Probing the QCD equation of state with fluctuations of conserved charges

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XXXII International Workshop on High Energy Physics "Hot problems of Strong Interactions"

November 12, 2020

Acknowledgements:

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Unterstützt von / Supported by



Alexander von Humboldt Stiftung/Foundation

QCD phase diagram



- Analytic crossover at vanishing net baryon density a first-principle result from lattice QCD
- Phase structure at finite density is largely unknown
- Probed by heavy-ion collisions

Fluctuations of conserved charges

Consider a random variable N

Cumulants: $K_N(t) = \ln \langle e^{tN} \rangle = \sum_{n=1}^{\infty} \kappa_n \frac{t^n}{n!}$ variance $\kappa_2 = \langle (\Delta N)^2 \rangle = \sigma^2$ widthskewness $\kappa_3 = \langle (\Delta N)^3 \rangle$ asymmetrykurtosis $\kappa_4 = \langle (\Delta N)^4 \rangle - 3 \langle (\Delta N^2) \rangle^2$ peak shape

Statistical mechanics:

Grand partition function

$$ln Z^{
m gce}(T, V, \mu) = ln \left[\sum_{N} e^{\mu N} Z^{
m ce}(T, V, N)\right],$$



Fluctuations probe finer details of the (QCD) equation of state

Common uses of thermal fluctuations

• Taylor expansion of the equation of state

$$\frac{p(T,\mu_B)}{T^4} = \frac{p(T,0)}{T^4} + \frac{\chi_2^B(T,0)}{2!}(\mu_B/T)^2 + \frac{\chi_4^B(T,0)}{4!}(\mu_B/T)^4 + \dots$$

 $\chi_n(T, \mu_B) = \frac{\partial^n(p/T^4)}{\partial(\mu_B/T)^n}$ - susceptibilities

• Fluctuation signals of the QCD critical point

$$\kappa_2 \sim \xi^2$$
, $\kappa_3 \sim \xi^{4.5}$, $\kappa_4 \sim \xi^7$, $\xi \to \infty$
[M. Stephanov, PRL '09]

• Chiral criticality at $\mu_B = 0$

[Friman, Karsch, Redlich, Skokov, EPJC '11]



Figure from VV, Anchishkin, Gorenstein, Poberezhnyuk, PRC '15

Constraining the excluded volume interactions in hadron resonance gas model

- Matching HRG to lattice QCD at heavy-ion freeze-out stage
- First step toward connecting thermal (grand-canonical) QCD fluctuations with experimental measurements

VV, M.I. Gorenstein, H. Stoecker, *Phys. Rev. Lett.* 118, 182301 (2017)
VV, A. Pasztor, S.D. Katz, Z. Fodor, H. Stoecker, *Phys. Lett. B* 775, 71 (2017)

Hadron resonance gas (HRG) model

• HRG model: free gas of known hadrons and resonances

$$p(T, \mu_B) = T \phi_M(T) + 2 T \phi_B(T) \cosh(\mu_B/T)$$

$$mesons$$

$$\phi_{M(B)}(T) = \sum \frac{d_i m_i^2 T}{2\pi^2} K_2(m_i/T)$$

- Hadronic interactions dominated by resonance formation*
- Single term in relativistic virial expansion
- Rich history in describing hadron yields in heavy-ion collisions

 $i \in M(B)$

- Matches well with lattice QCD below T_{pc}
- Net baryon fluctuations: Skellam distribution
 - $\chi^B_{2n} \propto \langle N_B \rangle + \langle N_{\bar{B}} \rangle, \ \chi^B_{2n-1} \propto \langle N_B \rangle \langle N_{\bar{B}} \rangle$
 - LQCD suggests breakdown of the model close to $T_{pc} \approx 155$ MeV



Excluded volume HRG model

EV-HRG model: incorporate repulsive baryon-baryon, antibaryon-antibaryon interactions [VV, Gorenstein, Stoecker, PRL '17; VV, Pasztor, Katz, Fodor, Stoecker, PLB '17]

$$p(\mathcal{T},\mu_B)=p^{\mathsf{id}}_{\mathcal{M}}(\mathcal{T})+p^{\mathsf{ev}}_{\mathcal{B}}(\mathcal{T},\mu_B)+p^{\mathsf{ev}}_{ar{\mathcal{B}}}(\mathcal{T},\mu_B)$$

$$p_{B(\bar{B})}^{\text{ev}} = p_{B(\bar{B})}^{\text{id}} e^{-bp_{B(\bar{B})}^{\text{ev}}/T} \quad \text{or} \quad p_{B(\bar{B})}^{\text{ev}} = \frac{T}{b} W[b \phi_B(T) e^{\pm \mu_B/T}]$$



Net baryon fluctuations no longer Skellam

$$\frac{\chi_4^B}{\chi_2^B} = \frac{1 - 8W(b\phi_B) + 6[W(b\phi_B)]^2}{[1 + W(b\phi_B)^4]} \simeq 1 - \frac{12b\phi_B(T)}{12b\phi_B(T)} + O(b^2\phi_B^2)$$

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 $p(T, \mu_B) = p_M^{\mathsf{id}}(T) + p_B^{\mathsf{ev}}(T, \mu_B) + p_{\bar{B}}^{\mathsf{ev}}(T, \mu_B)$

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EV-HRG vs lattice

WB: 1805.04445; HotQCD: 1708.04897



EV-HRG \approx **QCD** at $T \approx T_{pc}$

EV-HRG vs lattice



BONUS: Describes imaginary μ_B lattice data for the leading four virial coefficients

$$ho_B(T, \mu_B) = \sum_{k=1}^{\infty} b_k(T) \sinh(k\mu_B/T)$$

[VV, Pasztor, Katz, Fodor, Stoecker, PLB '



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Connecting fluctuations in heavy-ion collisions with grand-canonical susceptibilities

- Measurements affected by global conservation laws, thermal smearing, resonance decays, etc.
- Tackle these with a novel *particlization* routine for hydro

Heavy-ion collisions

- Heavy-ion collisions are commonly described with relativistic fluid dynamics
- Hydro expansion ends with a particlization of locally equilibrated QCD matter, this happens roughly at "chemical freeze-out", $T \approx 160$ MeV at highest energies
- Fairly successful in describing bulk observables like hadrochemistry, p_T spectra, flow etc.
- What about event-by-event fluctuations?
 - A (too) common approach: directly compare cumulant ratios measured in experiment with grandcanonical (lattice QCD) susceptibility ratios and hope for the best see e.g. [HotQCD coll., 2001.08530] and others



Theory vs experiment: Caveats

 proxy observables in experiment (net-proton, net-kaon) vs actual conserved charges in QCD (net-baryon, net-strangeness)

Asakawa, Kitazawa, PRC '12; VV, Jiang, Gorenstein, Stoecker, PRC '18

volume fluctuations

Gorenstein, Gazdzicki, PRC '11; Skokov, Friman, Redlich, PRC '13; Braun-Munzinger, Rustamov, Stachel, NPA '17

• non-equilibrium (memory) effects

Mukherjee, Venugopalan, Yin, PRC '15

• final-state interactions in the hadronic phase

Steinheimer, VV, Aichelin, Bleicher, Stoecker, PLB '18

accuracy of the grand-canonical ensemble (global conservation laws)

Jeon, Koch, PRL '00; Bzdak, Skokov, Koch, PRC '13; Braun-Munzinger, Rustamov, Stachel, NPA '17

coordinate vs momentum space (thermal smearing)

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When are the measured fluctuations grand-canonical?

- Consider event-by-event fluctuations of particle number in acceptance ΔY_{accept} around midrapidity
- Scales
 - ΔY_{accept} acceptance
 - ΔY_{total} full space
 - ΔY_{corr} rapidity correlation length (thermal smearing)
 - ΔY_{kick} diffusion in hadronic phase
- GCE applies if $\Delta Y_{total} \gg \Delta Y_{accept} \gg \Delta Y_{kick}, \Delta Y_{corr}$
- In practice neither $\Delta Y_{total} \gg \Delta Y_{accept}$ nor $\Delta Y_{accept} \gg \Delta Y_{corr}$ is satisfied
 - Corrections from global conservation are large [Bzdak et al., PRC '13]
 - $\Delta Y_{corr} \sim 1 \sim \Delta Y_{accept}$ [Ling, Stephanov, PRC '16]



Subensemble sampler

A novel particlization routine

 η_s

Partition into subensembles*

$$Z^{\text{tot}} = \prod_{i=1}^{N} \sum_{B_i} e^{\mu_i B_i / T} Z^{\text{ce}}(T_i, B_i, V_i) \times \delta(B_{\text{tot}} - \sum_{i=1}^{N} B_i)$$

*Following the idea put forward in [**VV**, Savchuk, Poberezhnyuk, Gorenstein, Koch, PLB '20]

- 1. Partition the hydro (blast-wave) particlization hypersurface into subvolumes along the space-time rapidity axis
- 2. Sample each subvolume grand-canonically, using the partition function of an *interacting* HRG (e.g. EV-HRG)
- 3. Reject the event if global conservation is violated
- 4. Sample the momenta of particles
- 5. Do resonance decays or plug into hadronic afterburner

- ✓ (event-by-event) hydro
- \checkmark locally grand-canonical fluctuations
- \checkmark global conservation
- ✓ thermal smearing
- ✓ resonance decays

VV, Koch, *to appear*

T

A case study: net proton/baryon fluctuations at LHC

Central Pb-Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV

- Particlization at T = 160 MeV, $\mu_B = 0$ •
- Rapidity axis partitioned into 96 slices, $\Delta \eta_s = 0.1$, $|\eta_s| < 4.8$ ٠
- Boost-invariant blast-wave hypersurface and flow profile •
- Sampling of (anti)baryons from the lattice-based EV-HRG ٠ model with global baryon conservation, 10^{10} events

 $P(N) \sim \frac{(V - bN)^{N}}{N!} \theta(V - bN)$ Poisson + rejection sampling details in VV, Gorenstein, Stoecker, 1805.01402

GCE baseline: $\frac{\kappa_2^B}{\langle B+\bar{B}\rangle} = 0.94$, $\frac{\chi_4^B}{\chi_2^B} = 0.69$, $\frac{\chi_6^B}{\chi_2^B} = -0.18 \leftarrow \text{compatible with lattice}$

BW parameters from Mazeliauskas, Vislavicius, 1907.11059



Net baryon fluctuations at LHC

- Global baryon conservation distorts the cumulant ratios already for one unit of rapidity acceptance
- Neglecting thermal smearing, effects of global conservation can be described analytically via SAM*

$$\frac{\kappa_2}{\langle B + \bar{B} \rangle} = (1 - \alpha) \frac{\kappa_2^{\text{gce}}}{\langle B + \bar{B} \rangle}, \qquad \alpha = \frac{\Delta Y_{\text{acc}}}{9.6}, \quad \beta \equiv 1 - \alpha$$
$$\frac{\kappa_4}{\kappa_2} = (1 - 3\alpha\beta) \frac{\chi_4^B}{\chi_2^B},$$
$$\frac{\kappa_6}{\kappa_2} = [1 - 5\alpha\beta(1 - \alpha\beta)] \frac{\chi_6^B}{\chi_2^B} - 10\alpha(1 - 2\alpha)^2\beta \left(\frac{\chi_4^B}{\chi_2^B}\right)^2$$

- Thermal smearing distorts the signal at $\Delta Y_{accept} \leq 1$
- Effect of resonance decays is negligible

*Subensemble acceptance method (SAM): VV, Savchuk, Poberezhnyuk, Gorenstein, Koch, PLB '20



Net baryon vs net proton

- Experiments measure protons as a proxy for baryons
- Protons form a subset of all baryons this dilutes the signal
- For example

 $\frac{\chi_4^B}{\chi_2^B}\Big|_{T=160MeV}^{\text{GCE}} \stackrel{\text{``lattice QCD''}}{\simeq 0.67} \neq \frac{\chi_4^B}{\chi_2^B}\Big|_{\Delta Y_{\text{acc}}=1}^{\text{HIC}} \simeq 0.56 \neq \frac{\chi_4^P}{\chi_2^P}\Big|_{\Delta Y_{\text{acc}}=1}^{\text{HIC}} \simeq 0.83$

- HIC net proton ≠ HIC net baryon ≠ LQCD net baryon
- Baryon cumulants can be reconstructed from proton cumulants via binomial (un)folding method of Kitazawa and Asakawa [Phys. Rev. C 85 (2012) 021901]



Comparison to ALICE data

- Model describes data within errors
- "Large" error bars and "small" $p_{\rm T}$ coverage do allow to distinguish the equation of state
- Future measurements will require larger acceptance





Summary

- HRG with a baryonic excluded volume matches lattice QCD susceptibilities in vicinity of the pseudocritical temperature
 - excluded-volume parameter $b = 1 \text{ fm}^3$
 - fluctuations at freeze-out in heavy-ion collisions
- Subensemble sampler is a novel particlization routine
 - event-by-event fluctuations in a fluid dynamical picture
 - global charge conservation, thermal smearing, resonance decays
- Quantitative analysis of fluctuations in Pb-Pb collisions at LHC
 - HIC net proton \neq HIC net baryon \neq LQCD/GCE net baryon
 - EV-HRG agrees with available experimental data





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Thanks for your attention!



