthesis

e.g. $p + n \leftrightarrow d + \gamma$ vs $p + n + \pi \leftrightarrow d + \pi$

Similarities:

- Inelastic nucleonic reactions freeze-out before nuclei formation
- Isentropic expansion of boson-dominated matter (photons in BBN vs mesons in HIC), baryon-to-boson ratio: $\eta_{BBN} \sim 10^{-10}$, $\eta_{LHC} \sim 0.05$
- Strong nuclear formation and regeneration reactions \rightarrow Saha equation

Differences:

- Time scales: 1-100 s in BBN vs $\sim 10^{-22}$ s in HIC
- Temperatures: $T_{BBN} < 1$ MeV vs $T_{HIC} \sim 100$ MeV
- Binding energies, proton-neutron mass difference, and neutron lifetime important in BBN, less so in HICs
- $\mu_B \approx 0$ at the LHC, $\mu_B \neq 0$ in BBN
- Resonance feeddown important at LHC, irrelevant in BBN

LHC nucleosynthesis: simplified setup

- Chemical equilibrium lost at $T_{ch} = 155$ MeV, abundances of nucleons are frozen and acquire effective fugacity factors: $n_i = n_i^{eq} e^{\mu_N/T}$
- Isentropic expansion driven by effectively massless mesonic d.o.f.

$$rac{V}{V_{\mathsf{ch}}} = \left(rac{T_{\mathsf{ch}}}{T}
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, $\mu_N \simeq rac{3}{2} \ T \ \mathsf{ln} \left(rac{T}{T_{\mathsf{ch}}}
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• Detailed balance for nuclear reactions, $X + A \leftrightarrow X + \sum_i A_i$, X is e.g. a pion

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Saha equation
$$X_{A} = d_{A} \left[(d_{M})^{A-1} \zeta(3)^{A-1} \pi^{\frac{1-A}{2}} 2^{-\frac{3+A}{2}} \right] A^{5/2} \left(\frac{T}{m_{N}} \right)^{\frac{3}{2}(A-1)} \eta_{B}^{A-1} \exp\left(\frac{B_{A}}{T} \right)$$
$$d_{M} \sim 11 - 13, \quad \eta_{B} \simeq 0.03$$

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BBN:
$$X_A = d_A \left[\zeta(3)^{A-1} \pi^{\frac{1-A}{2}} 2^{\frac{3A-5}{2}} \right] A^{\frac{5}{2}} \left(\frac{T}{m_N} \right)^{\frac{1}{2}(A-1)} \eta^{A-1} X_p^Z X_n^{A-Z} \exp\left(\frac{B_A}{T} \right)$$

[E. Kolb, M. Turner, "The Early Universe" (1990)] 8

(Simplified) Saha equation vs thermal model

Saha equation:

Thermal model:

Strong exponential dependence on the temperature is eliminated in the Saha equation approach

Further, quantitative applications require numerical treatment of full spectrum of *massive* mesonic and baryonic resonances

Full numerical implementation

Expansion of hadron resonance gas in partial chemical equilibrium at $T < T_{ch}$ [H. Bebie, P. Gerber, J.L. Goity, H. Leutwyler, Nucl. Phys. B '92]

Chemical composition of stable hadrons is fixed, kinetic equilibrium maintained through quasi-elastic resonance reactions $\pi\pi \leftrightarrow \rho$, $\pi K \leftrightarrow K^*, \pi N \leftrightarrow \Delta$, etc.

Effective chemical potentials:

 $\tilde{\mu}_j = \sum_{i \in \text{stable}} \langle n_i \rangle_j \mu_i, \quad \langle n_i \rangle_j - \text{mean number of hadron } i \text{ from decays of hadron } j, \quad j \in \text{HRG}$

Conservation laws:



Numerical implementation within (extended) **Thermal-FIST** package [V.V., H. Stoecker, arXiv:1901.05249, *Computer Physics Communications*, in print] **open source:** https://github.com/vlvovch/Thermal-FIST



Full calculation: parameters



"Initial conditions" from thermal fits with Thermal-FIST to 0-10% ALICE hadron yields $T_{ch} = 155 \text{ MeV}, V_{ch} = 4700 \text{ fm}^3$ $[V.V., Gorenstein, Stoecker, 1807.02079]_{11}$

Full calculation: deuteron yield



Resonance feed-down is important in precision studies

LHC deuteron-synthesis

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Editors' Suggestion

Featured in Physics

Microscopic study of deuteron production in PbPb collisions at $\sqrt{s} = 2.76$ TeV via hydrodynamics and a hadronic afterburner

Dmytro Oliinychenko,¹ Long-Gang Pang,^{1,2} Hannah Elfner,^{3,4,5} and Volker Koch¹ ¹Lawrence Berkeley National Laboratory, 1 Cyclotron Rd, Berkeley, California 94720, USA ²Physics Department, University of California, Berkeley, California 94720, USA ³Frankfurt Institute for Advanced Studies, Ruth-Moufang-Strasse 1, 60438 Frankfurt am Main, Germany ⁴Institute for Theoretical Physics, Goethe University, Max-von-Laue-Strasse 1, 60438 Frankfurt am Main, Germany ⁵GSI Helmholtzzentrum für Schwerionenforschung, Planckstr. 1, 64291 Darmstadt, Germany



FIG. 1. Deuteron-pion interaction cross sections from SAID database [40] and partial wave analysis [41] are compared to our parametrizations (Tables II and III in the Appendix). Inelastic $d\pi \leftrightarrow$



FIG. 5. Reaction rates of the most important $\pi d \leftrightarrow \pi pn$ reaction in forward and reverse direction.

Law of mass action at work

Full calculation: nuclei



Deviations from thermal model predictions are moderate despite significant cooling and dilution. Is this the reason for why thermal model works so well?

For $T = T_{kin}$ similar results reported in [X. Xu, R. Rapp, EPJA 55, 68 (2019)]

Full calculation: hypernuclei



Hypernuclei stay close to the thermal model prediction. An exception is a hypothetical $\Xi\Xi$ state \leftarrow planned measurement in Runs 3 & 4 at the LHC [LHC Yellow Report, 1812.06772]

Full calculation: resonances



At $T \approx T_{kin}$ the suppressed resonance yields agree quite well with ALICE data for 0-20% central Pb+Pb collisions [ALICE, 1404.0495; 1805.04361; 1805.04365] This implies significant resonance regeneration in the hadronic phase

Summary and outlook

- Nucleosynthesis in HICs at LHC via the Saha equation is in analogy to initial stages of big bang nucleosynthesis in the early universe.
- This would naturally explain why thermal model works for light (anti-)(hyper-)nuclei yields. It does *not* establish where exactly they are formed though, $any T < T_{ch}$ permitted!
- Who can give the answer?
 - Building of clusters (Hagedorn states)? [K. Gallmeister et al.]
 - Reconciliation with coalescence?
 - Reaction rates vs expansion rate and binding energies?
 - Transport description? [D. Oliinychenko et al.]
 - Internal consistency of model assumptions? [Cai, Cohen, Gelman, Yamaguchi, 1905.02753]
- Quantum mechanical treatment of bound systems in medium needed

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Thanks for your attention!