

Pion condensation in the early Universe at nonvanishing lepton flavor asymmetry

Volodymyr Vovchenko (LBNL)

TH Heavy Ion Coffee, CERN

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with B. Brandt, G. Endrödi (Bielefeld U.), F. Cuteri, F. Hajkarim, J. Schaffner-Bielich (Frankfurt U.)

based on [2009.02309](#)



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Bose-Einstein condensation

$$n(E_p) = \frac{1}{e^{(E_p - \mu)/T} - 1}$$

S.N. Bose, A. Einstein, 1924

$$\text{BEC at } T < T_c, \quad T_c \approx 3.31 \frac{n^{2/3}}{m}, \quad \frac{n_0}{n} = 1 - \left(\frac{T}{T_c} \right)^{3/2} \quad \text{for non-rel. Bose gas}$$

Velocity-distribution data for a gas of rubidium atoms confirming the discovery of BEC in 1995

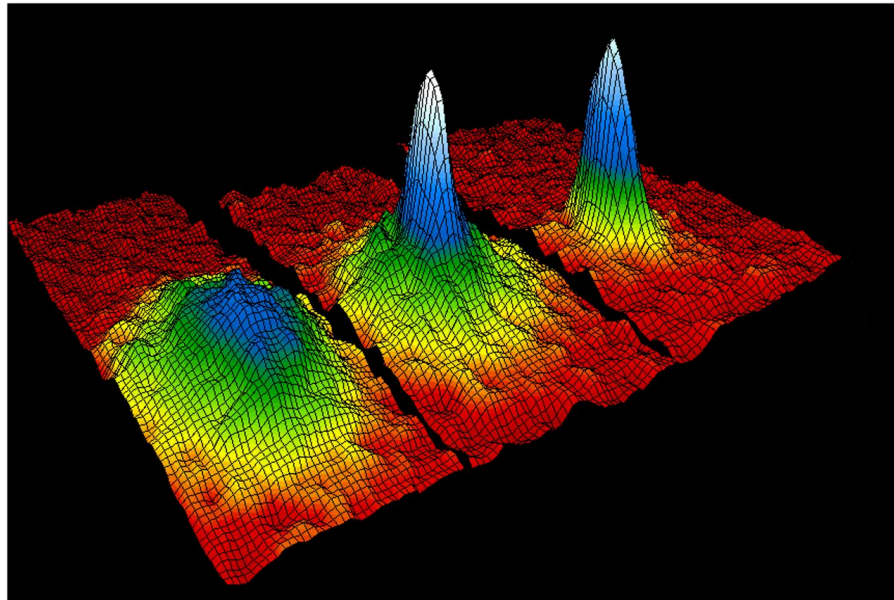


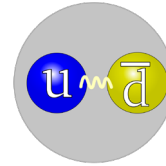
Image courtesy of NIST/JILA/CU-Boulder

Pion condensation

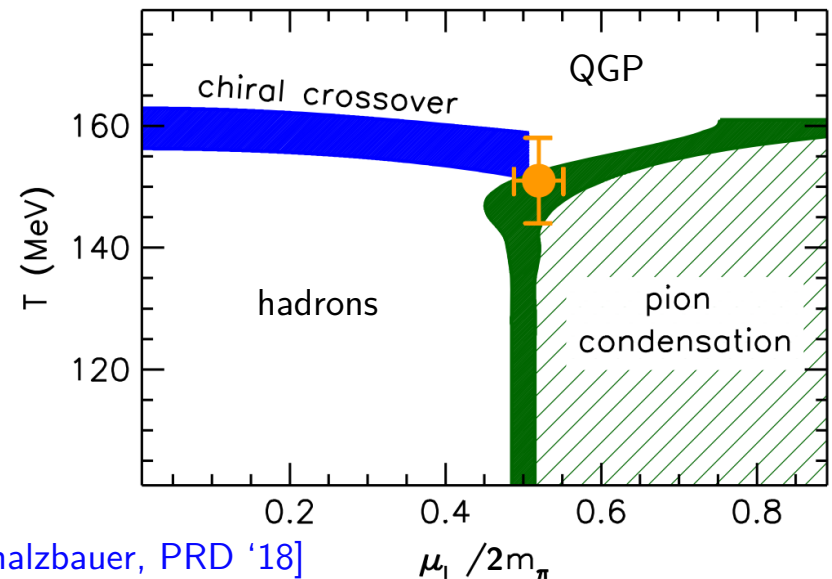
The relevant QCD degrees of freedom at low energies are **pions**

$$\mu_{\pi^\pm} = \pm \mu_I, \quad n_I = (n_u - n_d)/2$$

isospin



- chiral perturbation theory ($T=0$) [D.T. Son, M. Stephanov, PRL '01]
 - vacuum at $\mu_{\pi^\pm} < m_\pi$
 - **pion BEC** at $\mu_{\pi^\pm} \geq m_\pi$ (2nd order phase transition)
- Lattice QCD
 - no sign problem at finite μ_I
 - physical quark masses achieved
 - consistent with χ PT predictions

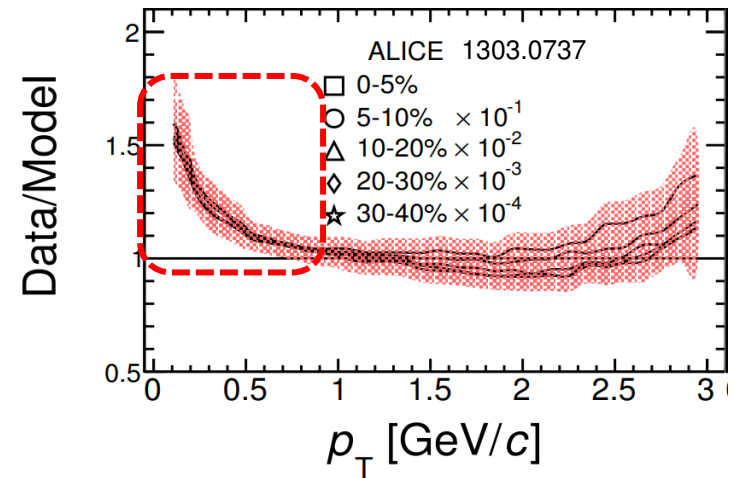


[Brandt, Endrodi, Schmalzbauer, PRD '18]

Pion condensation and heavy-ion collisions

- Low- p_T enhancement of pions produced in Pb-Pb collisions at LHC energies relative to hydro predictions

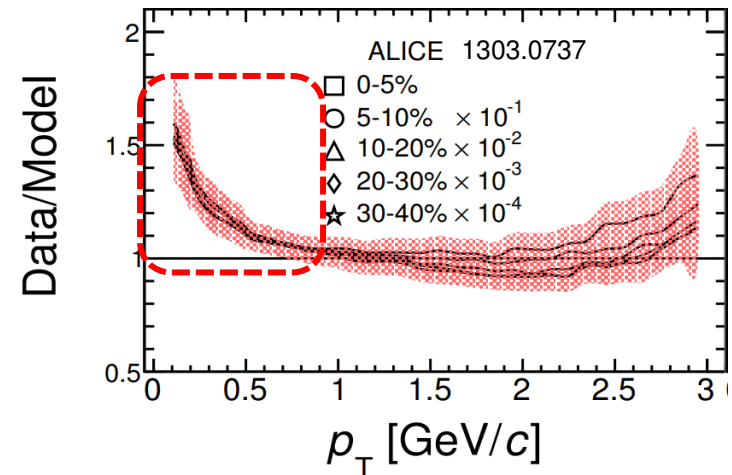
Figure from Devetak et al., JHEP '20



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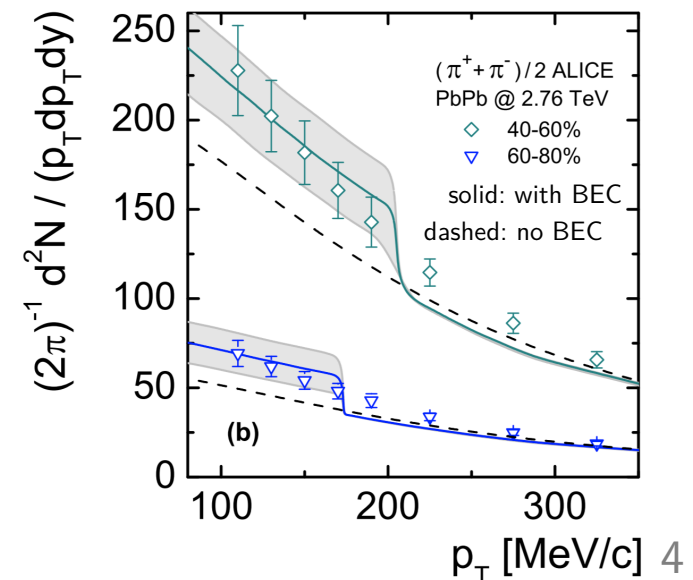
- Formation of a pion condensate may explain the data?

[Begun, Florkowski, Rybczynski, PRC '14, '15]

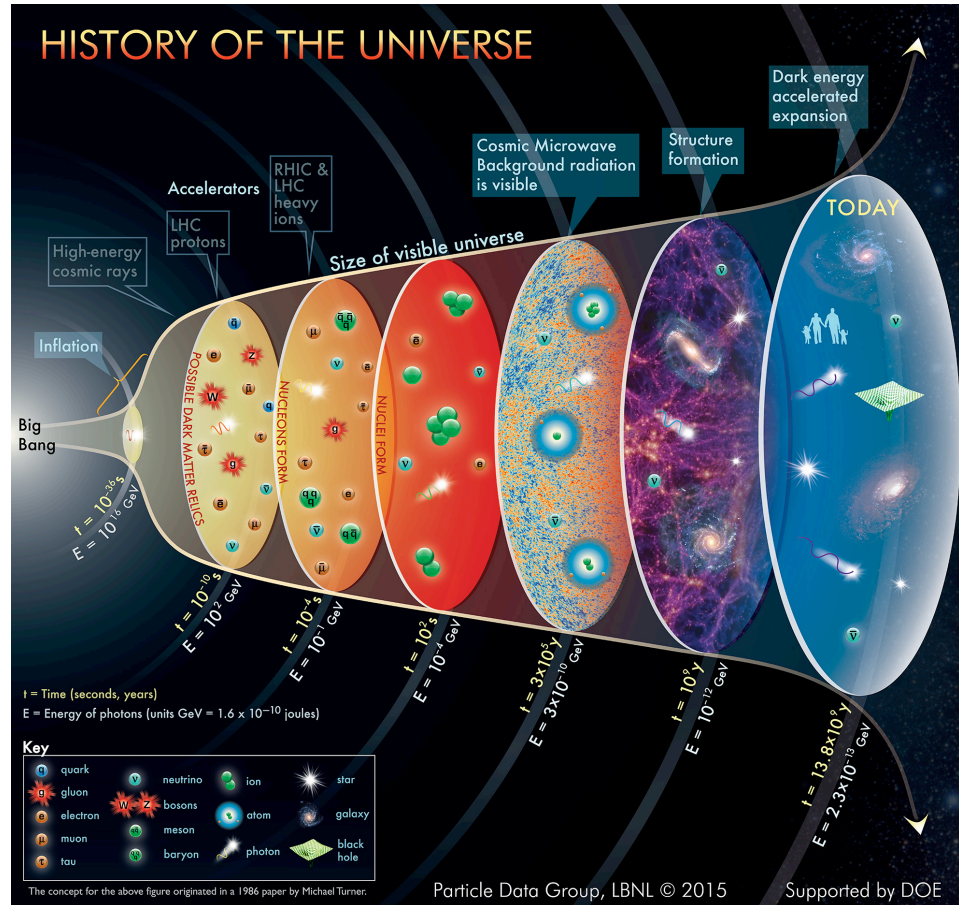
- But requires a large non-equilibrium pion chemical potential, e.g. an off-equilibrium hadronization of quark-gluon plasma

[Rafelski, Letessier, et al., EPJA '08, PRC '13]

Figure from Begun, Florkowski, PRC '14



Early Universe



QCD epoch: $\sim 10 \text{ MeV} < T < \sim 100 \text{ GeV}$

$\sim 10^{-11} \text{ s} < t < 1 \text{ s}$

Cosmic trajectories

- conservation equations for isentropic expansion

$$\frac{n_B}{s} = b, \quad \frac{n_Q}{s} = 0, \quad \frac{n_{L_\alpha}}{s} = l_\alpha \quad (\alpha \in \{e, \mu, \tau\})$$

- trajectory is a line in 6-dim space of temperature and chemical potentials

$$T, \quad \mu_B, \quad \mu_Q, \quad \mu_{L_\alpha}$$

- empirical constraints (CMB anisotropies)

$$b = (8.60 \pm 0.06) \cdot 10^{-11}$$

[Planck collab., 1502.01589]

$$|l_e + l_\mu + l_\tau| < 0.012$$

[Oldengott, Schwarz, 1706.01705]

- equation of state (QCD epoch)

$$p \approx p_{\text{QCD}} + p_{\text{leptons}} + p_{\text{photons}}$$

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Pion condensation may occur if $|\mu_Q| > m_\pi$ at $T < 160$ MeV

As pointed out in [M. Wygas et al., PRL '18]

Modeling the cosmic equation of state

$$p \approx p_{\text{QCD}} + p_{\text{leptons}} + p_{\text{photons}}$$

- leptons

$$p_{\text{leptons}}(T, \mu_Q, \mu_{L_\alpha}) = \sum_{\alpha \in \{e, \mu, \tau\}} [p_\alpha^{\text{id}}(T, \mu_Q, \mu_{L_\alpha}) + p_{\nu_\alpha}^{\text{id}}(T, \mu_{L_\alpha})] + \text{antiparticles}$$

- photons

$$p_\gamma(T) = \frac{\pi^2}{45} T^4$$

- QCD?

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The typical model of choice for hadronic matter is hadron resonance gas (HRG)

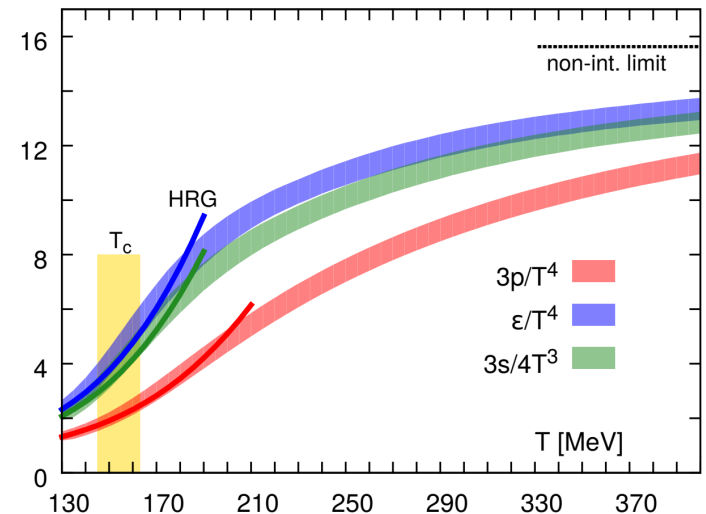


Figure from HotQCD coll., PRD '14

Strategy: Implement pion-pion interactions into the HRG model to account for the pion-condensed phase

Effective mass model for pion condensation

- **A quasiparticle picture:** pion interactions are driven by *effective mass*:

$$p_{\pi}^{\text{EM}}(T, \mu_{\pi}; m^*) = p_{\pi}^{\text{id}}(T, \mu_{\pi}; m^*) + p_f(m^*)$$

rearrangement term

$$m^*(T, \mu_{\pi}) \text{ from gap equation, } \frac{\delta p_{\pi}}{\delta m^*} = 0: \quad p'_f(m^*) = n_{\sigma}^{\text{id}}(T, \mu_{\pi}; m^*)$$

scalar density

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scalar density

- Onset of **pion condensation** takes place when chemical potential becomes equal to the effective mass, $\mu_{\pi} = m^*$. This gives the **Bose-Einstein condensation line**:

$$T_{\text{cond}}(\mu_{\pi}) : \quad p'_f(\mu_{\pi}) = n_{\sigma}^{\text{id}}[T_{\text{cond}}(\mu_{\pi}), \mu_{\pi}; m^* = \mu_{\pi}]$$

- $T < T_{\text{cond}}$: a fraction of pions forms a Bose-Einstein condensate, $n_{\pi} = n_{\pi}^{\text{th}} + n_{\pi}^{\text{BEC}}$

$$n_{\pi}^{\text{th}} = n^{\text{id}}(T, \mu_{\pi}; m^* = \mu_{\pi}) \quad n_{\pi}^{\text{BEC}} = p'_f(\mu_{\pi}) - n_{\sigma}^{\text{id}}(T, \mu_{\pi}; m^* = \mu_{\pi})$$

thermal pions *condensed pions*

The specific form of the rearrangement term $p_f(m^*)$ defines the model

Effective mass model: $T = 0$

No thermal excitations at $T = 0$, only condensed pions at $\mu_\pi > m_\pi$

$$n_\pi^{\text{EM}}(T = 0, \mu_\pi) = p'_f(\mu_\pi) \theta(\mu_\pi - m_\pi)$$

$$\chi\text{PT:} \quad n_{\chi\text{PT}}(T = 0, \mu_\pi) = \frac{\mu_\pi f_\pi^2}{2} \left[1 - \frac{m_\pi^4}{\mu_\pi^4} \right] \theta(\mu_\pi - m_\pi)$$

[D.T. Son, M. Stephanov, PRL '01]

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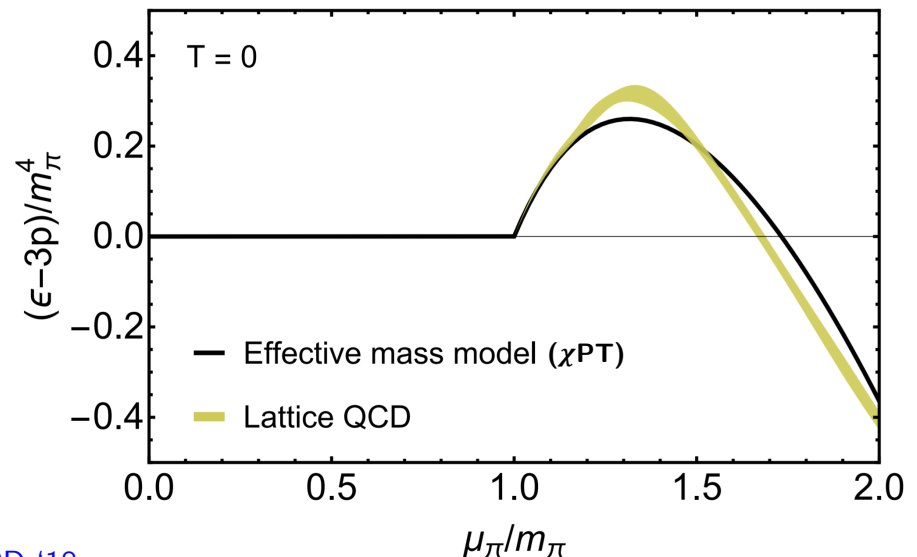
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[D.T. Son, M. Stephanov, PRL '01]

Match the effective mass model to chiral perturbation theory at $T = 0$:

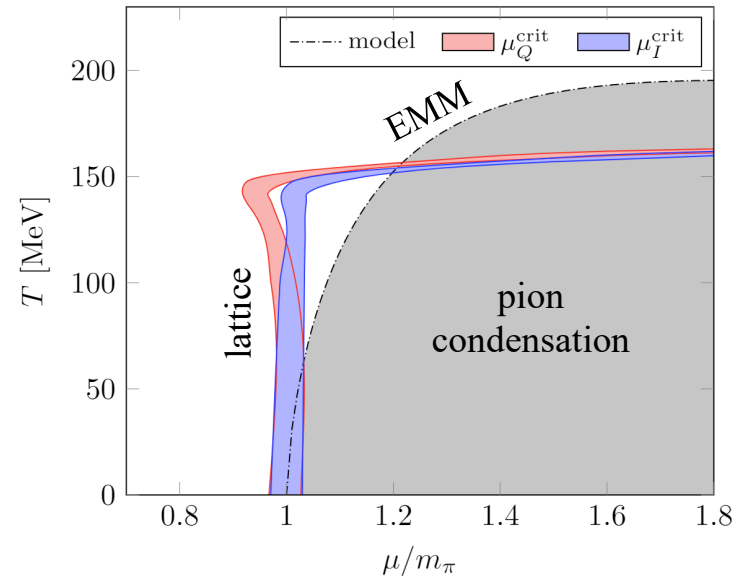
$$p_f(\mu_\pi) = \frac{\mu_\pi^2 f_\pi^2}{4} \left[1 - \frac{m_\pi^2}{\mu_\pi^2} \right]^2$$

$$f_\pi = 133 \text{ MeV}$$

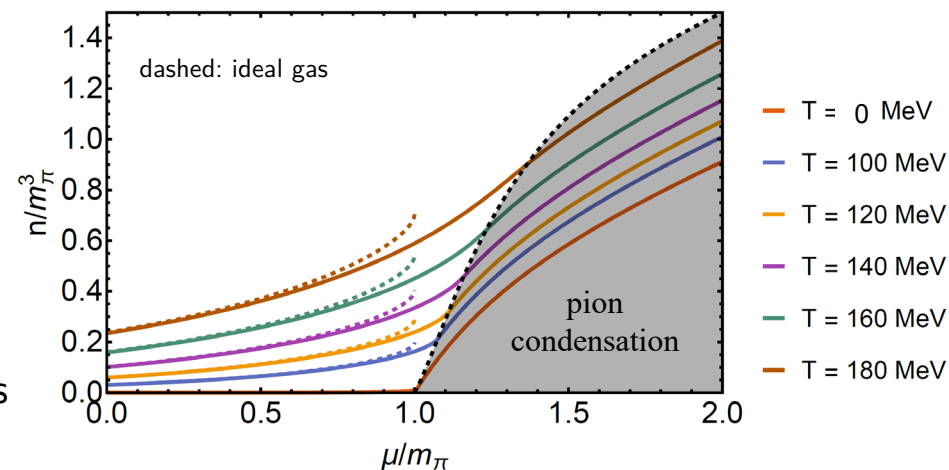


Effective mass model: Phase diagram

- Pion condensation boundary
 - Qualitatively similar to lattice QCD
 - Not as abrupt leveling off as on lattice*
 - Model has no deconfinement, thus not reliable at $T > 160$ MeV



- Order of the transition
 - Kink in $n_\pi(\mu_\pi)$ at zero temperature \rightarrow 2nd order phase transition
 - Does not turn 1st order at finite T
 - Consistent with lattice QCD observations



*See PQM type models for a more involved modeling of the transition line e.g. [\[Adhikari, Andersen, Kneschke, 1805.08599\]](#)

HRG model with pion interactions

$$p_{\text{QCD}}(T, \mu_B, \mu_Q) \approx \sum_{\substack{i \in \pi^\pm, \pi^0 \\ \text{interacting pions}}} p_i^{\text{EM}}(T, \mu_i) + \sum_j p_j^{\text{id}}(T, \mu_j).$$

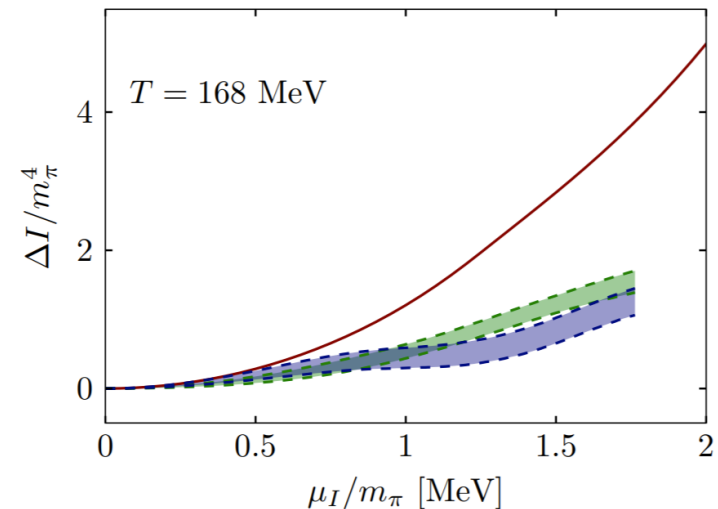
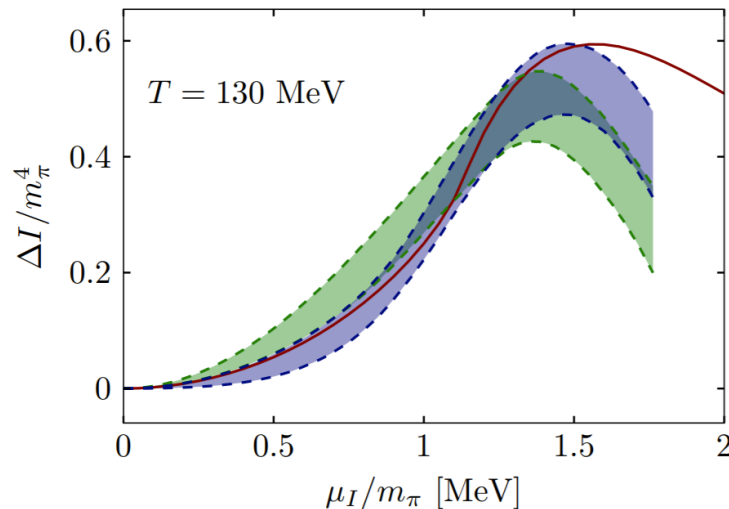
free hadrons and resonances

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free hadrons and resonances

- $\Delta I = I(T, \mu_I) - I(T, 0)$, $I \equiv \varepsilon - 3p$
- Two lattice spacings: $N_t = 10$, $N_t = 12$

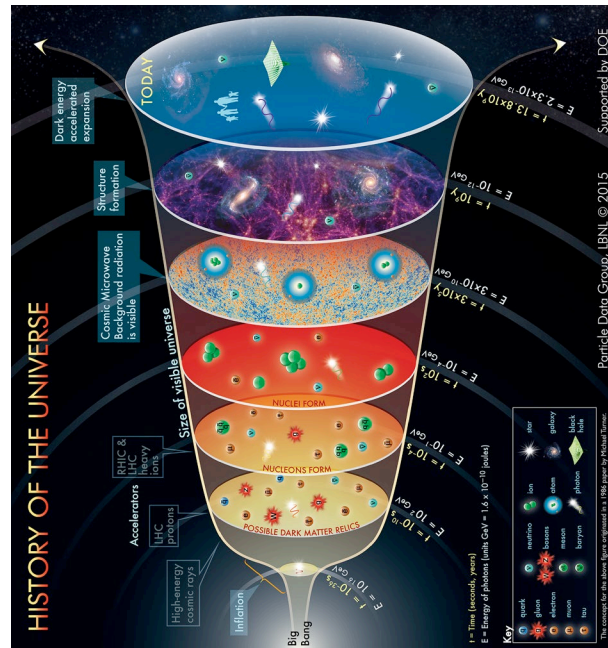


- Validity range of the model:

$$T \lesssim 160 \text{ MeV}, \quad \mu_I \lesssim 1.5 m_\pi$$

Calculating the cosmic trajectories

Early universe



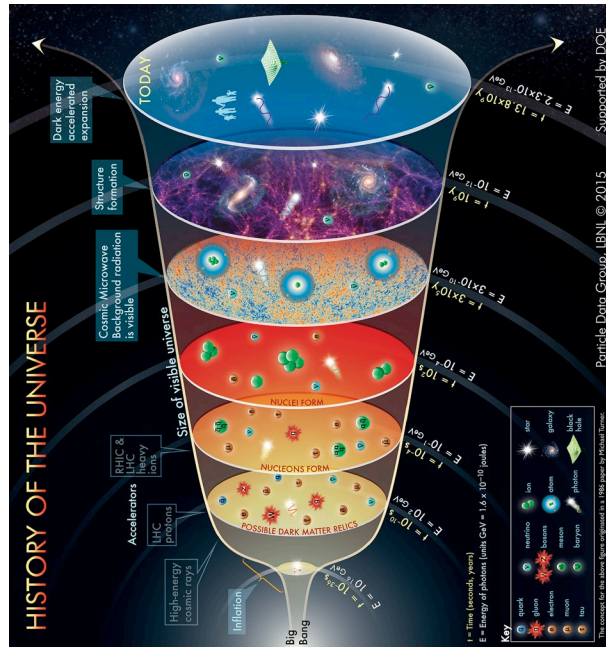
PDG

$$\frac{n_B}{s} = b, \quad \frac{n_Q}{s} = 0, \quad \frac{n_{L\alpha}}{s} = l_\alpha$$

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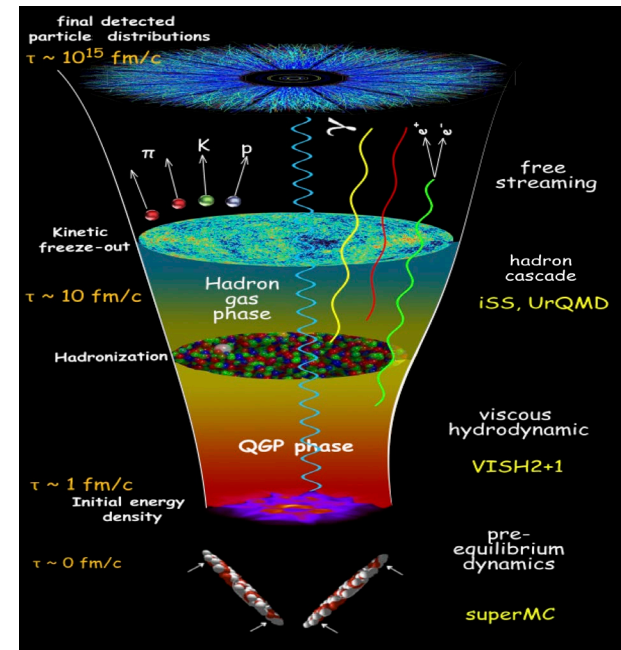


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Heavy-ion collision



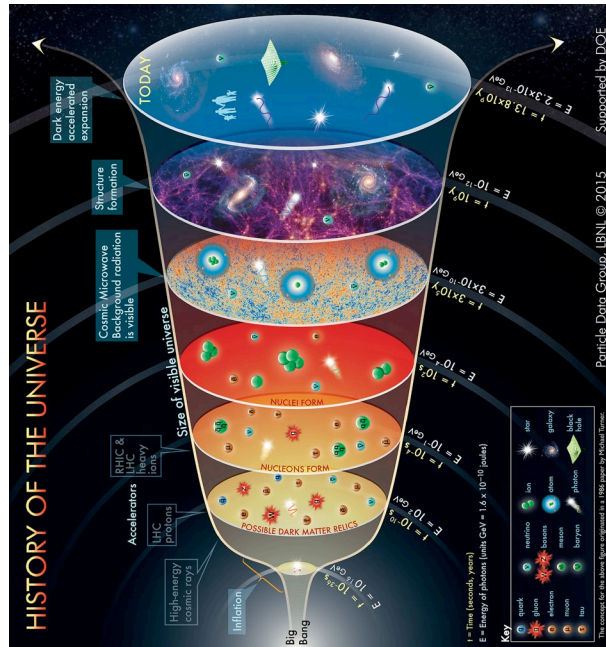
C. Shen, Ohio State U.

$$\frac{n_B}{s} = b, \quad \frac{n_Q}{s} = q, \quad \frac{n_S}{s} = 0$$

$$p = p_{\text{QCD}}$$

Calculating the cosmic trajectories

Early universe

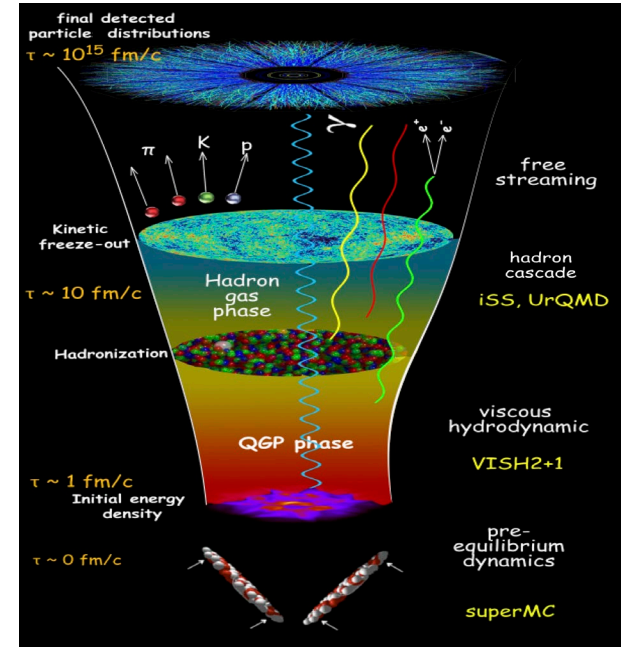


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Cosmic trajectories implemented within (extended) **Thermal-FIST** package

[V.V., H. Stoecker, *Computer Physics Communications* **244**, 295 (2019); [github link](#)]

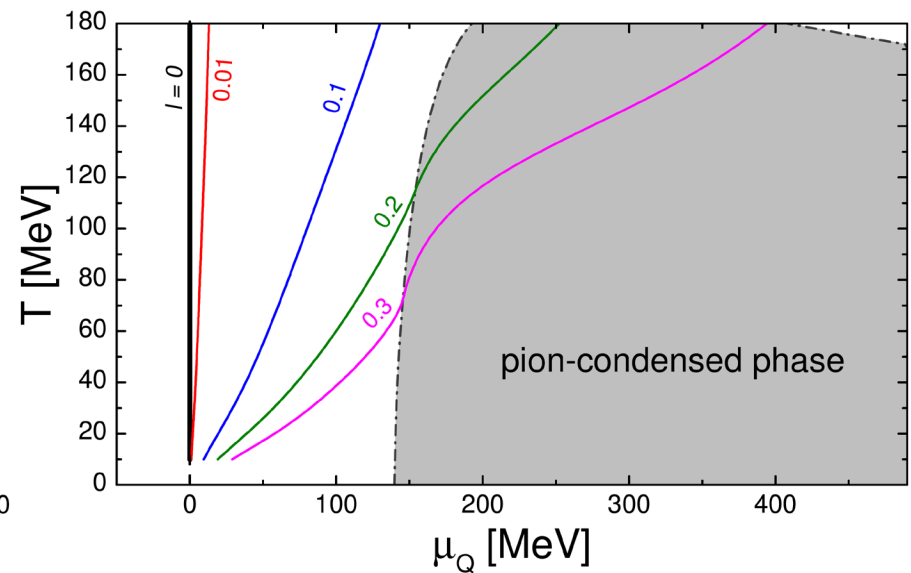
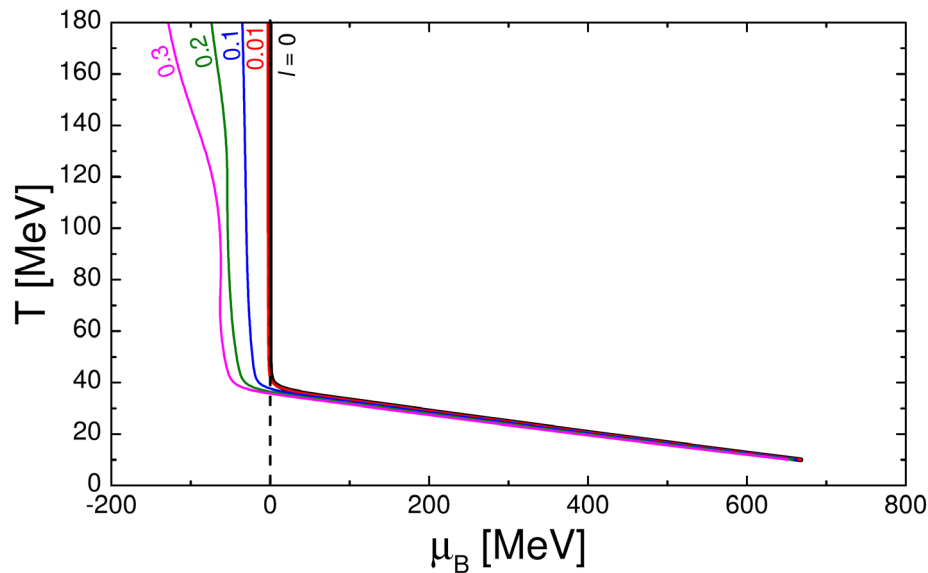


Using a *heavy-ion* tool in *cosmology*

Trajectories: Lepton-flavor symmetric case

Fix $b = 8.6 \cdot 10^{-11}$ and do a parametric scan in lepton asymmetries

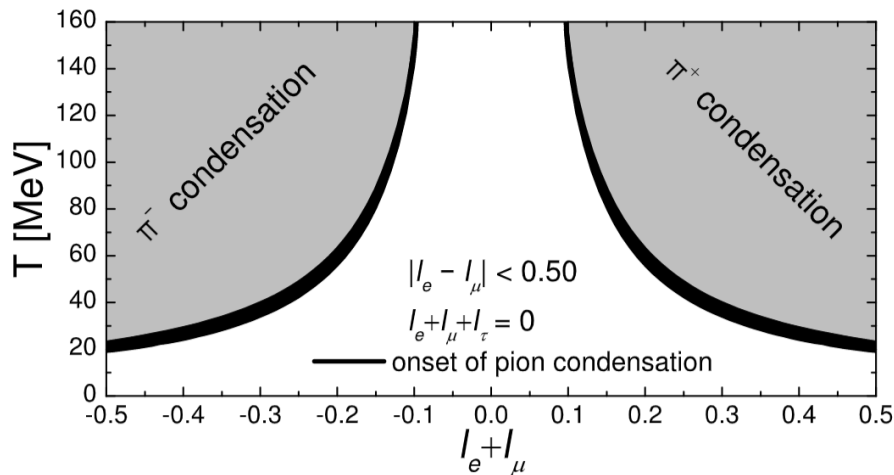
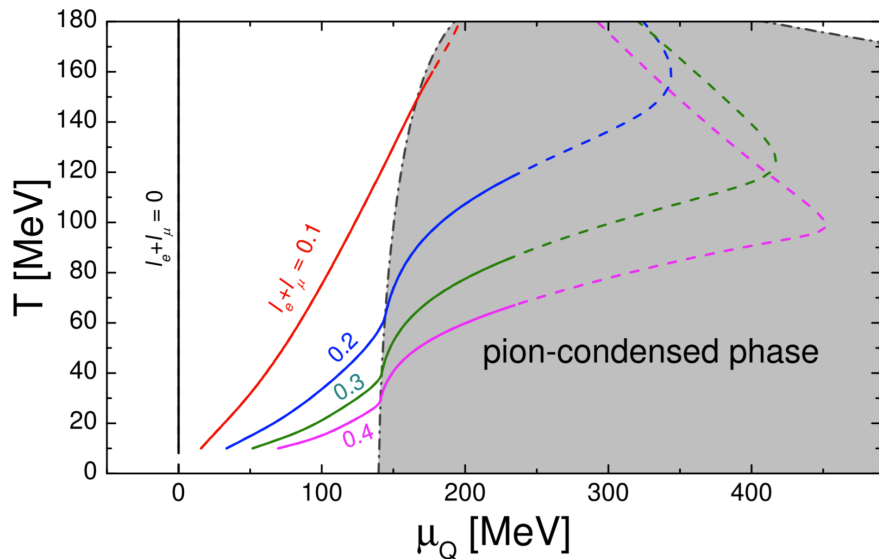
First consider $l_e = l_\mu = l_\tau = l/3$



- Pion condensation in the symmetric scenario occurs if $|l| > \sim 0.15$
- However, this violates the empirical constraint $|l| < 0.012$

Outside pion-condensed region reproduces HRG model results of [M. Wygas et al., PRL '18; 2009.00036]

Trajectories: Lepton-flavor asymmetric case

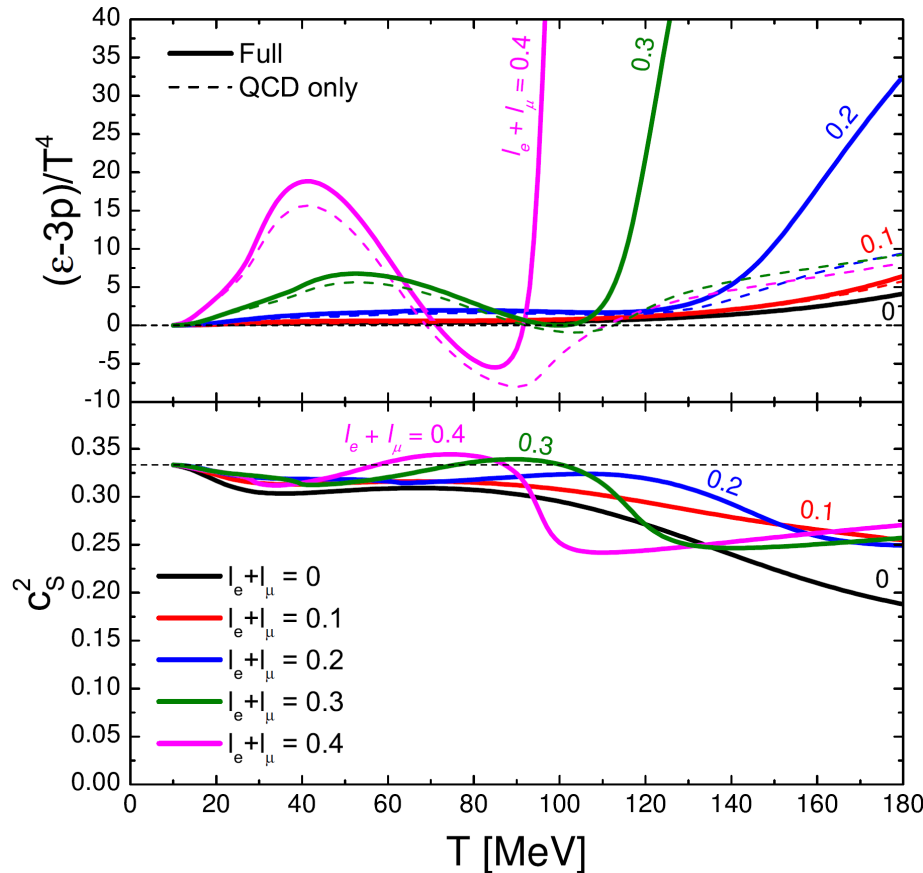


- Individual lepton flavor asymmetries are much less constrained
- Set total lepton asymmetry to zero but vary individual flavor ones
 $l_e + l_\mu + l_\tau = 0$ but $l_e \neq l_\mu \neq l_\tau$
- 2D scan in $(l_e + l_\mu, l_e - l_\mu)$

Pion condensation occurs if

$$|l_e + l_\mu| \gtrsim 0.1$$

Lepton-flavor asymmetric case: Cosmic EoS



- Cosmic equation of state affected strongly by large lepton asymmetry
- Pion condensation leads to (nearly) negative interaction measure and $c_s^2 > 1/3$
- At higher temperatures large I/T^4 driven by large lepton chemical potentials

The changed **EoS** has cosmological implications

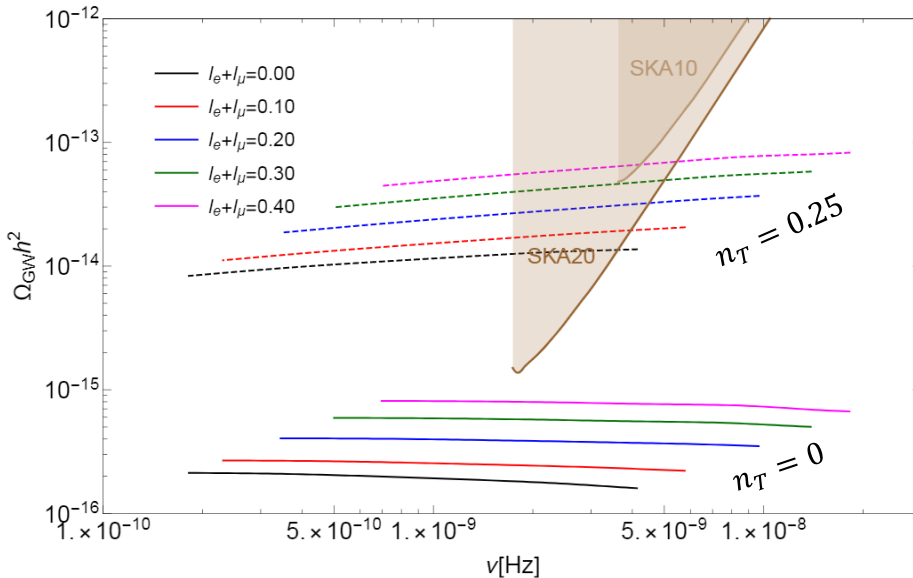
$$s(T)[a(T)]^3 = \text{const} \quad H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3} \epsilon$$

scale factor *Hubble rate*

$$\frac{a''}{a} = \frac{4\pi G}{3} (\epsilon - 3p)$$

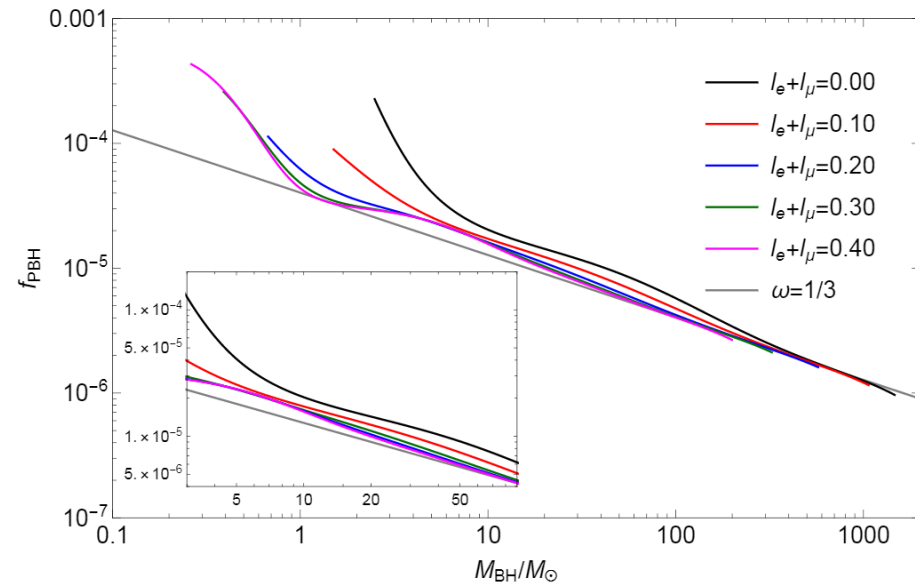
Primordial gravitational waves and black holes

Primordial gravitational waves from inflationary scenario



- Enhanced relic density of primordial gravitational waves (relative to amplitude at $l_e + l_\mu = 0$)
- Possibly reachable by pulsar-timing arrays, e.g. Square Kilometer Array (SKA) over 10-20 years of operation

Fraction of primordial black holes relative to cold dark matter



- Changed fraction of primordial black holes heavier than solar mass
- Pion condensation epoch is a source of PBHs?
- Speculation: BHs merger events, e.g. LIGO GW190521, ...

Production mechanism for pion stars

- Pion stars are gravitationally bound objects whose main constituent is the Bose-Einstein condensate of charged pions
[Carignano et al., 1610.06097; Brandt et al., 1802.06685; Andersen, Kneschke, 1807.08951]
- Pion condensation serves as a primordial production mechanism

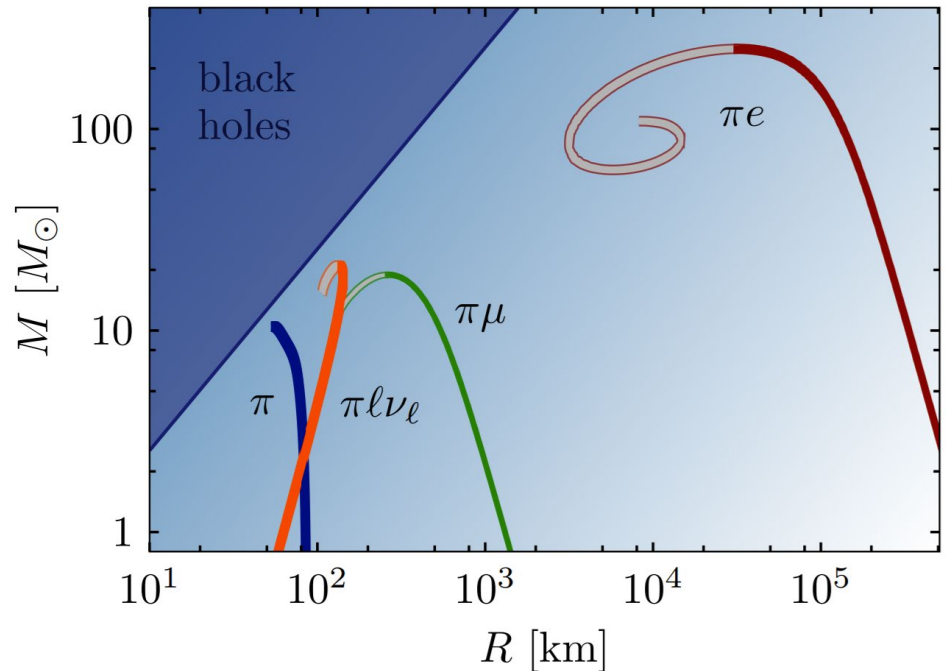


Figure from Brandt et al., 1802.06685

- If pion stars decay around the time of big bang nucleosynthesis, the produced high energy leptons can influence the primordially produced nuclei

Summary

- The early universe passes through a pion-condensed phase if electron and muon lepton asymmetry is sufficiently large:

$$|l_e + l_\mu| > 0.1$$

- **Implications:**

- Large effect on the pre-BBN equation of state
- Enhanced relic density of primordial gravitational waves (relative to amplitude at $l_e + l_\mu = 0$)
- Changed fraction of primordial black holes with mass larger than M_\odot
- Possible formation and decay of pion stars, effect on big bang nucleosynthesis

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Backup slides

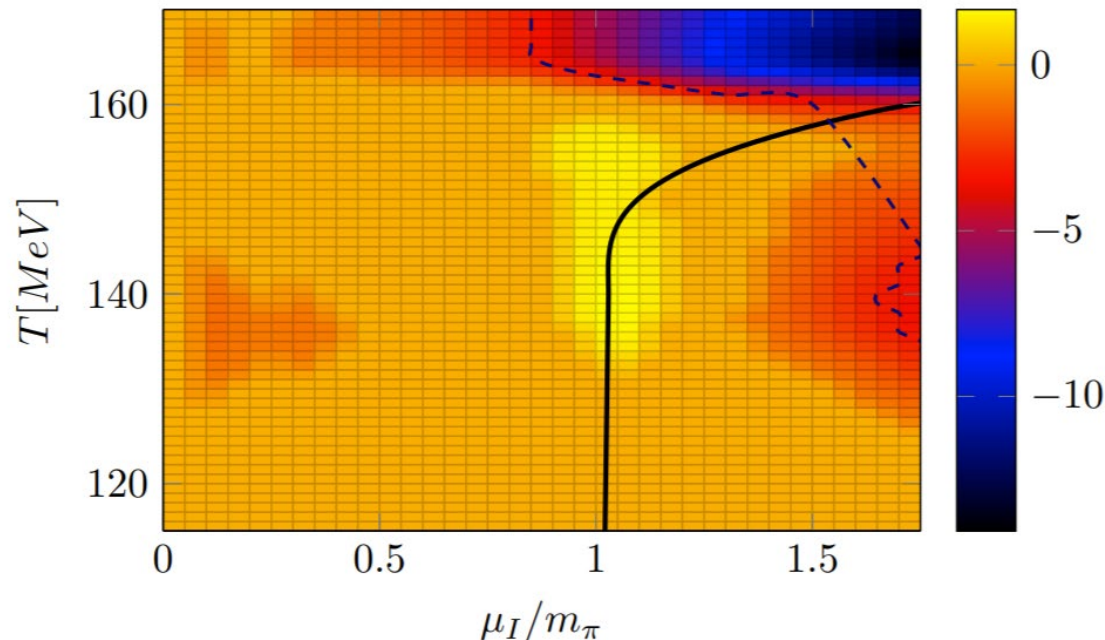
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$$(\Delta I - \Delta I^{\text{EM}})/\sigma(\Delta I)$$



Lepton chemical potentials

