

QCD phase diagram with constant entropy contours and fluctuations in heavy-ion collisions

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YITP workshop “Buenas Ideas on the QCD Phase Diagram”, Kyoto, Japan

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Thanks to: M. Kahangirwe, Y. Fujimoto, G. Pihan, R. Poberezhniuk, H. Shah, ...



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QCD under extreme conditions

What we know

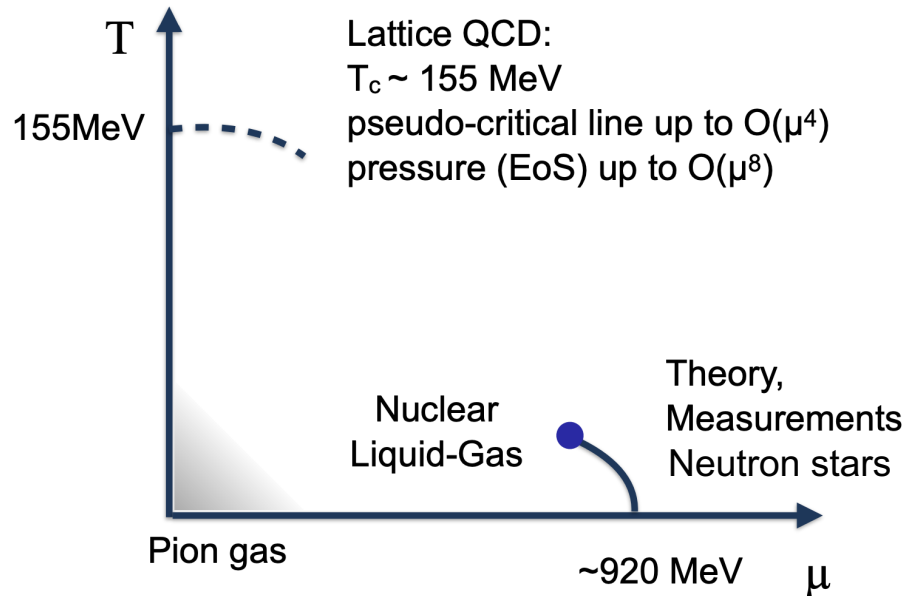


Figure courtesy of V. Koch

What we hope to know

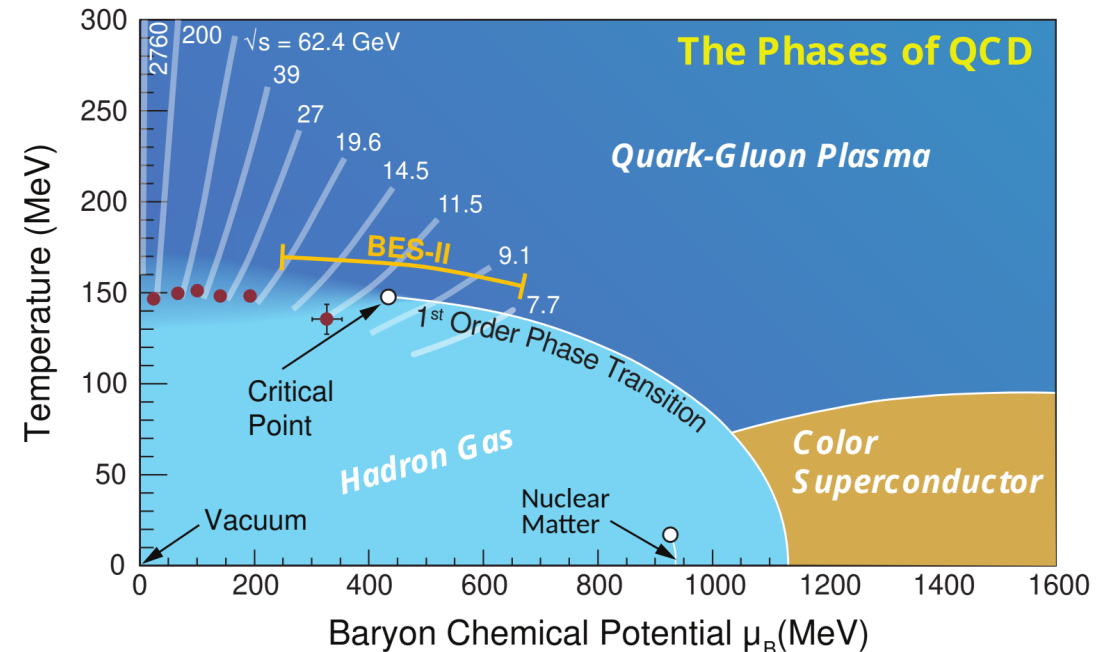


Figure from Bzdak et al., Phys. Rept. '20 & 2015 US Nuclear Long Range Plan

- Dilute hadron gas at low T & μ_B due to confinement, quark-gluon plasma high T & μ_B
- Nuclear liquid-gas transition in cold and dense matter, lots of other phases conjectured
- Chiral crossover at $\mu_B = 0$ which may turn into a *first-order phase transition* at finite μ_B

QCD under extreme conditions

What we know

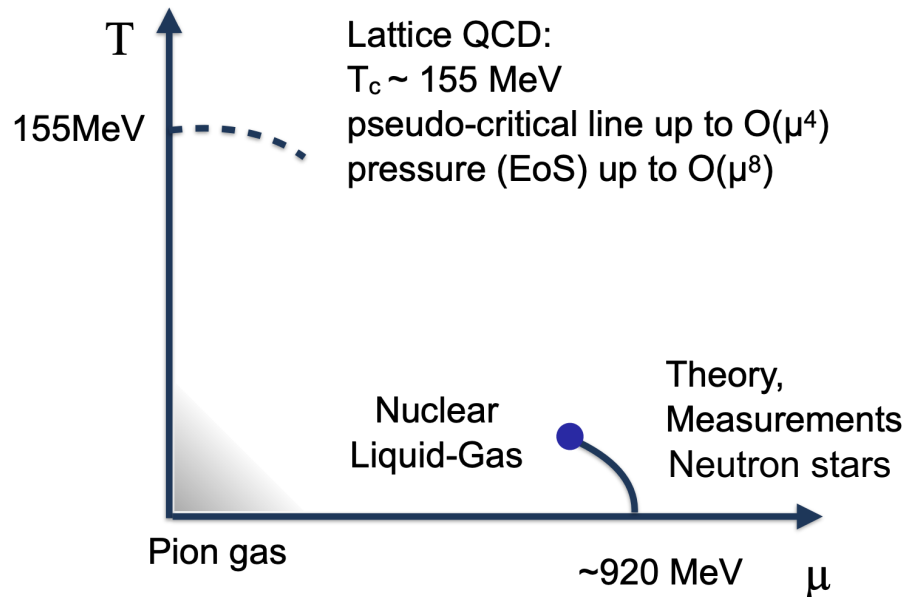
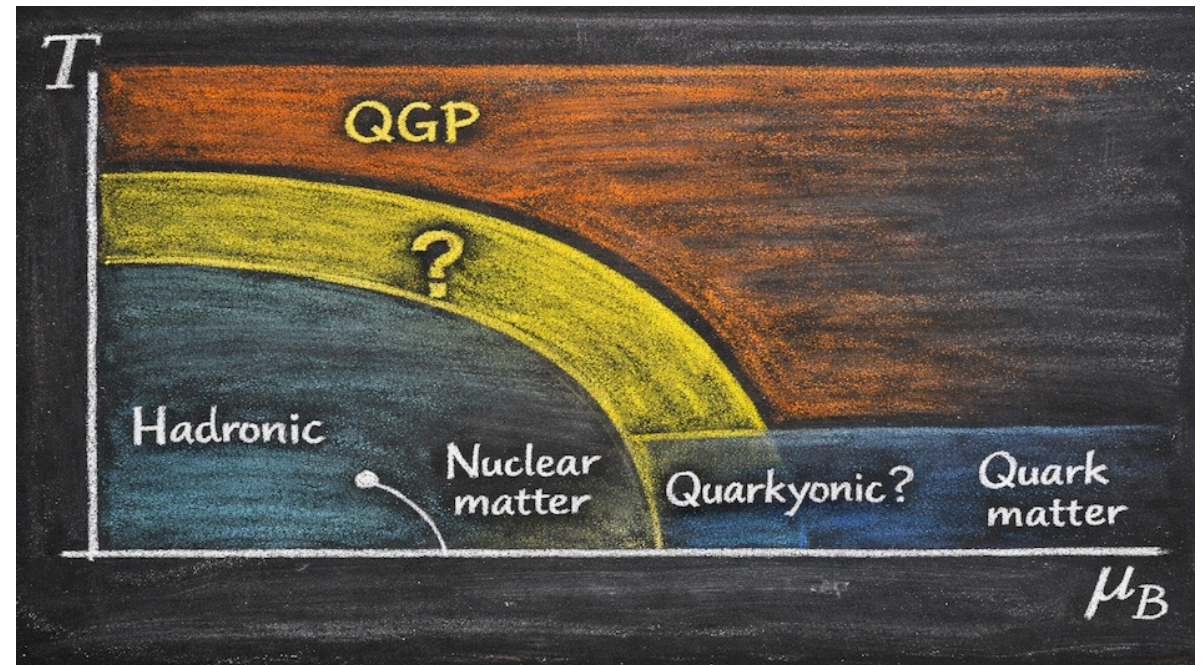


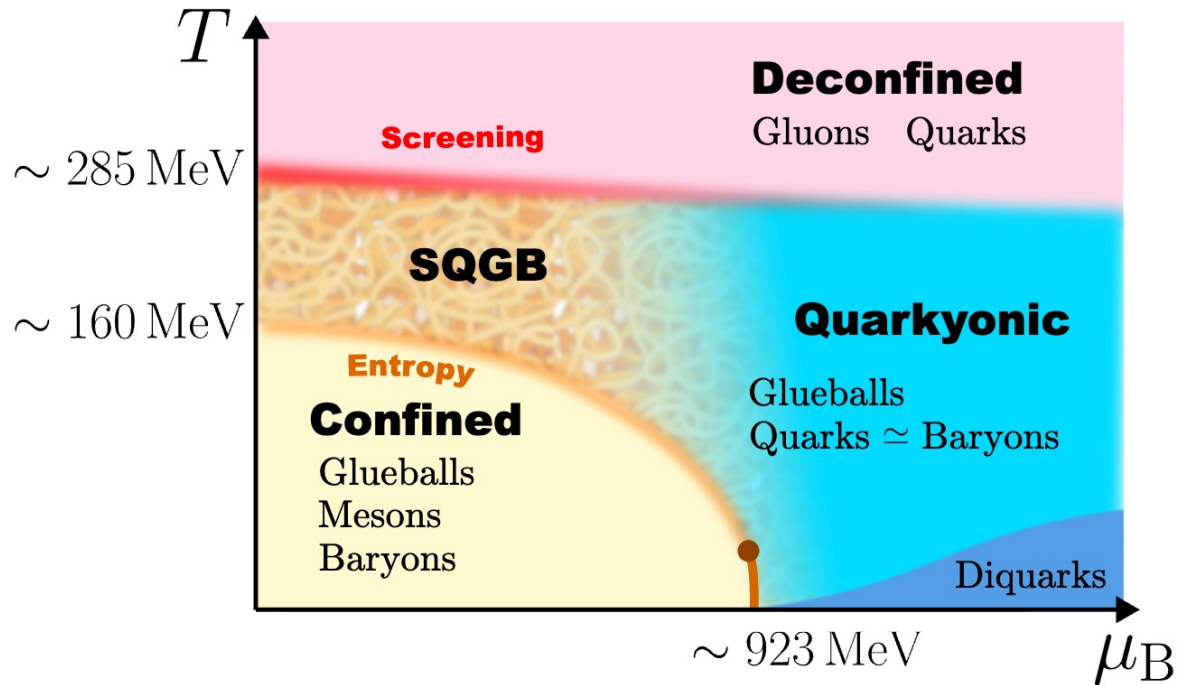
Figure courtesy of V. Koch

What we hope to know

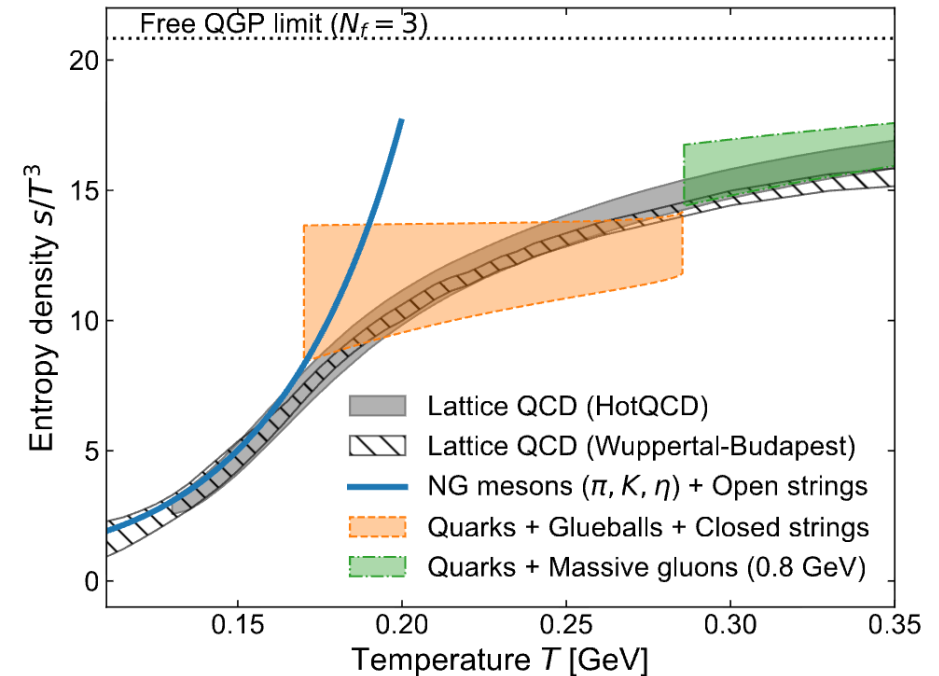


Conference poster "Buenas Ideas on the QCD Phase Diagram"

- Dilute hadron gas at low T & μ_B due to confinement, quark-gluon plasma high T & μ_B
- Nuclear liquid-gas transition in cold and dense matter, lots of other phases conjectured
- Chiral crossover at $\mu_B = 0$ which may turn into a *first-order phase transition* at finite μ_B
- Or perhaps a different state of matter between T_c and $(2-3)T_c$



Fujimoto, Fukushima, Hidaka, McLerran, PRD 112, 074006 (2025)

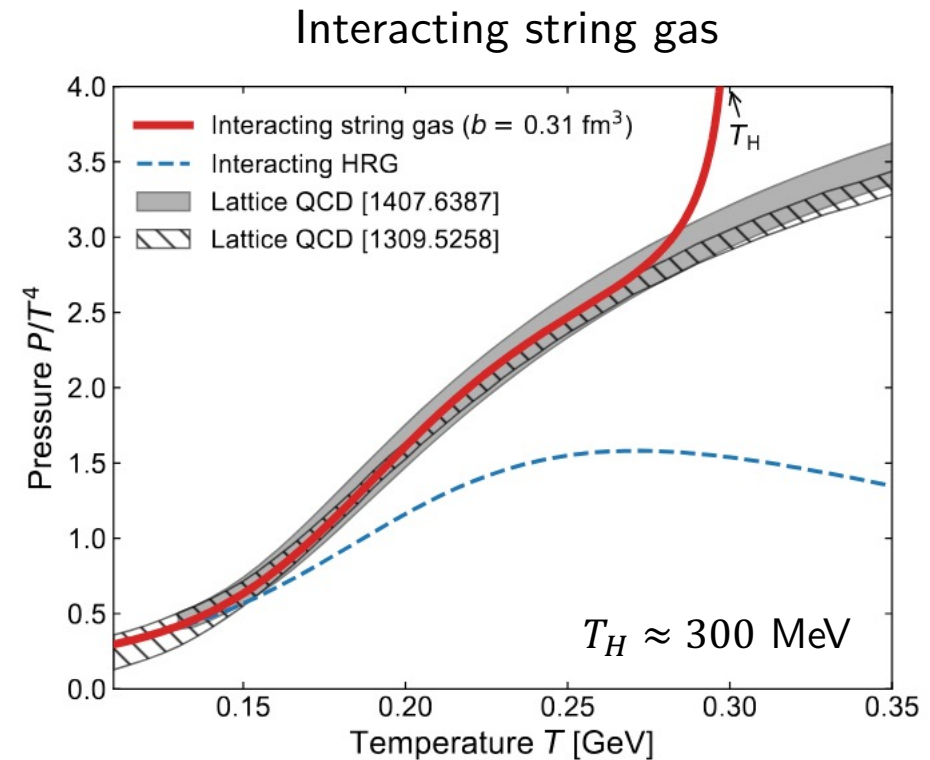
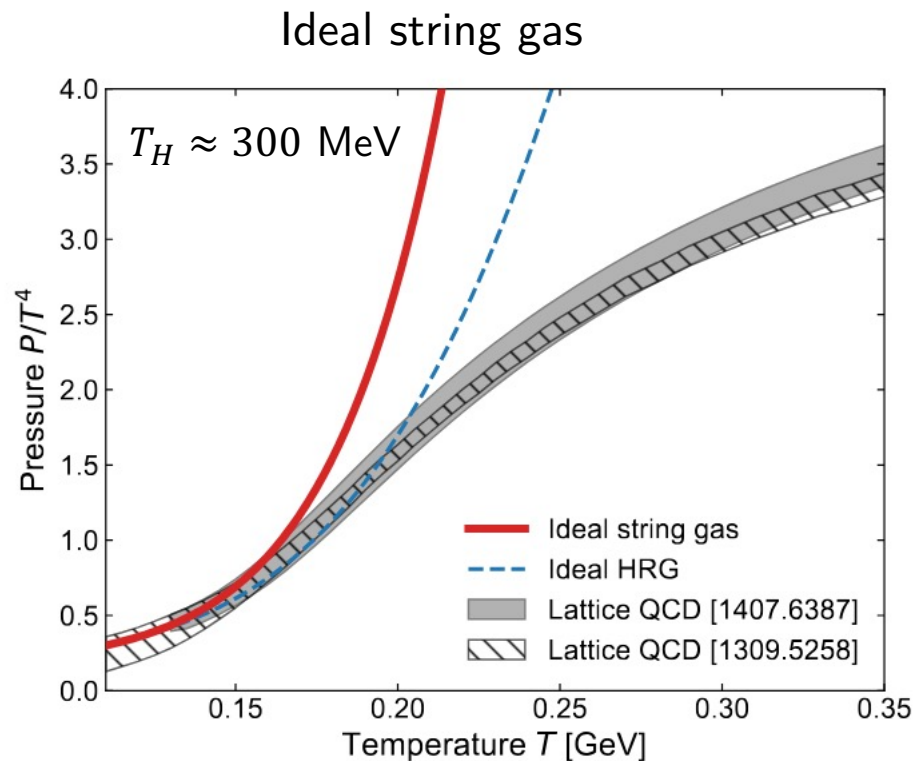


- At low T the HRG description well approximated by a gas of confining strings with a single Hagedorn-like temperature $T_H \sim 280 - 340$ MeV
- Works for glueballs [H. Meyer, PRD 80, 051502 (2009)], mesons [Marczenko et al., PRD 112, 096010, (2025)], baryons [Fujimoto, PRD 113, 054044 (2026)] and perhaps other flavors [Marczenko et al., 2603.28668].
- IM dilepton measurements point to $T_{IM} \approx T_H \approx 300$ MeV across broad range of energies

HRG as a gas of confining strings with interaction

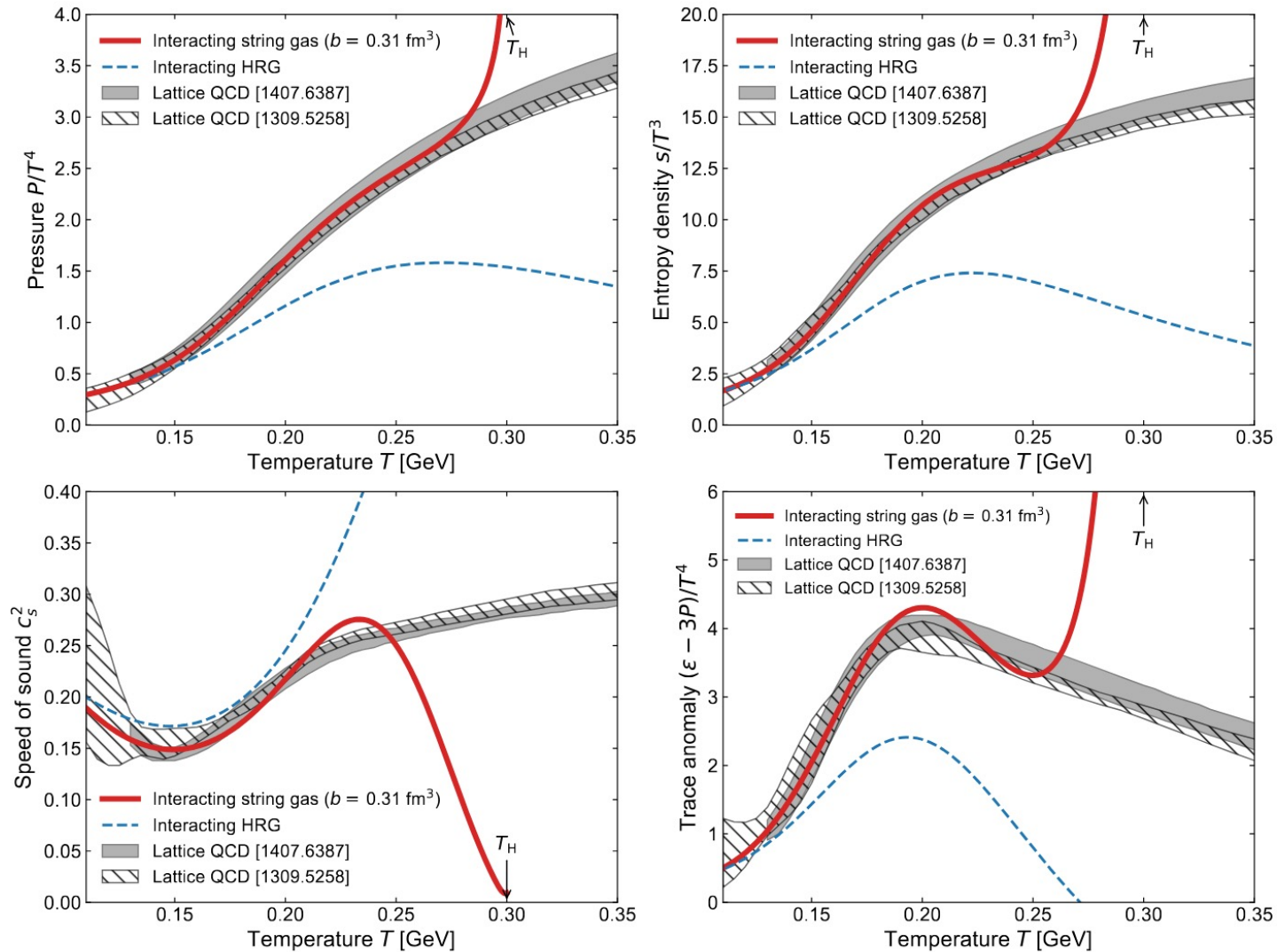
- Strings with excluded volume, a la EV/VDW-HRG model [VV, Gorenstein, Stoecker, PRL 118, 182301 (2017)]

$$P(T) = T n_{\text{id}}(T) \exp\left(-\frac{bP}{T}\right) \quad T\phi_i(T) = \int_{m_{\text{str},i}}^{\infty} dm \rho_{\text{str}}(m) d_{\text{str},i} P(m) \quad \rho_{\text{str}}(m) = \frac{\sqrt{2\pi}}{6T_H} \left(\frac{m}{T_H}\right)^{-3/2} e^{m/T_H}$$



Description of thermodynamics lowers the Hagedorn temperature to $T_H \approx 300$ MeV relative to $T_H \approx 340$ MeV from PDG fits

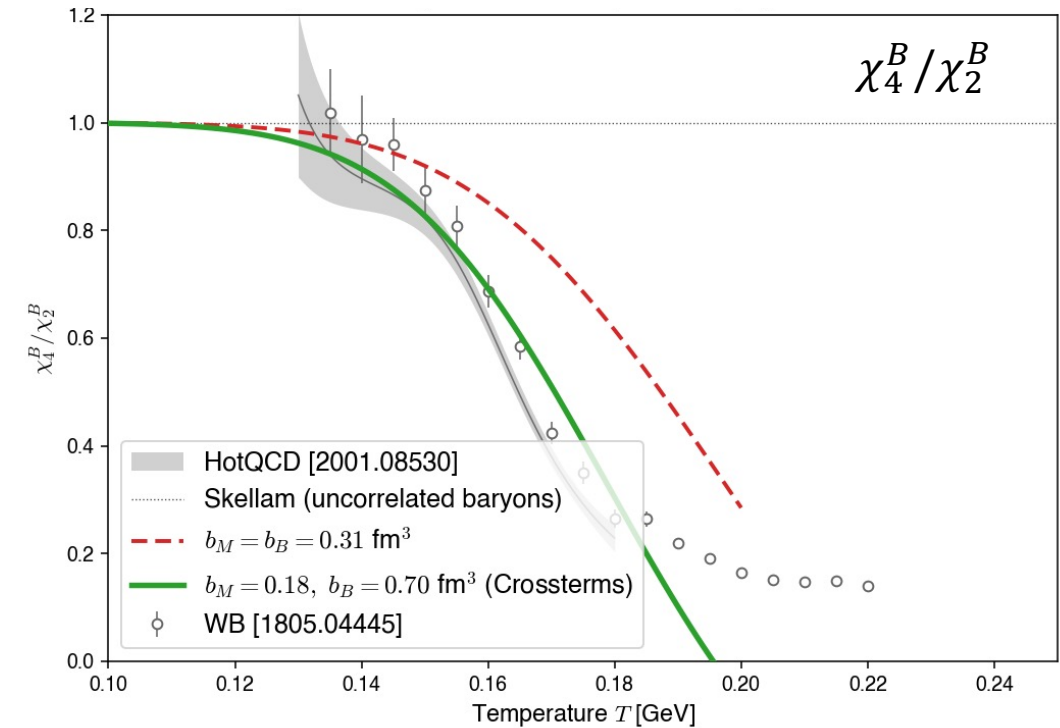
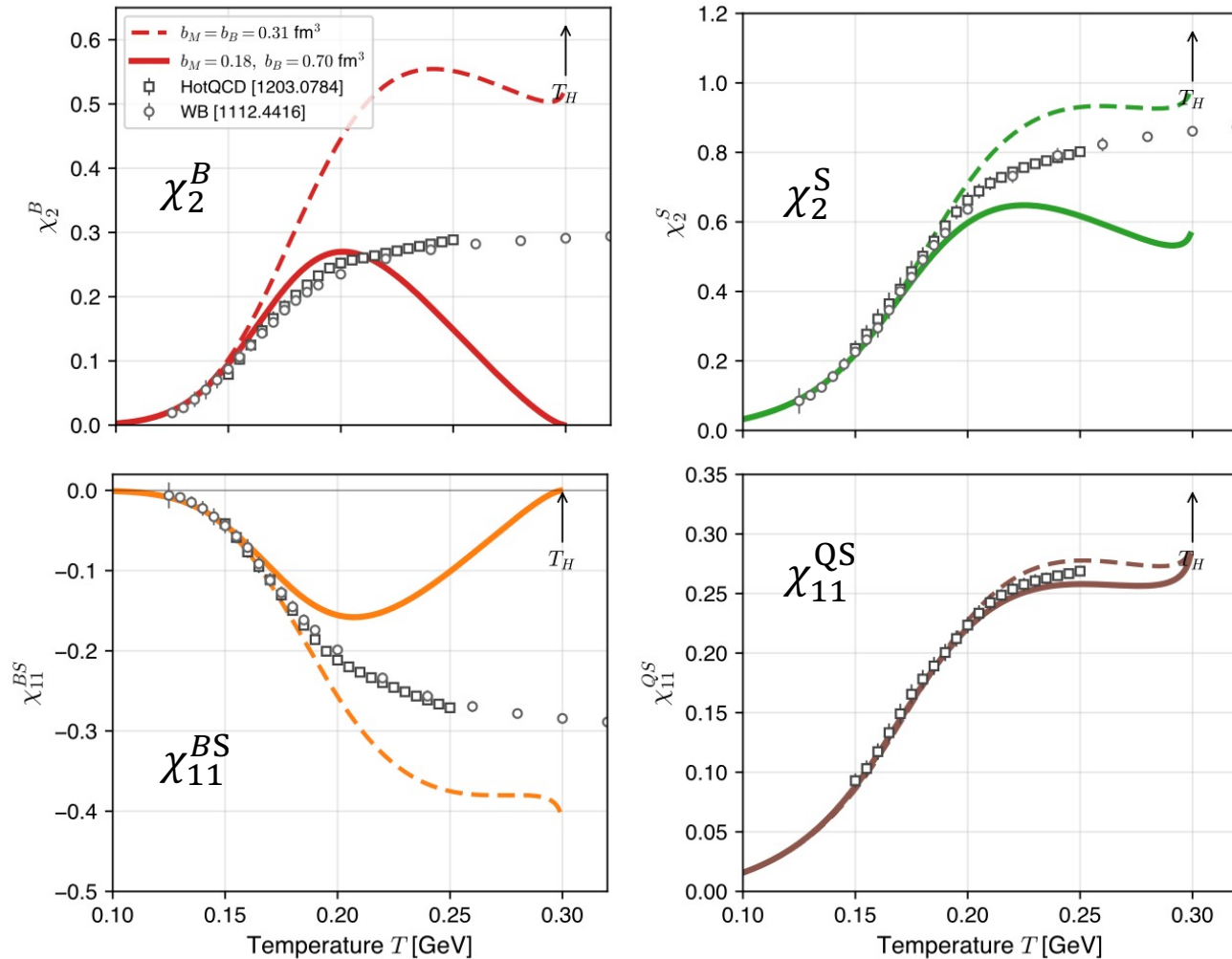
Thermodynamics of an interacting string gas



Uniform description of thermodynamics up to $T \approx 200 - 220$ MeV

Fujimoto, VV, to appear

Susceptibilities of an interacting string gas



Susceptibilities reveal a preference for a larger EV parameter for baryonic strings

Fujimoto, VV, to appear

QCD phase diagram with constant entropy contours

H. Shah, M. Hippert, J. Noronha, C. Ratti, VV, PRC 113, L012201 (2026) & 2601.08823

H. Shah, **T. Gyure**, et al., *to appear*

Critical point predictions from some years ago

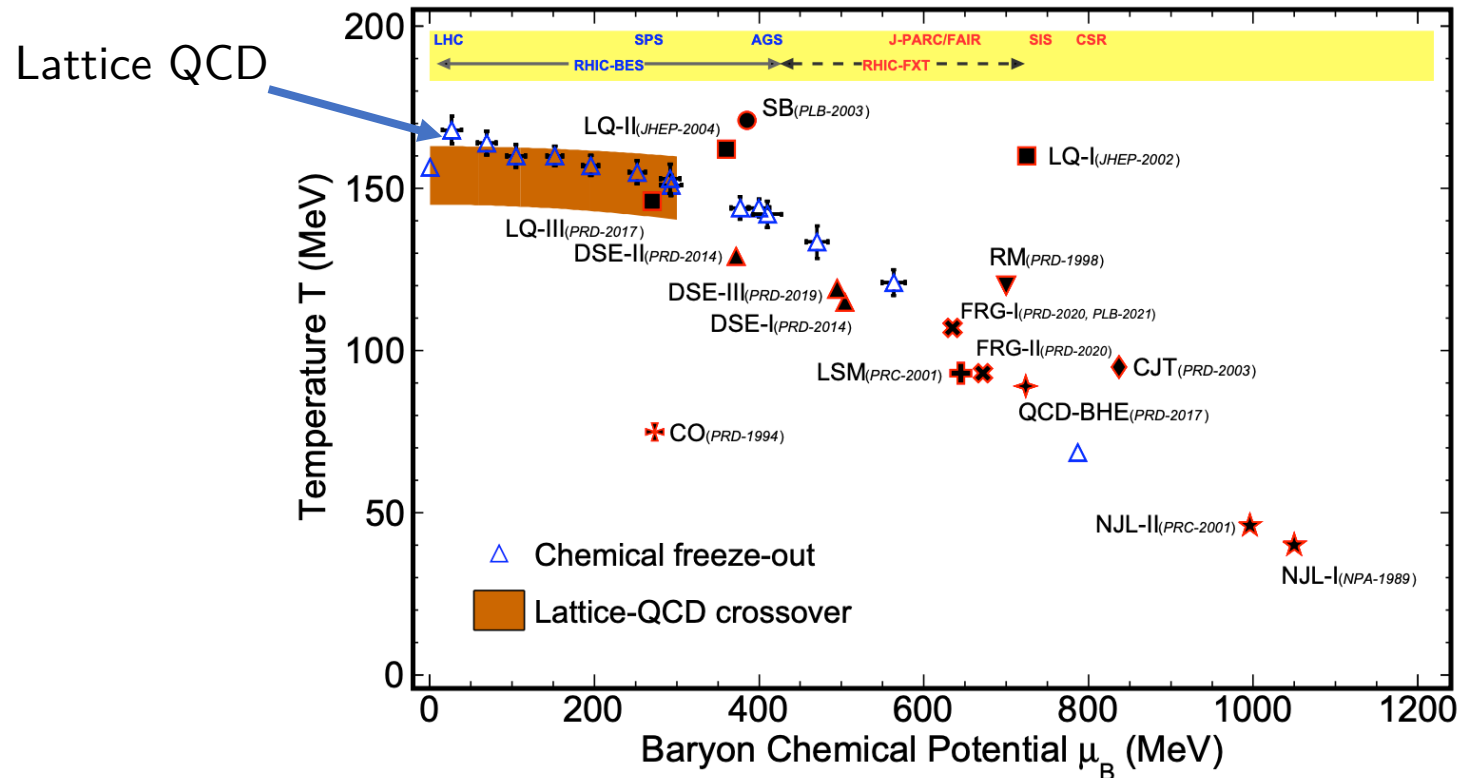


Figure adapted from A. Pandav, D. Mallick, B. Mohanty, Prog. Part. Nucl. Phys. 125 (2022)

- Including the possibility that the QCD critical point does not exist

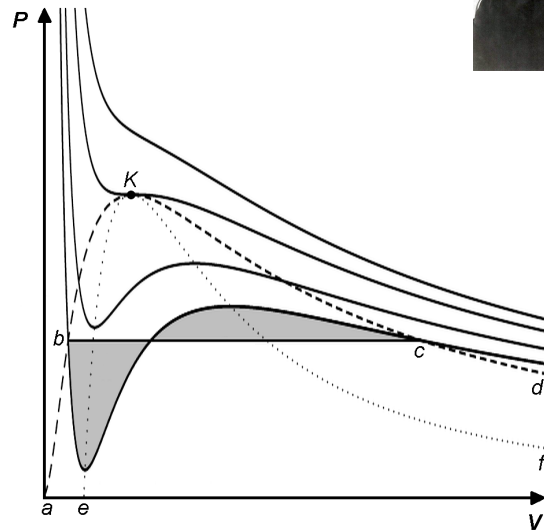
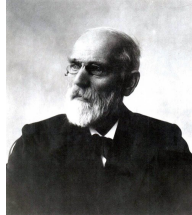
[de Forcrand, Philipsen, JHEP 01, 077 \(2007\)](#); [VV, Steinheimer, Philipsen, Stoecker, PRD 97, 114030 \(2018\)](#)

- Lattice QCD excludes the CP at $\mu_B < 450$ MeV on (one-sided) 2σ level

[Borsanyi et al., PRD 112, L111505 \(2025\)](#)

Extrapolating critical point from lattice QCD

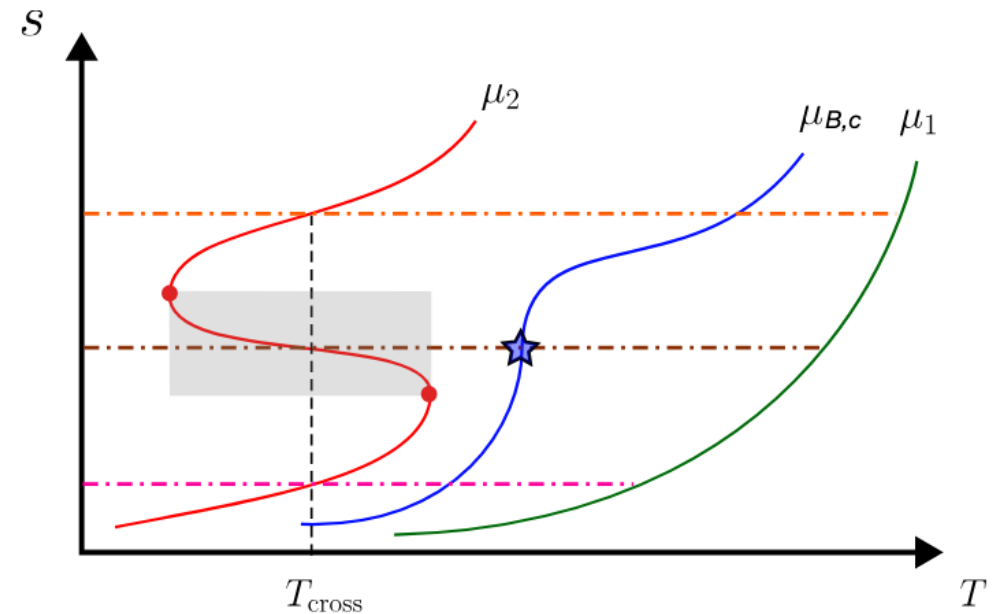
van der Waals (1873)



change of variables



Shah, Hippert, Noronha, Ratti, VV, PRC 113, L012201 (2026)



Critical Point:

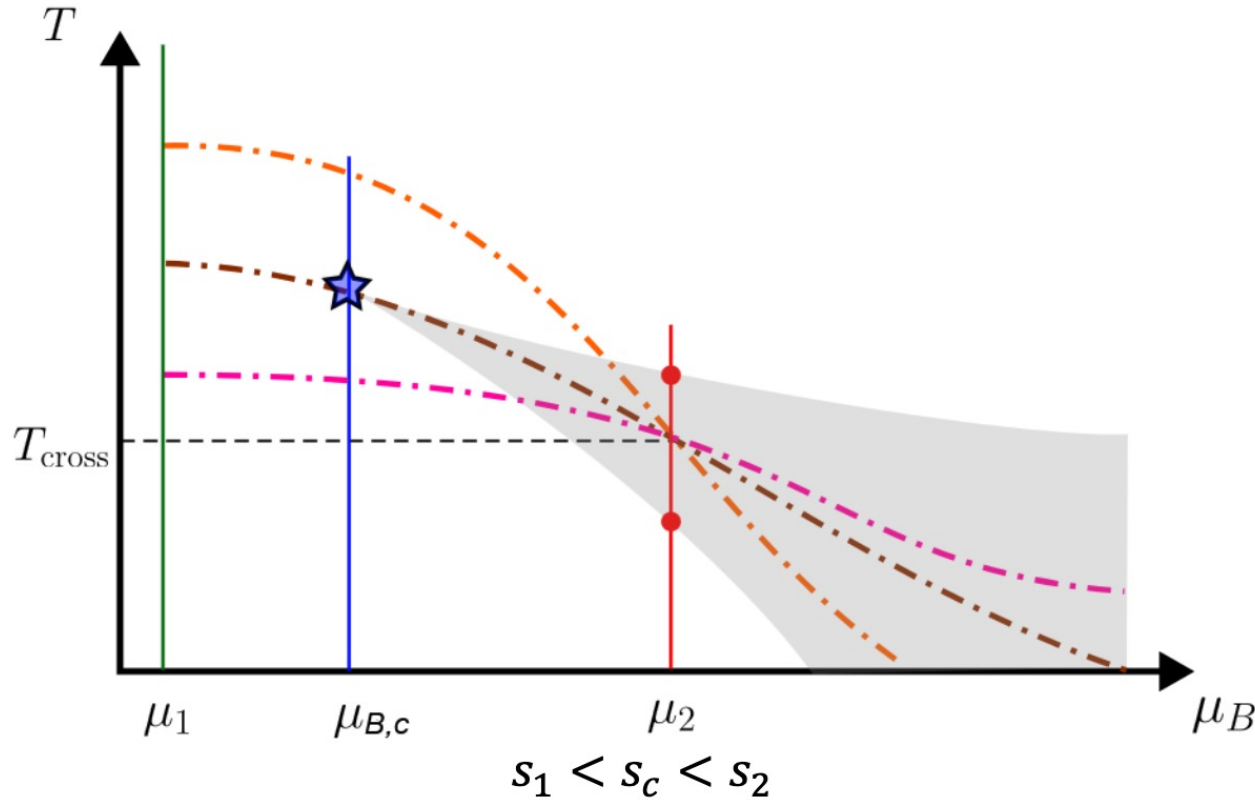
$$\left(\frac{\partial P}{\partial \rho_B}\right)_T = 0, \quad \left(\frac{\partial^2 P}{\partial \rho_B^2}\right)_T = 0.$$

$$\left(\frac{\partial T}{\partial s}\right)_{\mu_B} = 0, \quad \left(\frac{\partial^2 T}{\partial s^2}\right)_{\mu_B} = 0.$$

Strategy: Follow contours of constant entropy density away from $\mu_B = 0$ and look for crossings

Building the contours

Shah, Hippert, Noronha, Ratti, VV, PRC 113, L012201 (2026)



Define T -vs- μ_B line of constant entropy contour as Taylor series

$$T_s(\mu_B; T_0) = T_0 + \alpha_2(T_0) \frac{\mu_B^2}{2} \quad \text{order } O(\mu_B^2)$$

$$\alpha_2(T_0) = -\frac{2T_0\chi_2^B(T_0) + T_0^2\chi_2^{B'}(T_0)}{s'(T_0)}$$

- Allows multiple T_0 values to correspond to same T_s at finite μ_B
- In contrast to standard Taylor, permits the description of a first-order phase transition (mean-field-like behavior due to analyticity)

$$T_s(\mu_B; T_0) = T_0 + \alpha_2(T_0) \frac{\mu_B^2}{2}$$

expansion

$$\left(\frac{\partial T}{\partial s}\right)_{\mu_B} = 0 \quad \Rightarrow$$

$$\mu_{B,c} = \sqrt{-\frac{2}{\alpha_2'(T_{0,c})}}$$

spinodals

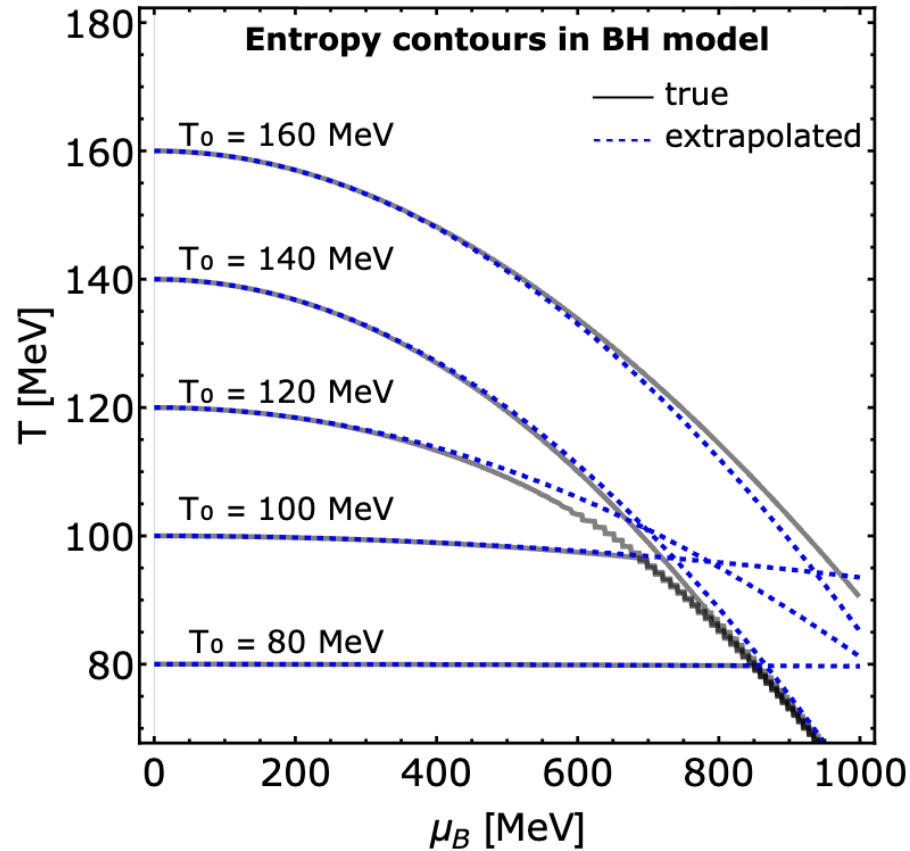
$$\left(\frac{\partial^2 T}{\partial s^2}\right)_{\mu_B} = 0 \quad \Rightarrow$$

$$\alpha_2''(T_{0,c}) = 0$$

critical point

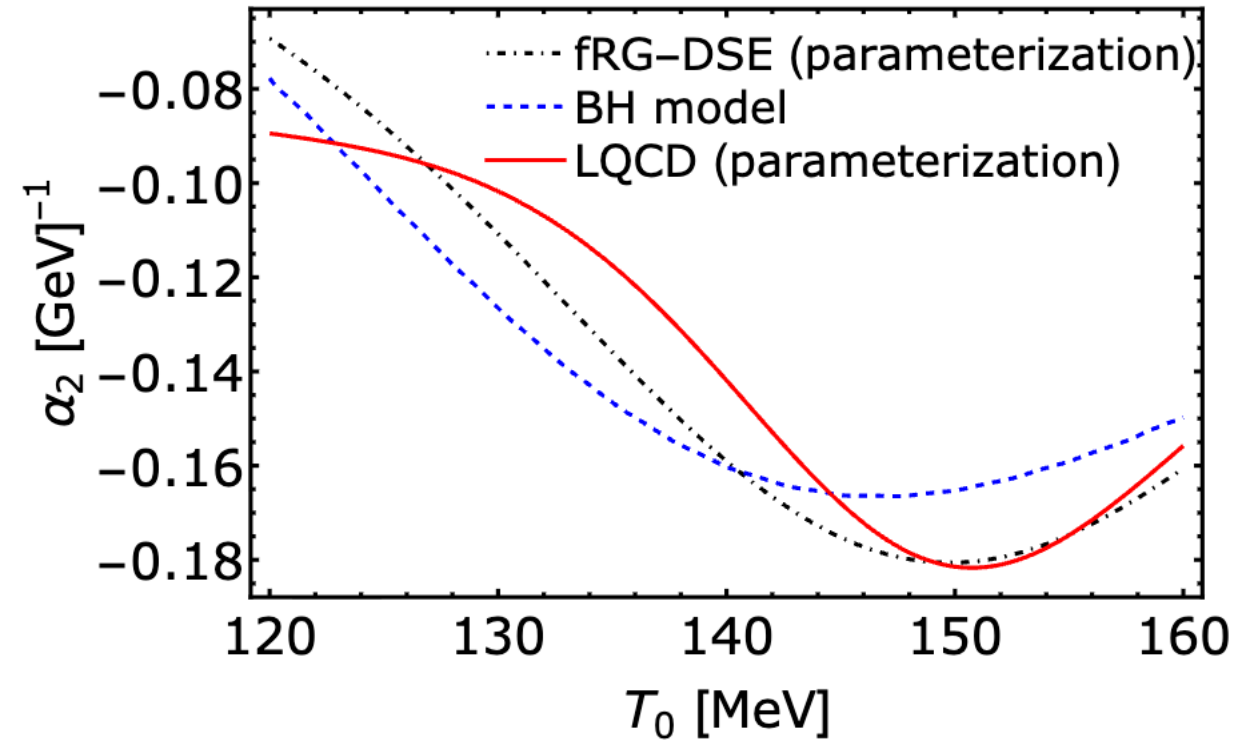
Testing with effective theories

Shah, Hippert, Noronha, Ratti, VV, arXiv:2601.08823



fRG-DSE: Lu, Gao, Liu, Pawłowski, PRD 113, 054019 (2026)

BH model: Hippert et al. PRD 110, 094006 (2024)



true CP

extrapolated CP

✓ fRG-DSE: (103,660) (108,629)

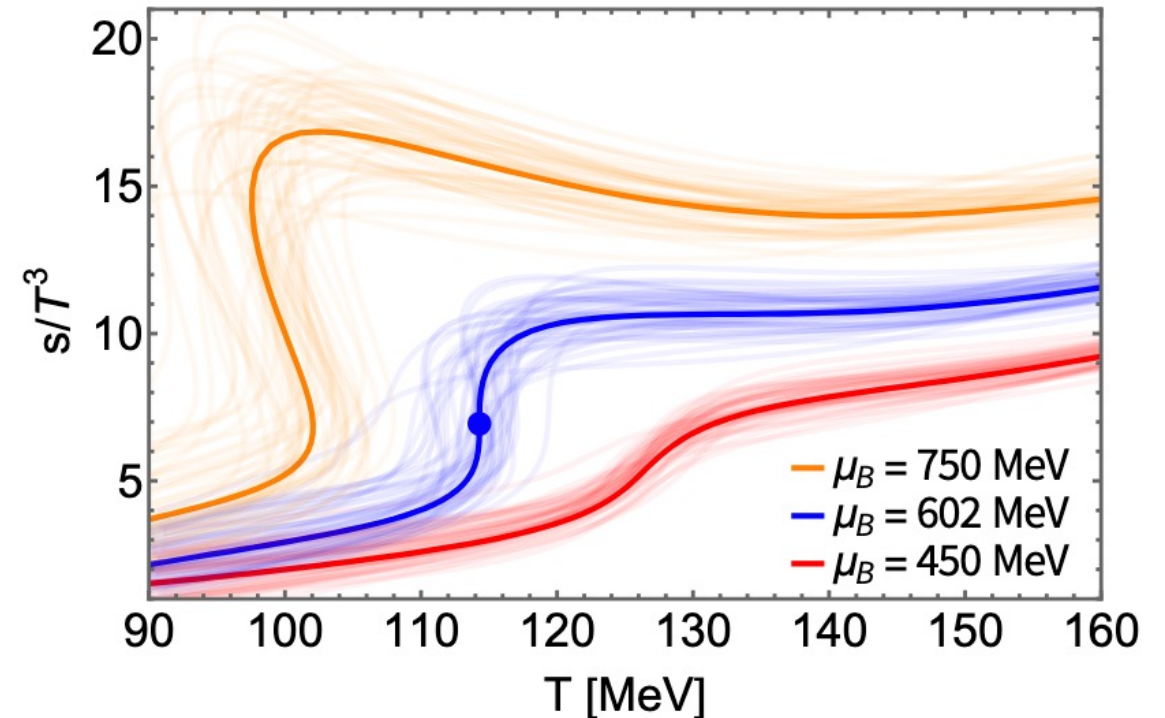
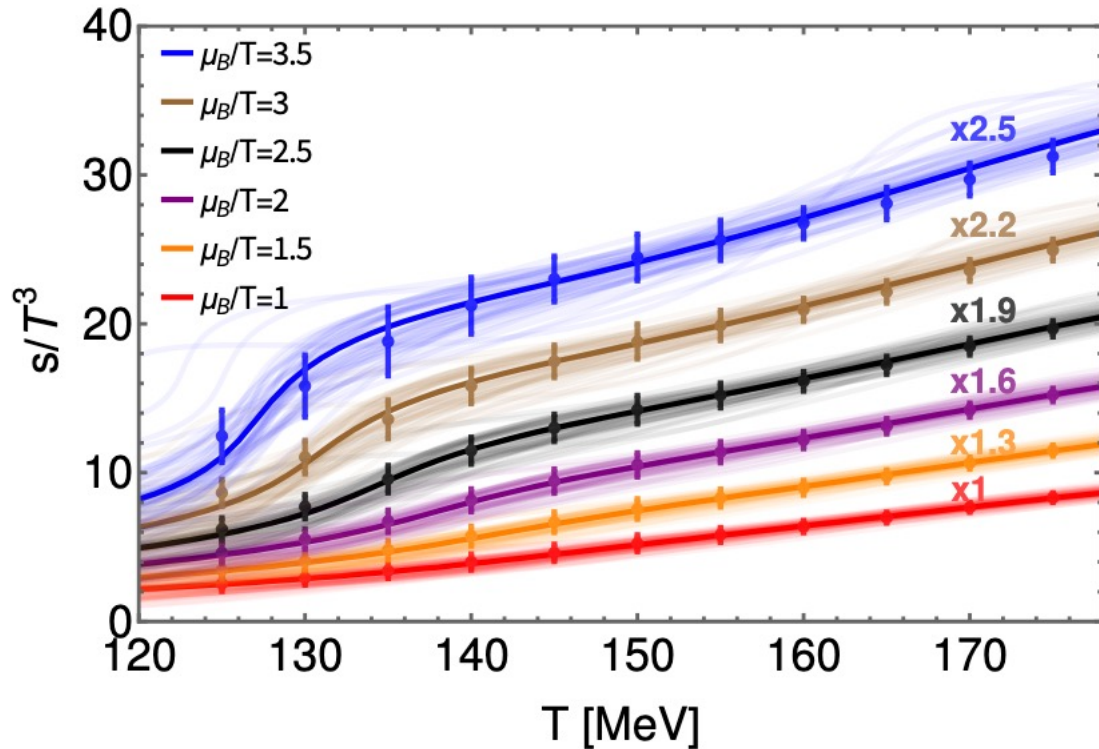
✓ BH model: (103,599) (104,637)

✓ **free QGP:** no true CP and no extrapolated CP

⚠ **HRG:** no true CP but fake extrapolated CP at $\sim(80,830)$ MeV

Entropy density at finite μ_B at $O(\mu_B^2)$ from lattice QCD

- Parametrized lattice QCD input from $\mu_B = 0$ [$s(T)$ and $\chi_2^B(T)$]

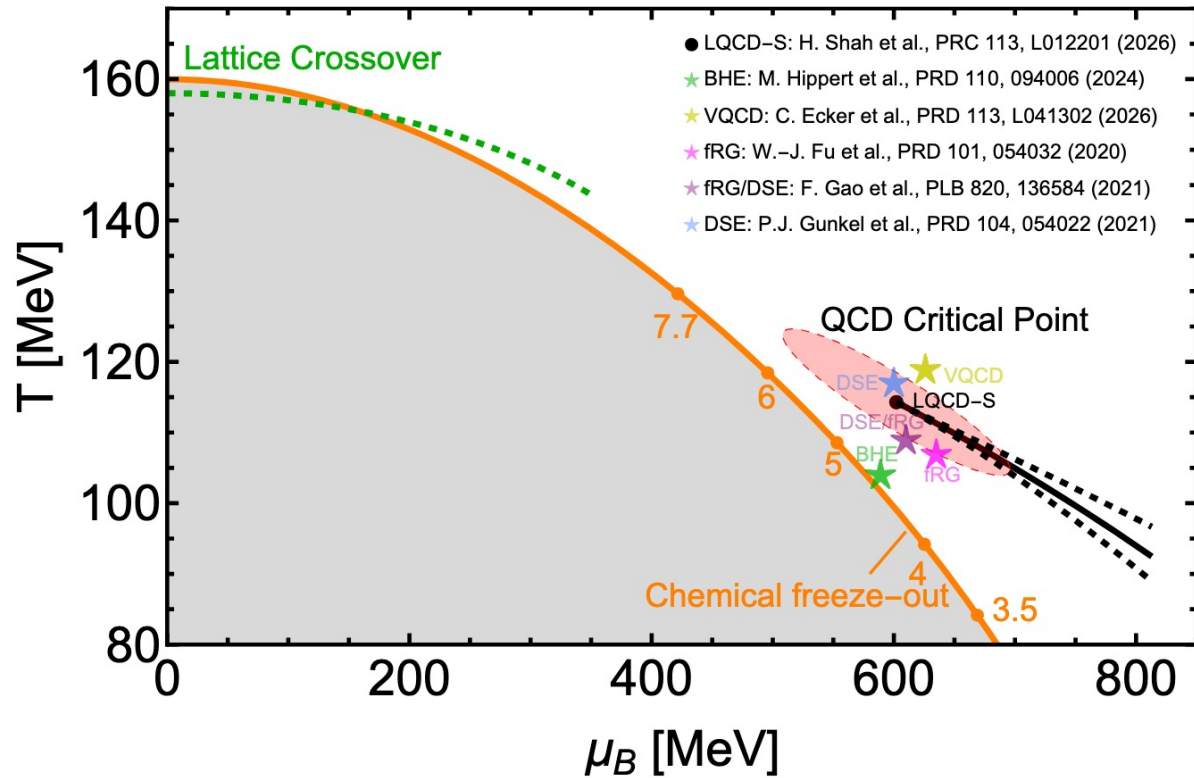


- Excellent agreement at low μ_B/T with available lattice QCD constraints from T' -expansion
- First-order phase transition emerges at $\mu_B > 600$ MeV

Borsanyi et al., PRL 126, 232001 (2021)

Shah, Hippert, Noronha, Ratti, VV, PRC 113, L012201 (2026)

Critical point estimates



Critical point estimate at $O(\mu_B^2)$:

$$T_c = 114 \pm 7 \text{ MeV}, \quad \mu_B = 602 \pm 62 \text{ MeV}$$

Estimates from recent literature:

YLE-1: D.A. Clarke et al. (Bielefeld-Parma), PRD 112, L091504 (2025)

YLE-2: G. Basar, PRC 110, 015203 (2024)

BHE: M. Hippert et al., PRD 110, 094006 (2024)

fRG: W.-J. Fu et al., PRD 101, 054032 (2020)

DSE/fRG: Gao, Pawłowski., PLB 820, 136584 (2021)

DSE: P.J. Gunkel et al., PRD 104, 052022 (2021)

VQCD: C. Ecker et al., PRD 113, L041302 (2026)

Optimist's view: Different estimates converge onto the same region because QCD CP is likely there

Pessimist's view: Different estimates converge onto the same region because it's the closest not yet ruled out by LQCD

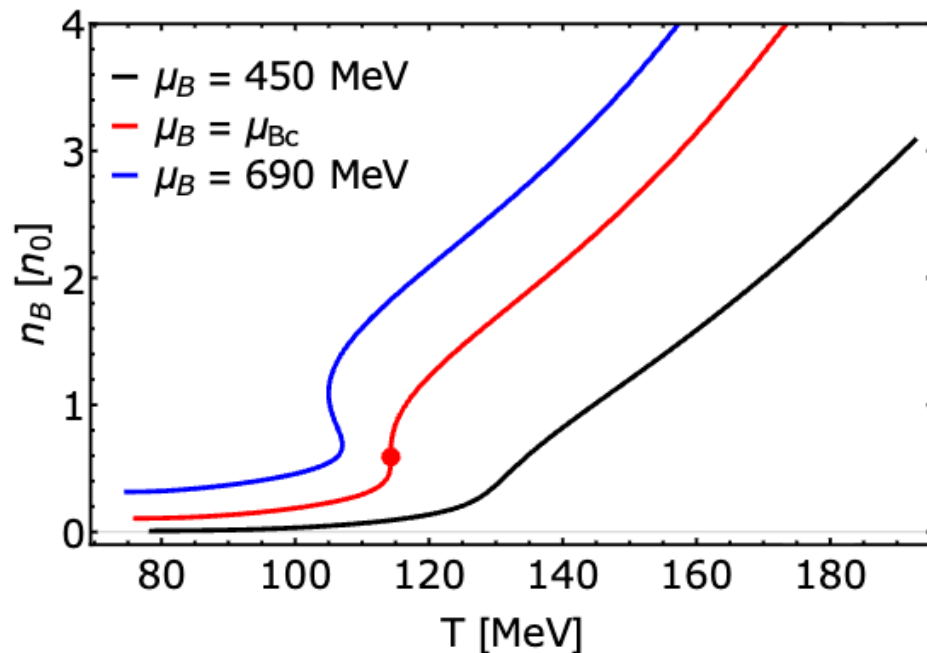
Equation of state

The method yields entropy density $s(T, \mu_B)$ across the (T, μ_B) phase diagram

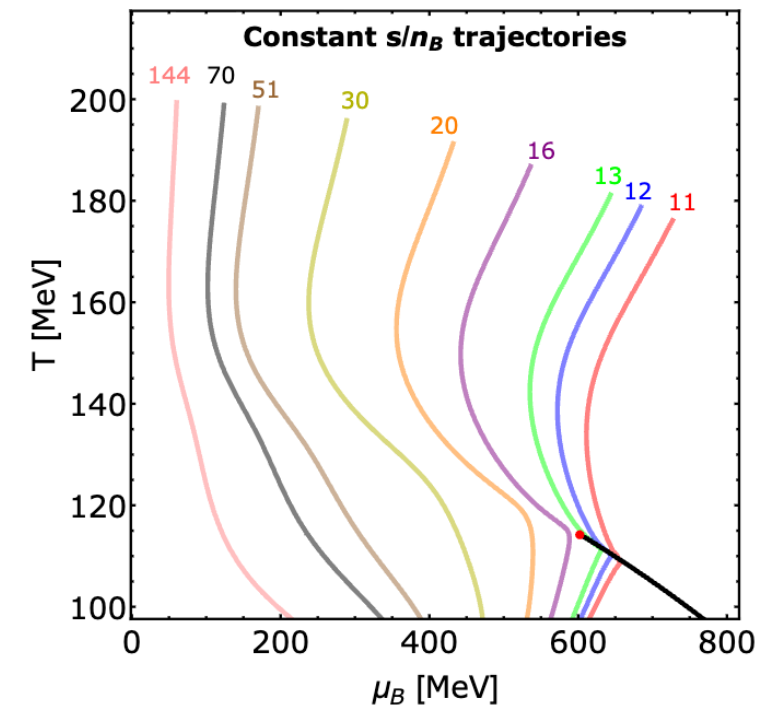
→ Integrate over T to obtain the pressure $P(T, \mu_B)$ and the rest of the thermodynamics

$$\int_{T_{\text{low}}}^T s(T', \mu_B) dT' = P(T, \mu_B) - P(T_{\text{low}}, \mu_B)$$

Integration constant at $T_{\text{low}} = 80$ MeV
matched to van der Waals HRG



Baryon density at the CP is $n_B \simeq 0.59 n_0 \simeq 0.09 \text{ fm}^{-3}$
 n_B in mixed phase region increases at lower T

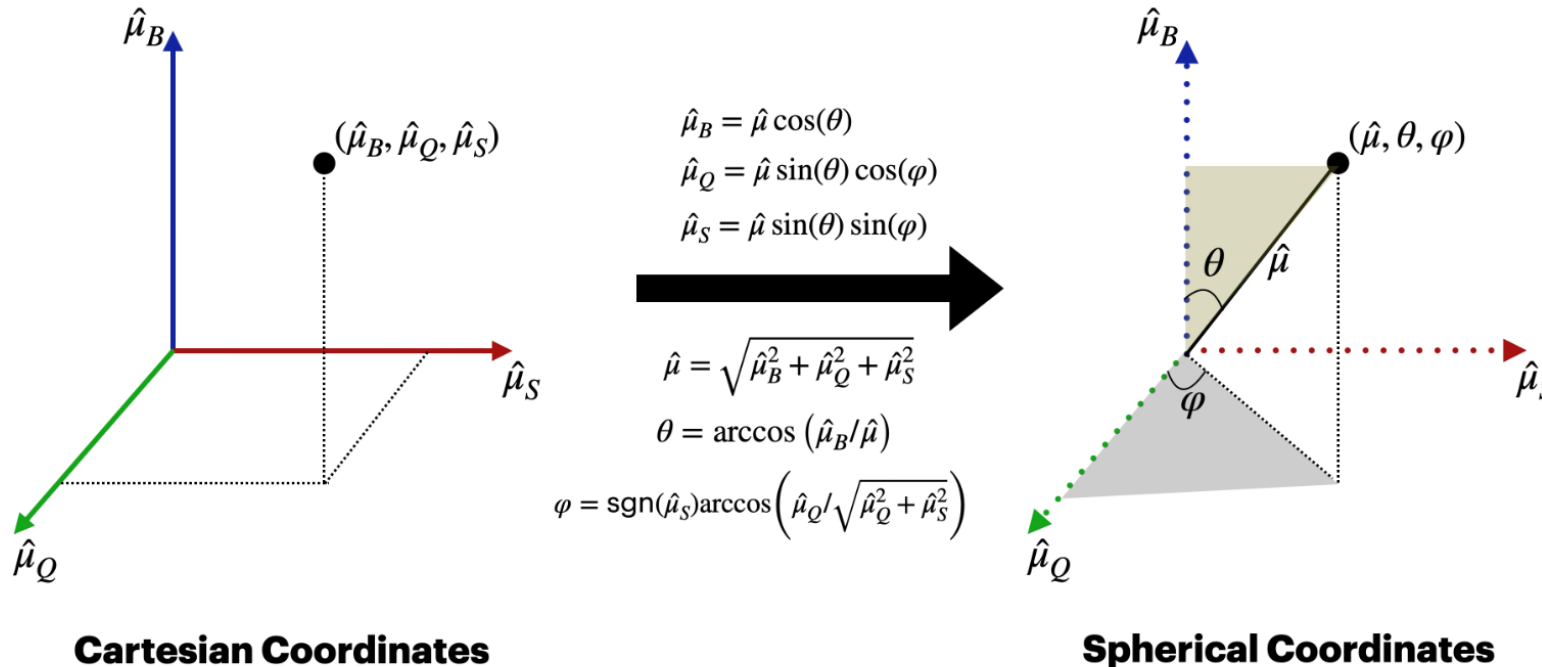


Critical point in 4D

So far we restricted the analysis to $\mu_S = \mu_Q = 0$ conditions

General case: fix the direction in 3D space of (μ_B, μ_Q, μ_S) using spherical coordinates

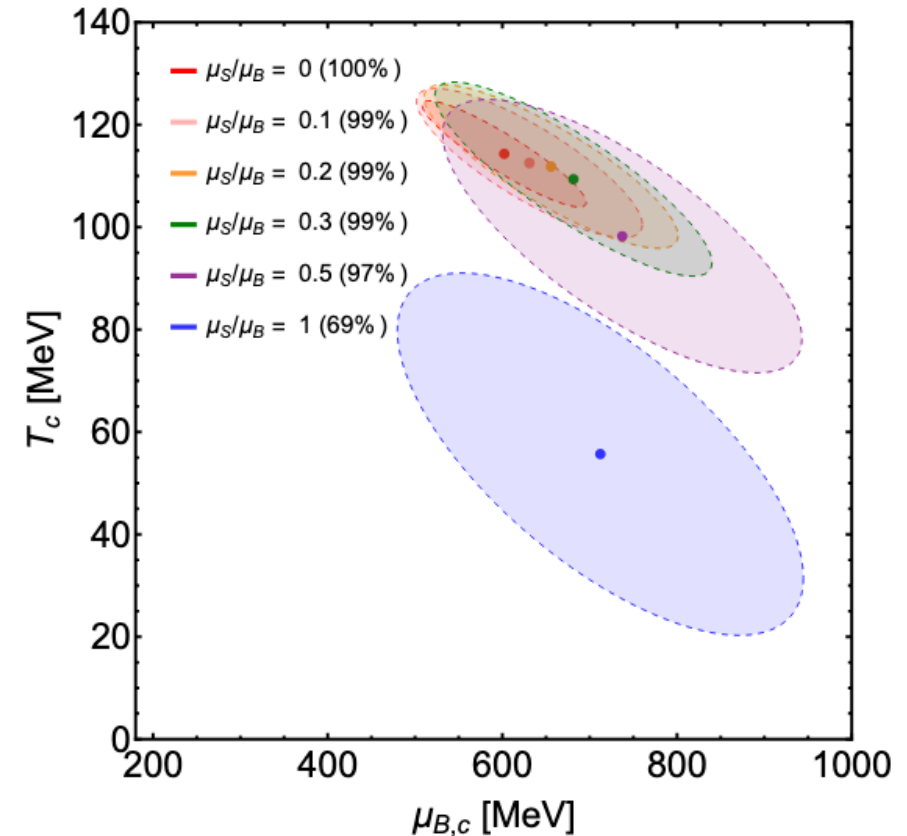
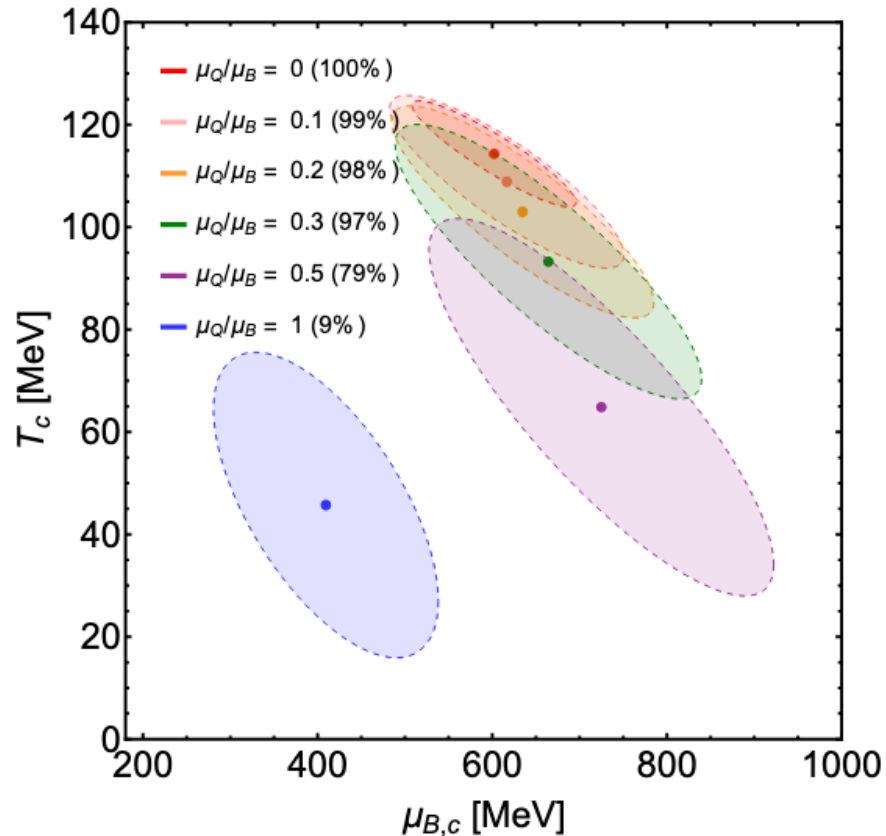
A. Abuali, S. Borsanyi, et al., PRD 112, 054502 (2025)



$$\chi_2^B(T) \longrightarrow X_2^{\theta, \varphi}(T) = c_\theta^2 \chi_2^B(T) + s_\theta^2 c_\varphi^2 \chi_2^Q(T) + s_\theta^2 s_\varphi^2 \chi_2^S(T) + 2c_\theta s_\theta c_\varphi \chi_{11}^{BQ}(T) + 2c_\theta s_\theta s_\varphi \chi_{11}^{BS}(T) + 2s_\theta^2 c_\varphi s_\varphi \chi_{11}^{QS}(T)$$

Using the lattice input for susceptibilities from [A. Abuali, et al., PRD 112, 054502 (2025)]

The brackets denote the percentage of MC samples giving the critical point.



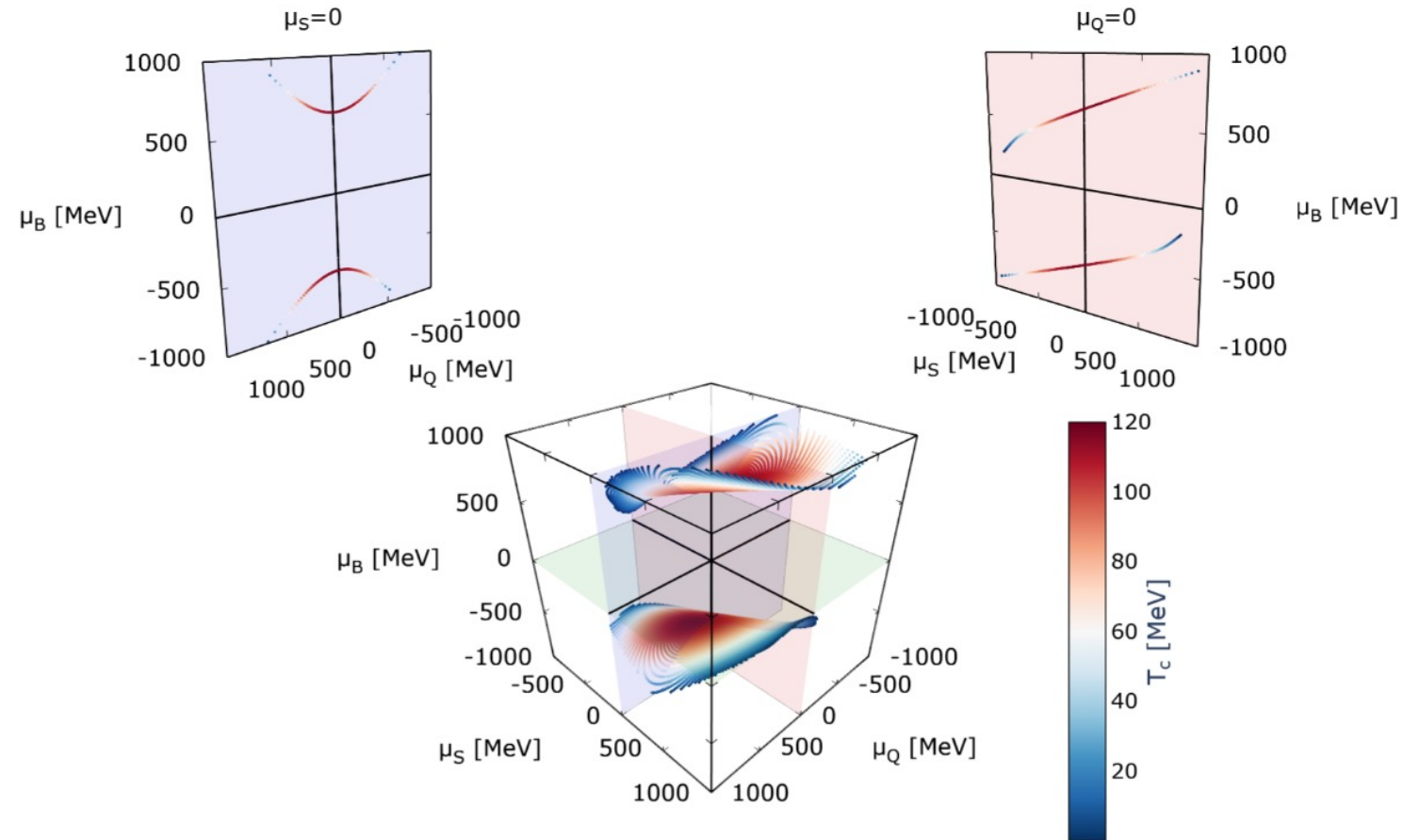
Strangeness neutrality ($\mu_S \approx \mu_B/3$): $\mu_{B,c}$ shifts to the right.

From $(T, \mu_B) = (114, 602)$ MeV at $\mu_S = 0$ to $(T, \mu_B) = (109, 683)$ MeV at $\mu_S = \mu_B/3$

Similar effect of an increased $\mu_{B,c}$ in recent functional QCD analysis [W.J. Fu et al., 2603.13455]

Critical point in 4D: full result

H. Shah, T. Gyure, F. Di Clemente, et al., to appear



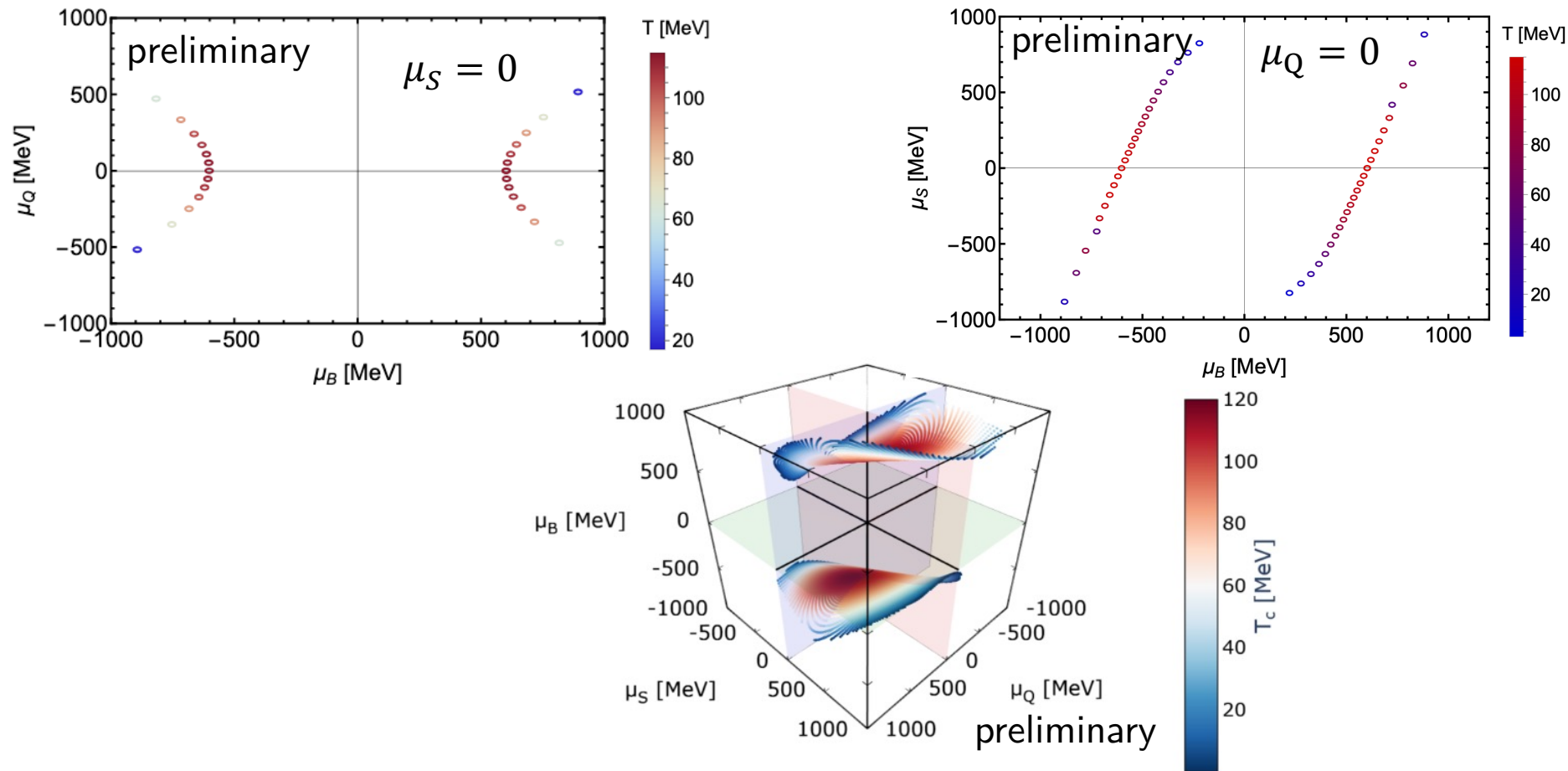
Pure μ_Q direction: No evidence for a CP. Consistent with lattice QCD at finite isospin

Caveats: Pion and kaon condensation

Applications: Heavy-ion collisions, cosmic trajectories with large lepton asymmetries,...

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Fluctuations in heavy-ion collisions

Critical point and heavy-ion collisions

Control parameters

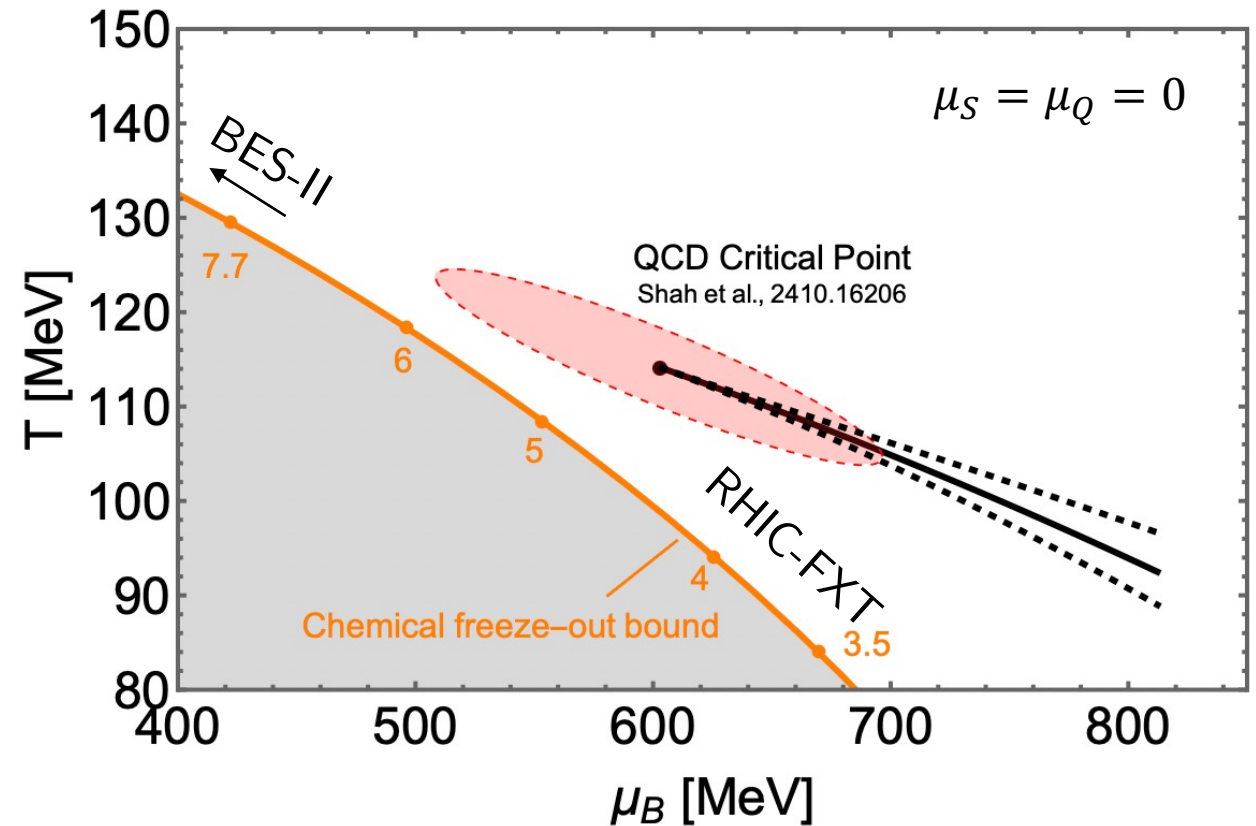
- Collision energy $\sqrt{s_{NN}} = 2.4 - 5020$ GeV
 - Scan the QCD phase diagram
- Size of the collision region
 - Expect stronger signal in larger systems

Measurements

- Final hadron abundances and momentum distributions **event-by-event**

Chemical freeze-out curve and CP

- Sets the **lower bound** on the temperature of the CP [\[Lysenko, Poberezhnyuk, Gorenstein, VV, PRC 111, 054903 \(2025\)\]](#)
- **Caveats:** strangeness neutrality ($\mu_S \neq 0$), uncertainty in the freeze-out curve
- If the $\mu_B \approx 600$ MeV predictions are true, the CP may be close to freeze-out line at $\sqrt{s_{NN}} \sim 3.5 - 5$ GeV



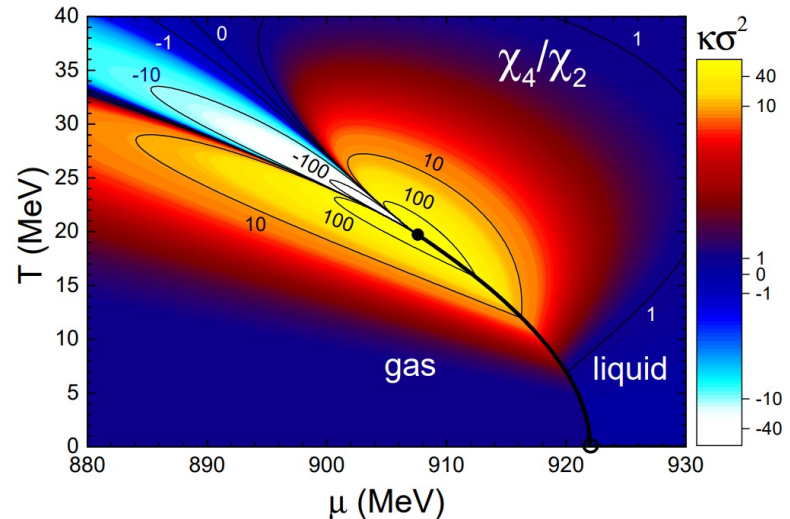
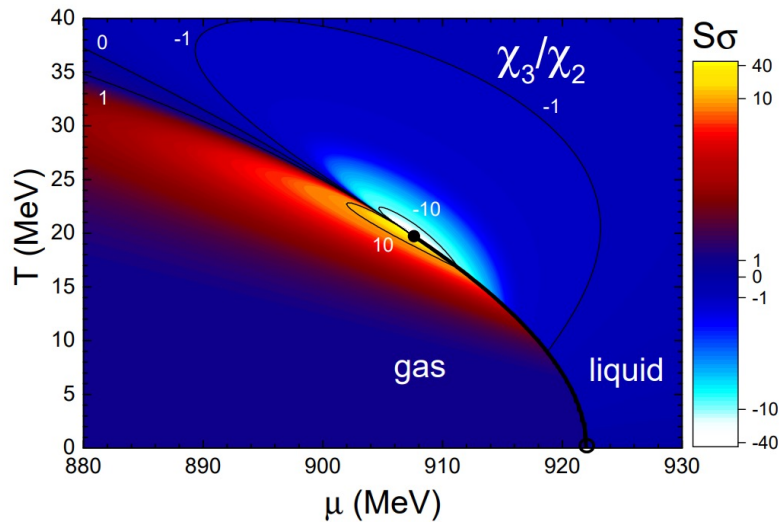
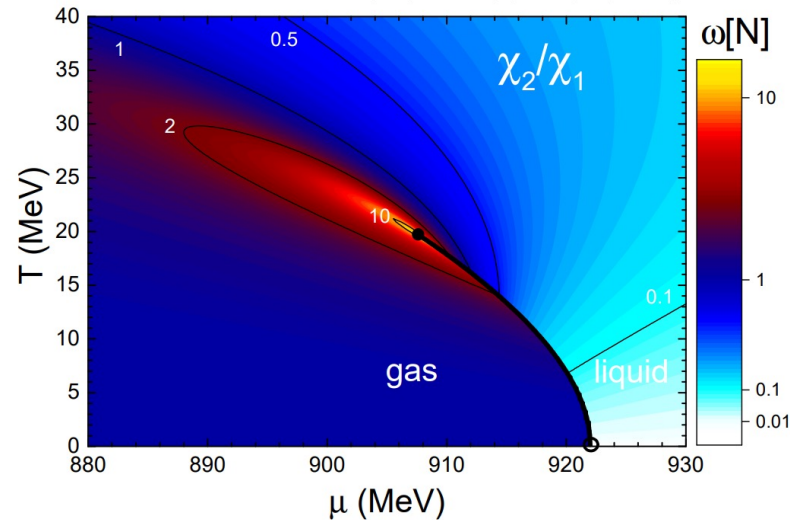
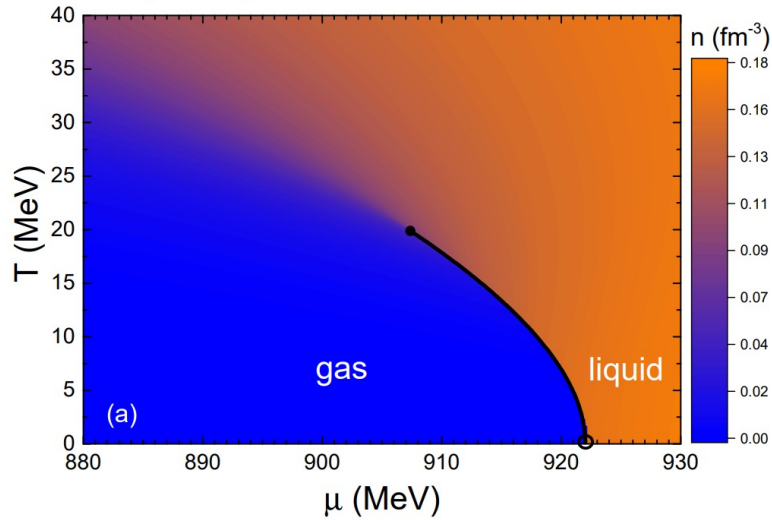
Example: (Nuclear) Liquid-gas transition

(QCD) critical point: large correlation length and fluctuations

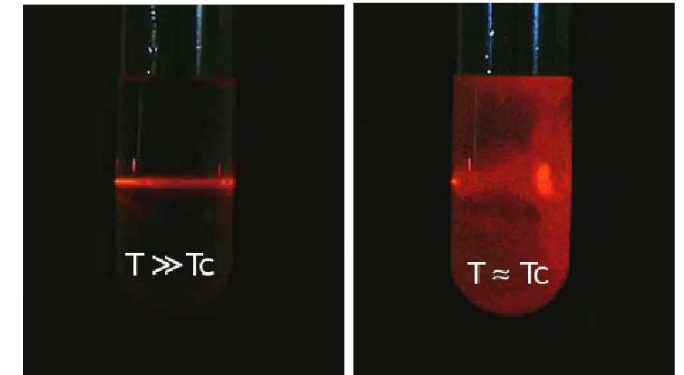
$$\kappa_2 \sim \xi^2, \quad \kappa_3 \sim \xi^{4.5}, \quad \kappa_4 \sim \xi^7$$

$$\xi \rightarrow \infty$$

M. Stephanov, PRL '09, '11



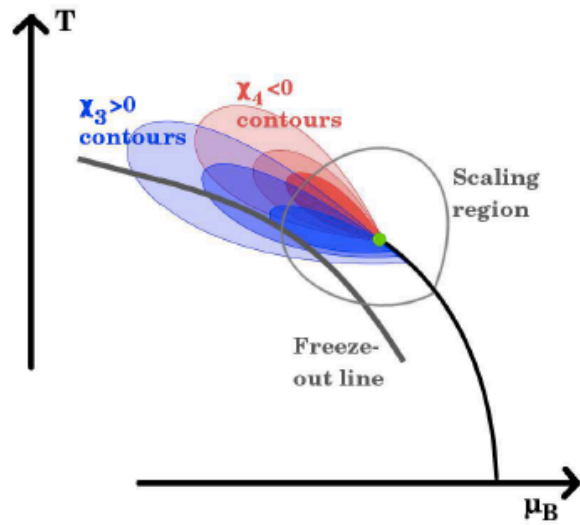
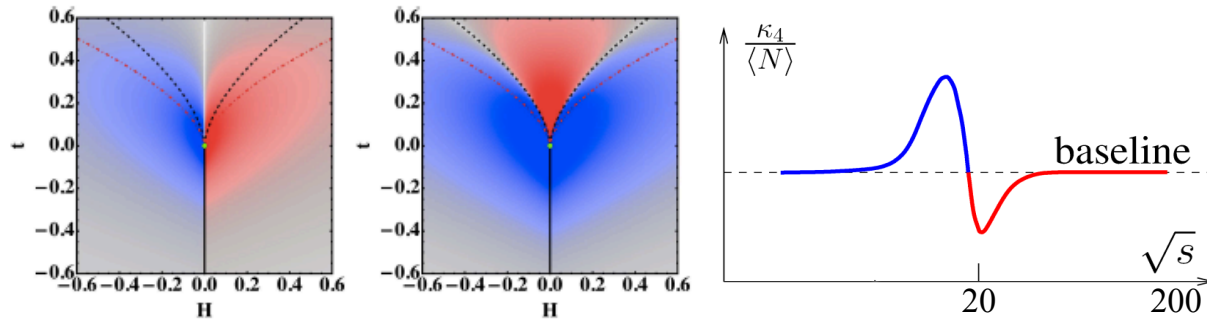
Critical opalescence



$$\langle N^2 \rangle - \langle N \rangle^2 \sim \langle N \rangle \sim 10^{23}$$

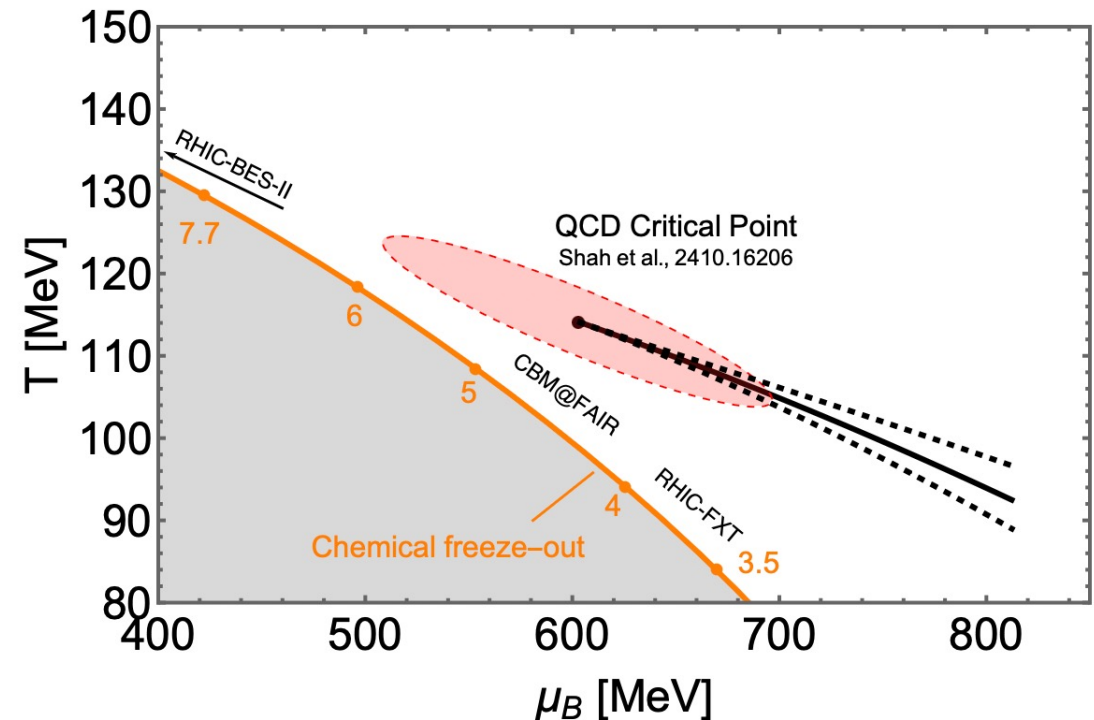
in equilibrium

Expectation from Calculations



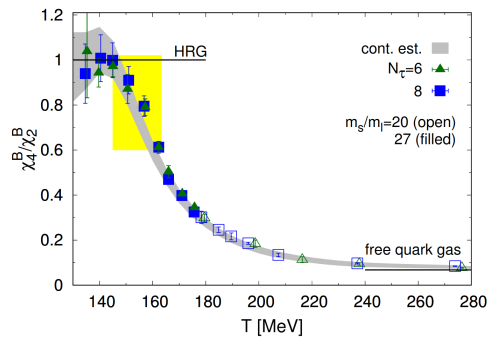
Characteristic “Oscillating pattern” is expected for the QCD critical point but *the exact shape depends on the location of freeze-out with respect to the location of CP*

- M. Stephanov, *PRL* **107**, 052301(2011)
- V. Skokov, Quark Matter 2012
- J.W. Chen, J. Deng, H. Kohyama, arXiv: 1603.05198, Phys. Rev. **D93** (2016) 034037

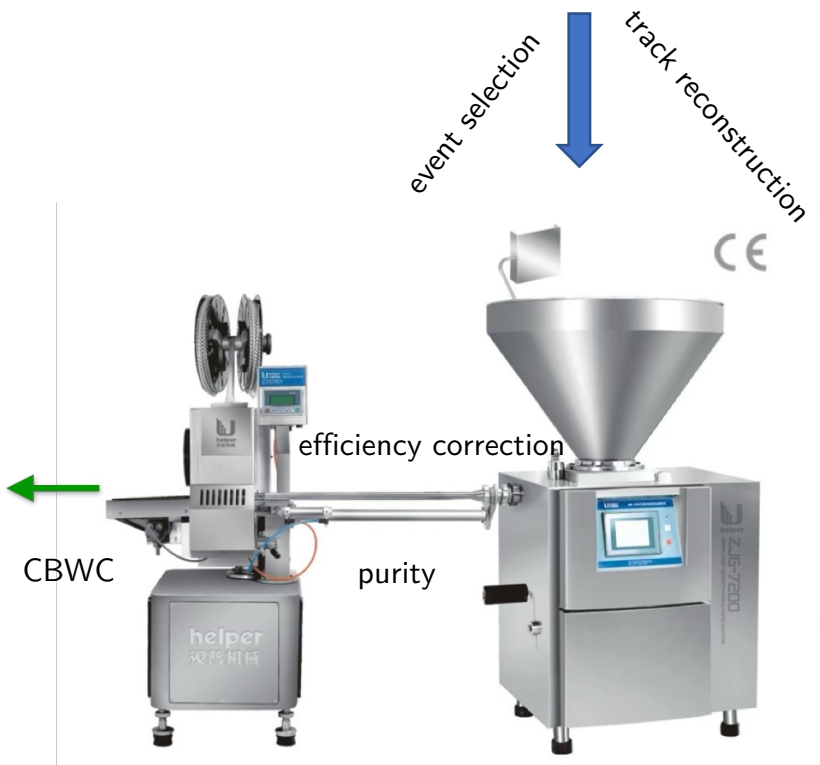
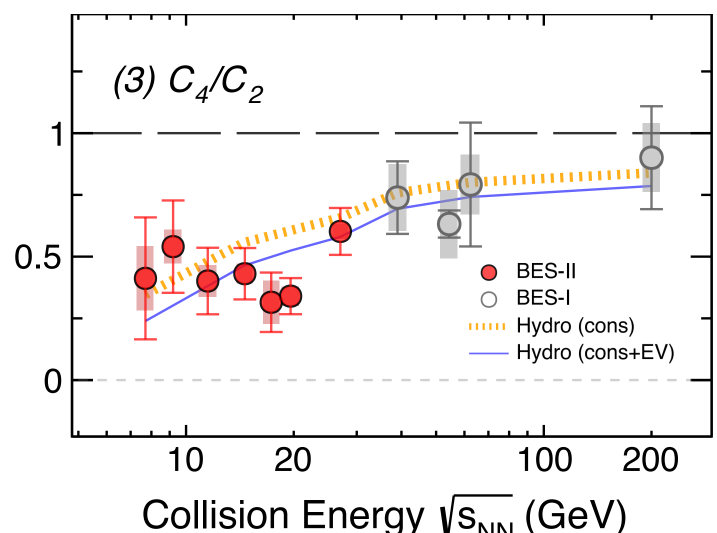
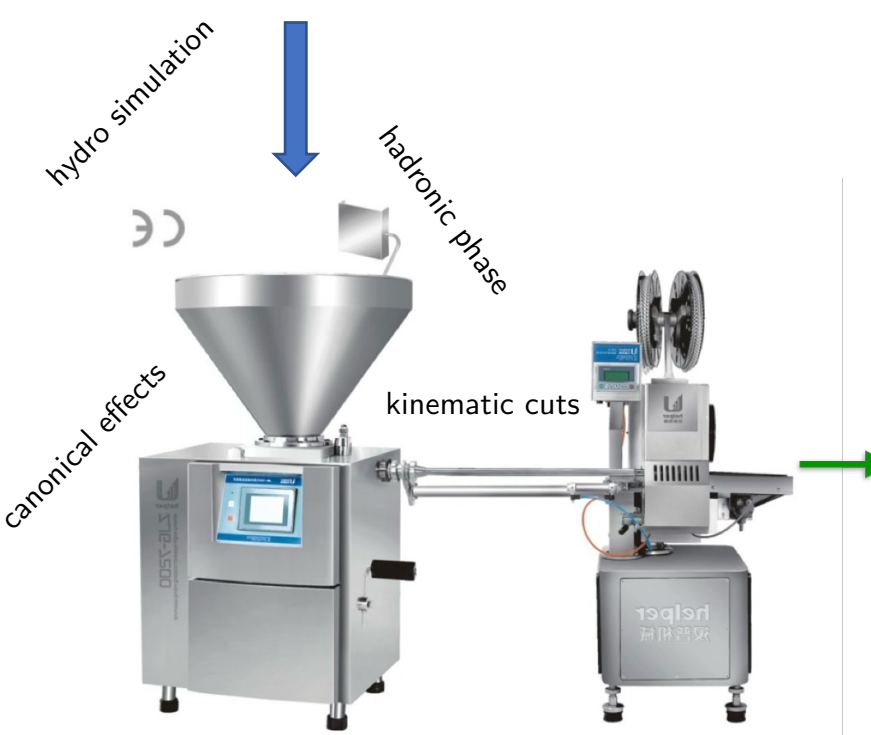
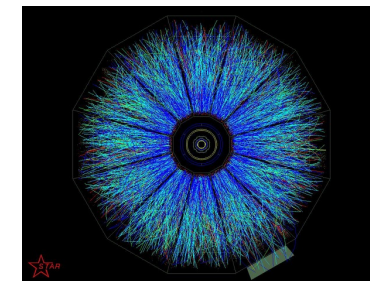


Theory vs experiment

guidance from theory (e.g. lattice)



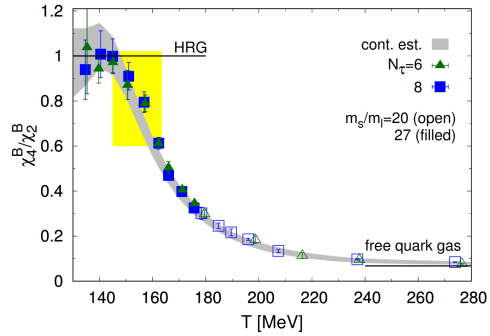
experiment (the real thing)



Theory vs experiment

guidance from theory (e.g. lattice)

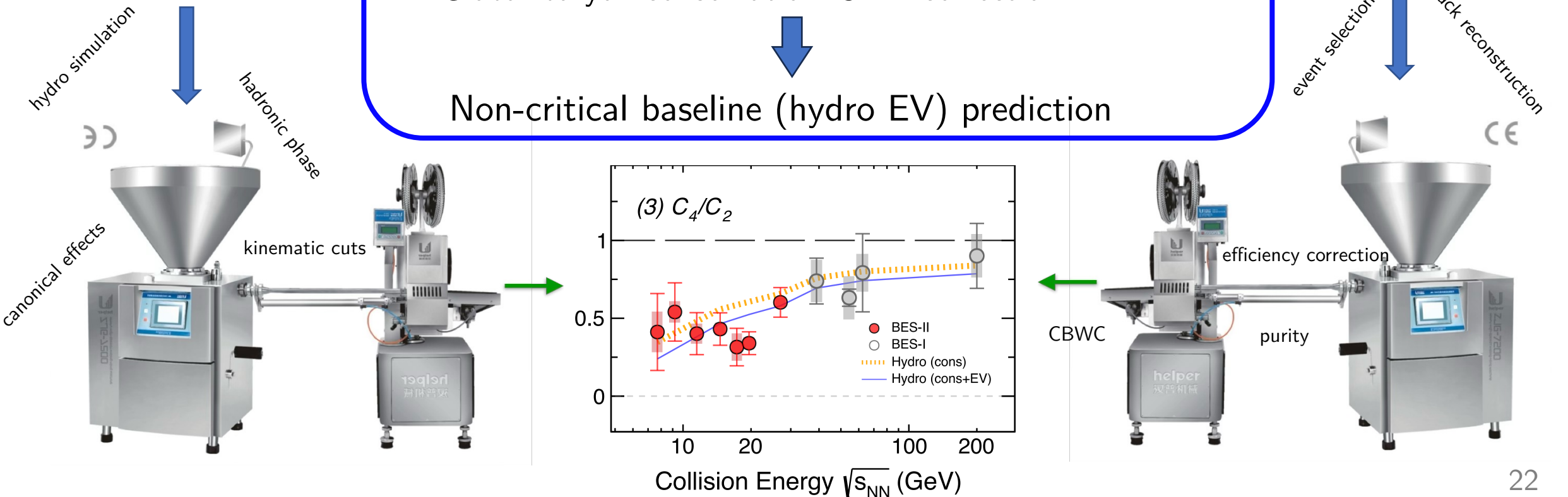
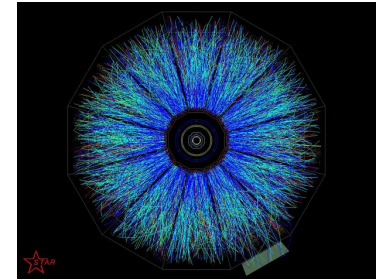
experiment (the real thing)



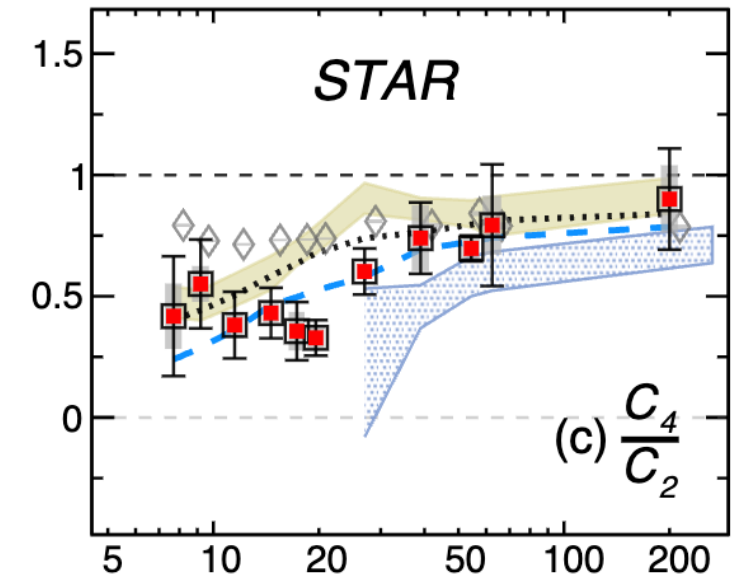
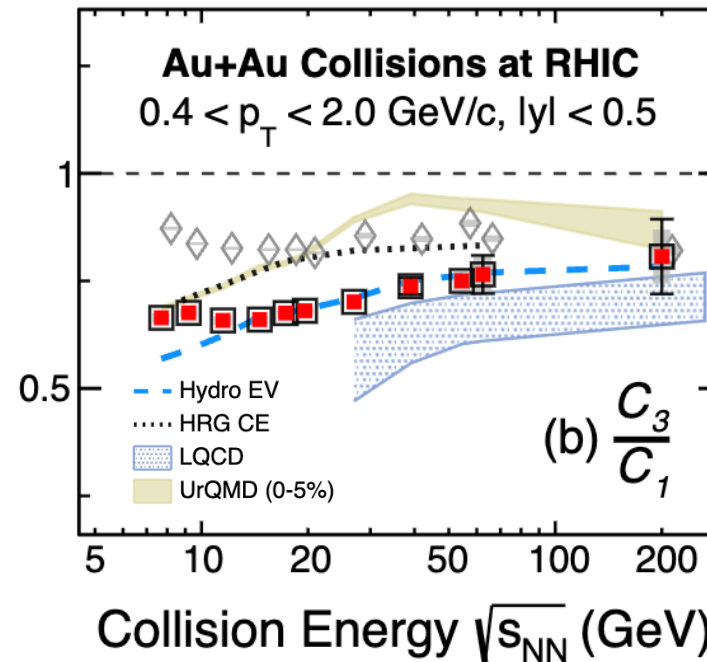
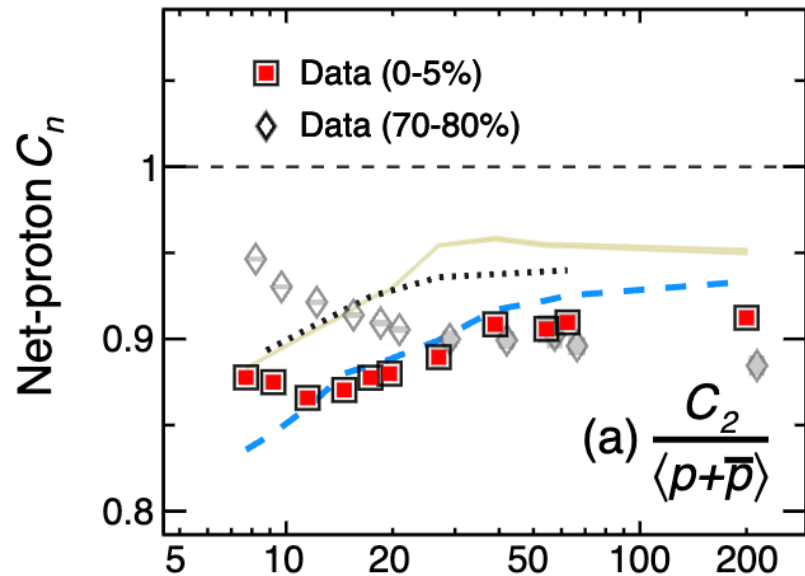
This was done in [VV, V. Koch, C. Shen, Phys. Rev. C 105, 014904 (2022)]

- Full hydro simulation
- Lattice QCD-like baryon susceptibilities (interacting HRG)
- Experimental kinematic cuts
- Global baryon conservation: SAM correction

Non-critical baseline (hydro EV) prediction



Net-proton cumulant ratios



STAR, PRL 135, 142301 (2025)

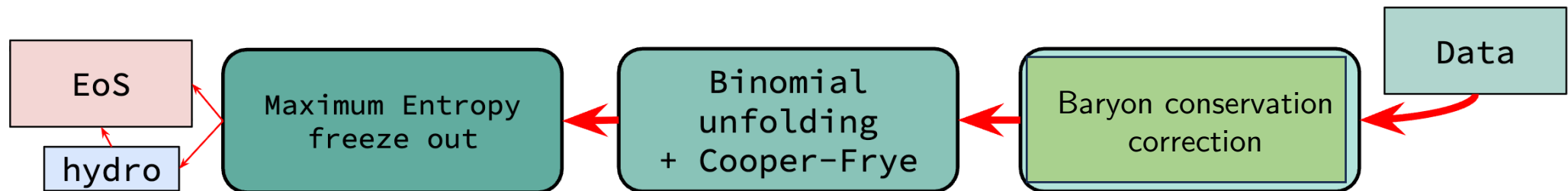
Hydro EV: [VV, V. Koch, C. Shen, Phys. Rev. C 105, 014904 \(2022\)](#)

Agreement with the baseline above $\sqrt{s_{NN}} \sim 10 - 20 \text{ GeV}$

Deviations at lower energies. More pronounced in factorial cumulants.

Inference of baryon number susceptibilities from data

G. Pihan, R. Poberezhniuk, VV, to appear



Hydro EV [VV, V. Koch, C. Shen, Phys. Rev. C 105, 014904 (2022)]

- Particlization at a constant energy density $\sim 0.26 \text{ GeV}/\text{fm}^3$
- Canonical (baryon conservation) statistical hadronization with **excluded volume repulsion**
- **Drawback:** Cannot incorporate arbitrary EoS, in particular one with a CP

Maximum entropy freeze-out [M. Pradeep et al., PRL 130, 162301 (2023); arXiv:2508.19237]

- Incorporate single hydrodynamic mode – **baryon density fluctuations**
- Maximum entropy method defines local **baryon/antibaryon joint susceptibilities** G. Pihan et al., to appear

$$\hat{\chi}_{nm}^{+-}(x) = \delta_{m0}\delta_{n1}\bar{\chi}_1^+ + \delta_{n0}\delta_{m1}\bar{\chi}_1^- + (-1)^m \underbrace{\hat{\Delta}\chi_{n+m}^B(x)}_{\text{difference to HRG}} \frac{\bar{\chi}_1^+(x)^n \bar{\chi}_1^-(x)^m}{(\bar{\chi}_1^+(x) + \bar{\chi}_1^-(x))^{n+m}} \quad \bar{\chi}_n^B(x) = \bar{\chi}_1^+(x) + (-1)^n \bar{\chi}_1^-(x)$$

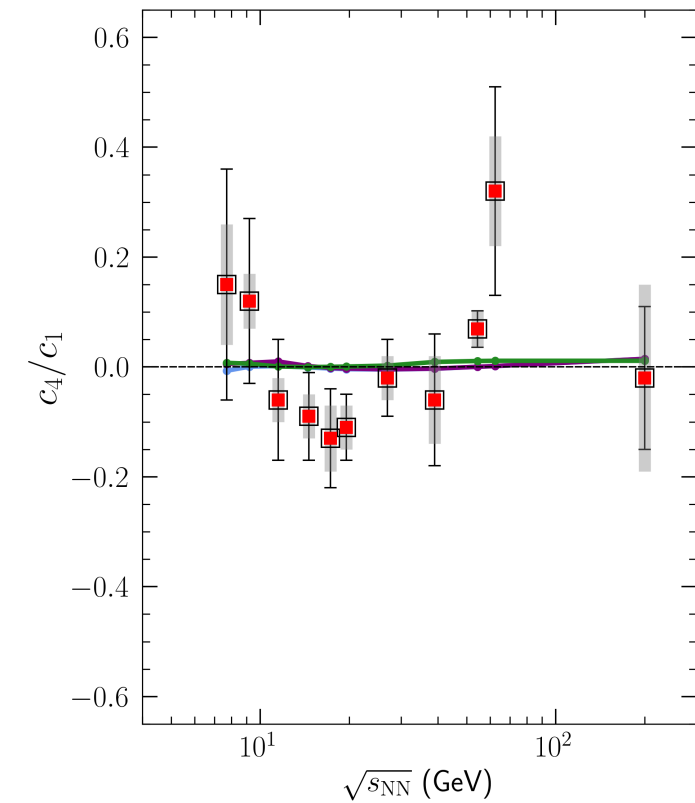
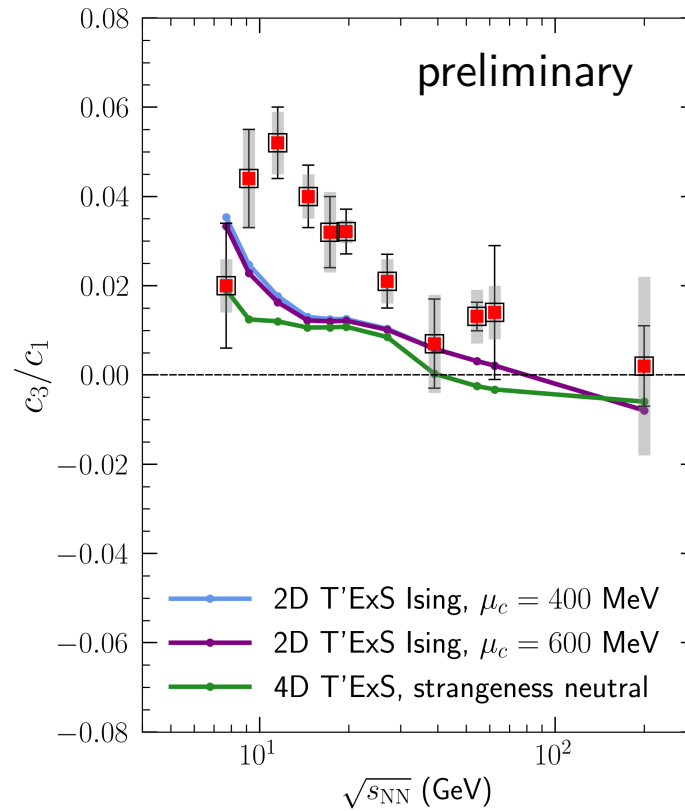
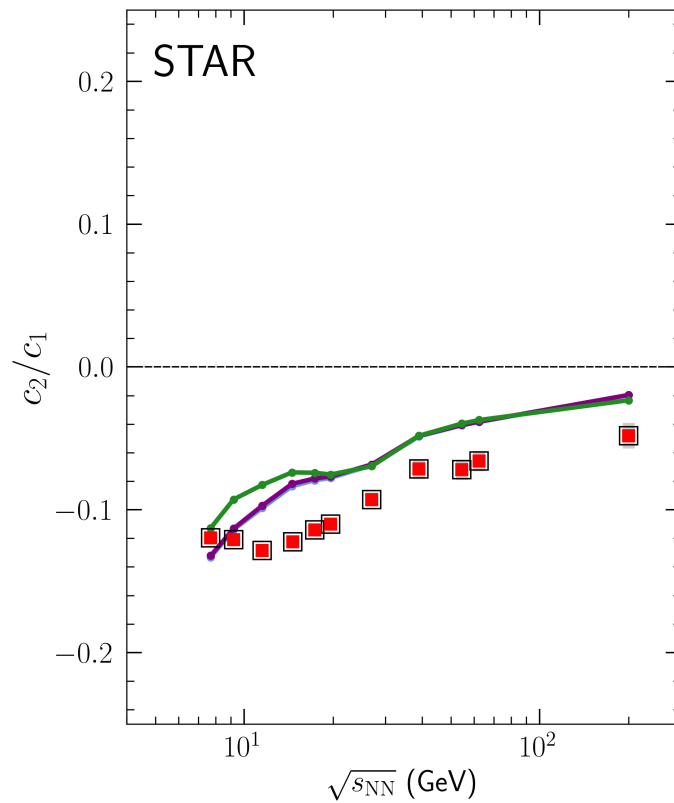
+ is baryon, – is antibaryon

HRG (Poisson) baseline

- Use the χ_n^B input from **arbitrary EoS** (LQCD expansion, s-contours, functional QCD, ...)

More details in talks of G. Pihan and J. Karthein next week at “QCD Critical Point and Hydrodynamic Evolution”

Pihan, Kahangirwe, Shah, et al., in preparation



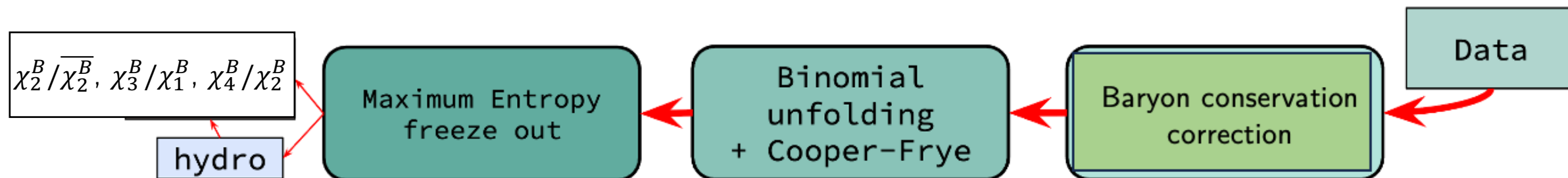
- 2D T'ExS Ising: critical point in μ_B plane ($\mu_Q = \mu_S = 0$) [Kahangirwe et al., PRD 109, 094046 (2024)]
- 4D T'ExS: non-critical lattice based EoS [A. Abuali, et al., PRD 112, 054502 (2025)]
- Switching energy density $e_{sw} = 0.18$ GeV/fm³ (to be varied, lower than chem. freeze-out)

Jahan, Roch, Shen, PRC 113, 024919 (2026)

Bayesian EoS inference

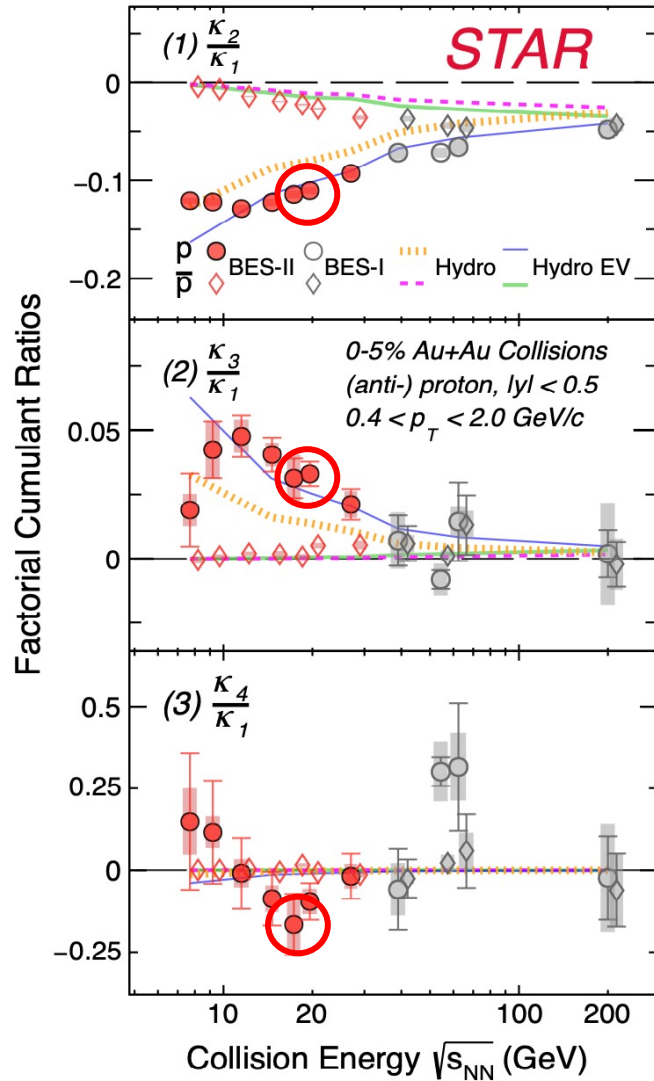
Instead of using χ_n^B as an input, extract χ_n^B 's that fit data directly at each energy

Specifically, the ratios $\chi_2^B / \overline{\chi_2^B}$, χ_3^B / χ_1^B , χ_4^B / χ_2^B where $\overline{\chi_2^B}$ is the HRG baseline

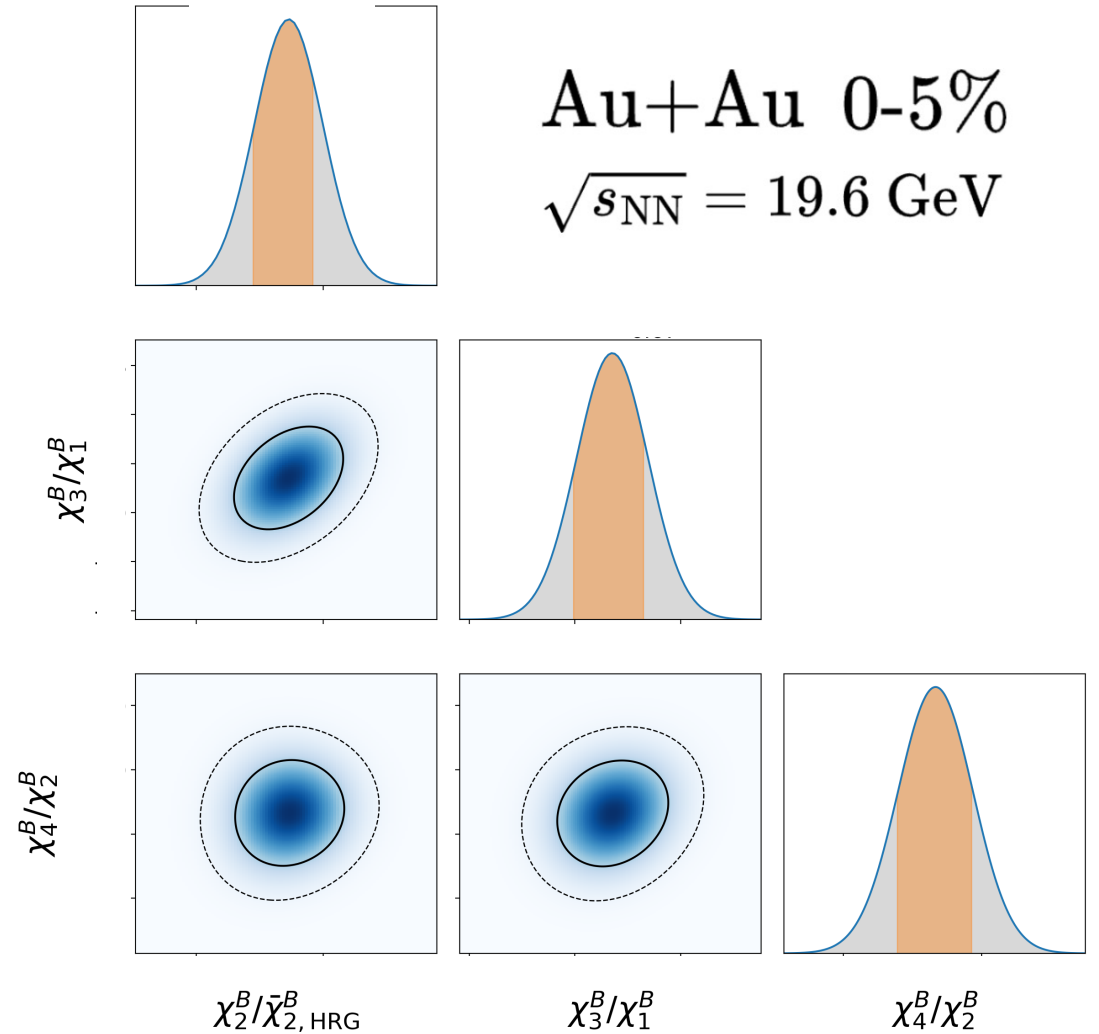


Adapted from G. Pihan, SQM2026

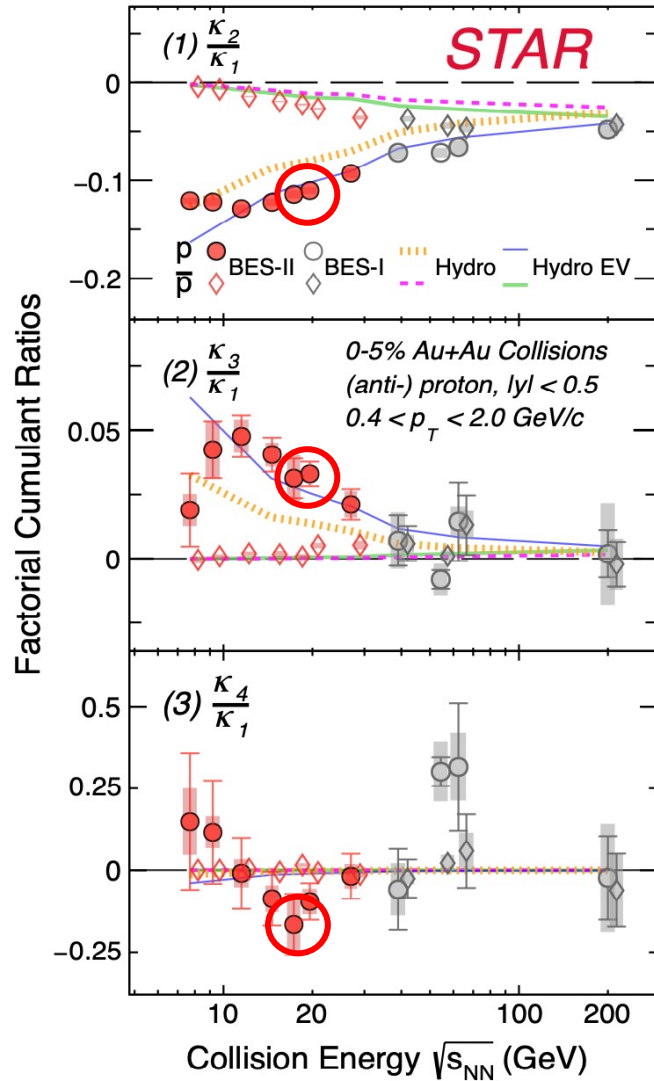
Bayesian inference: 19.6 GeV



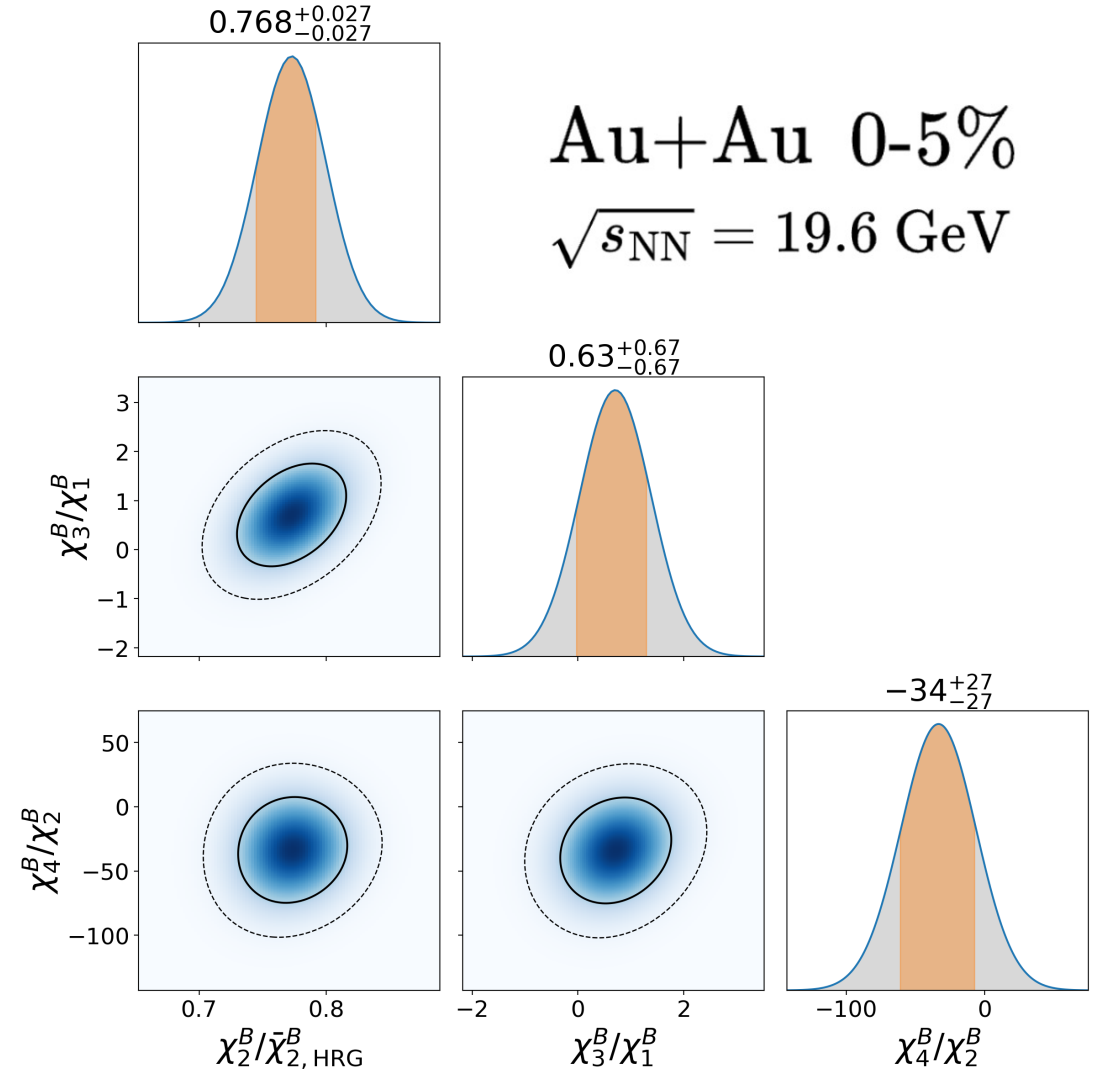
Unfolding



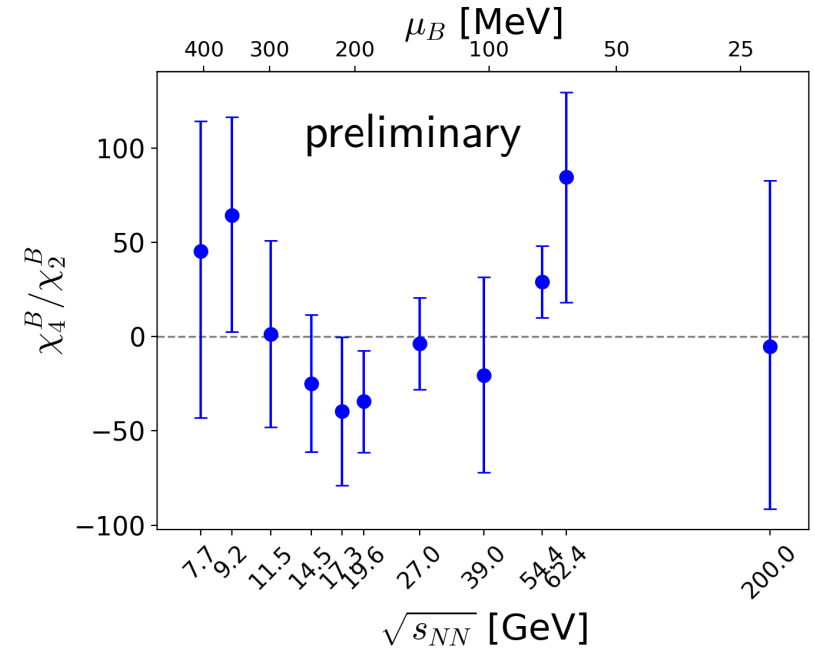
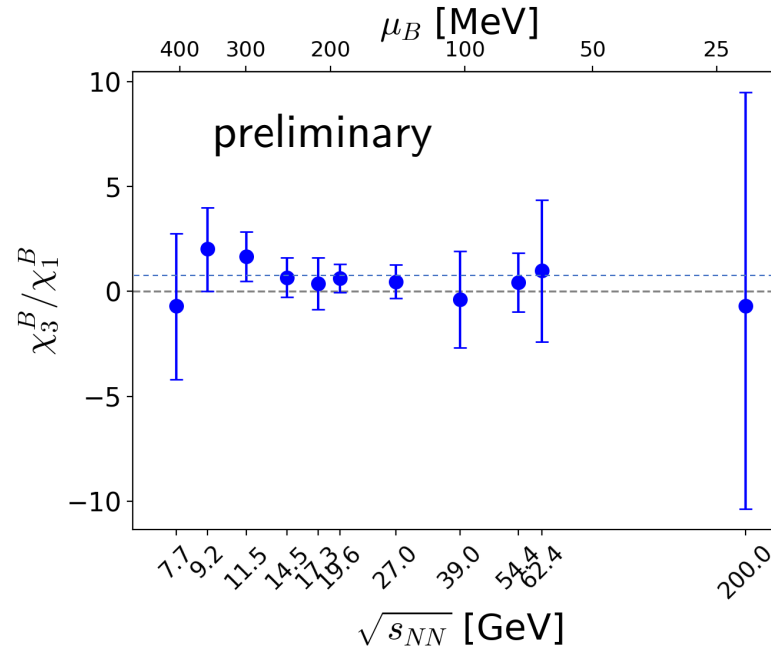
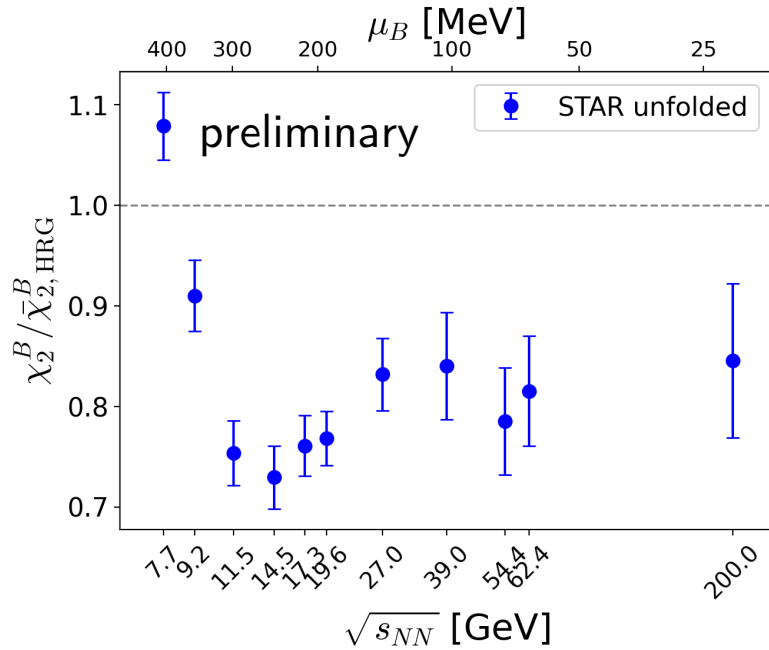
Bayesian inference: 19.6 GeV



Unfolding



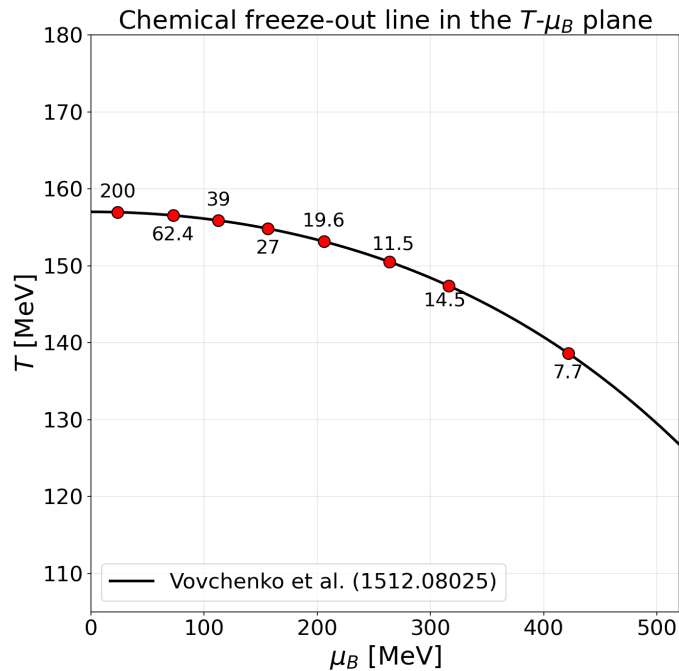
Extracted susceptibilities



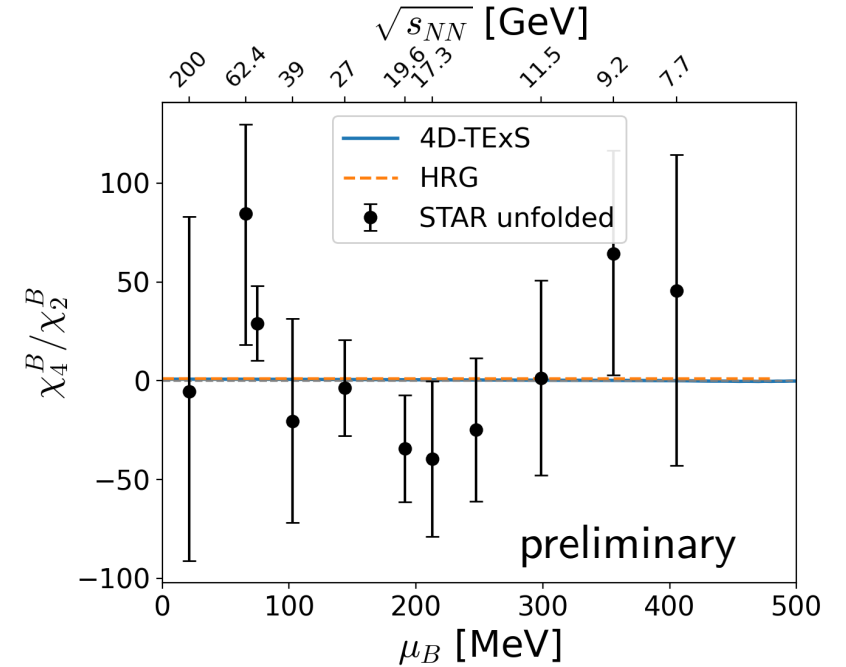
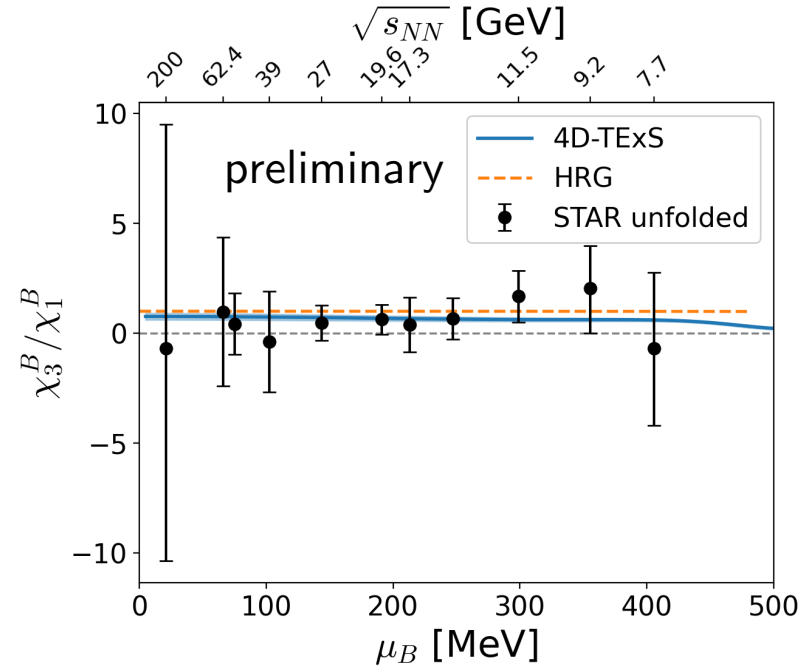
- Tight constraints on χ_2^B value relative to HRG
- Uncertainties in 3rd and 4th order susceptibilities are large---propagation of data uncertainty
 - Intermediate steps (net-B to B^+ , kinematic cuts, $B^+ \rightarrow p$) acts akin to efficiency correction
 - Reducing uncertainties in the data would lead to reduced extraction errors
- Results appear weakly sensitive to e_{sw} , sensitivity to other hydro parameters to be studied

Comparing to lattice QCD

To map the results to QCD phase diagram use chemical freeze-out parametrization [VV et al., PRC 93, 064906 (2015)]



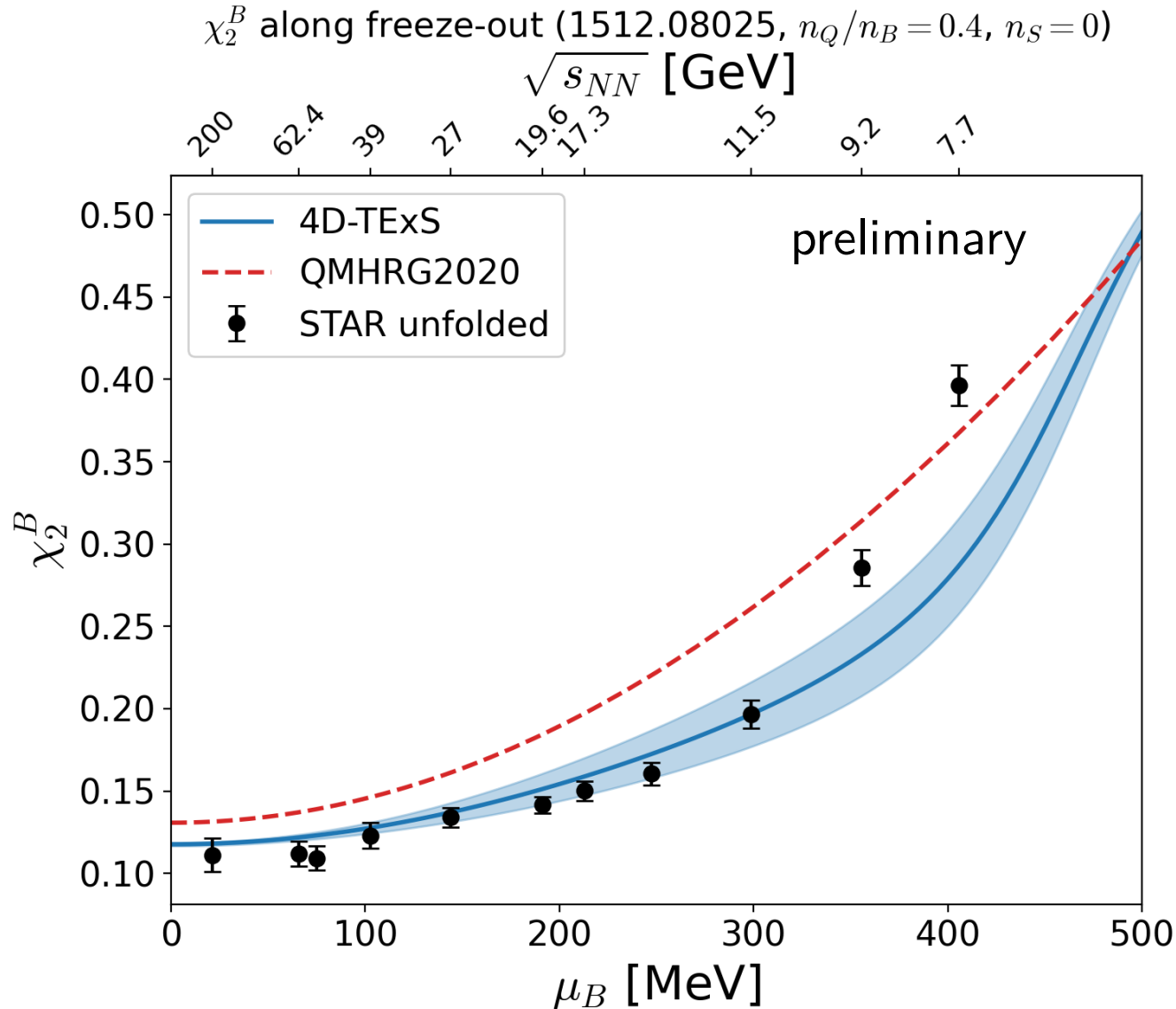
$$\varepsilon_{SW} \approx 0.35 - 0.40 \text{ GeV}/\text{fm}^3$$



4D-TExS: [A. Abuali et al., Phys. Rev. D 112, 054502 (2025)]

- Lattice-based EoS at finite baryon density
- Computed under HIC conditions ($n_Q/n_B = 0.4, n_S = 0$)

Baryon susceptibility χ_2^B



4D-TEoS: [A. Abuali et al., Phys. Rev. D 112, 054502 (2025)]

- To obtain bare χ_2^B multiply by HRG value
 - QMHRG motivated by lattice QCD studies at $\mu_B = 0$ [HotQCD, PRD 104, 074512 (2021)]
- Agreement with lattice QCD at $\mu_B < 300$ MeV
- Enhancement relative to lattice QCD at $\mu_B > 300$ MeV

⚠ Caveats:

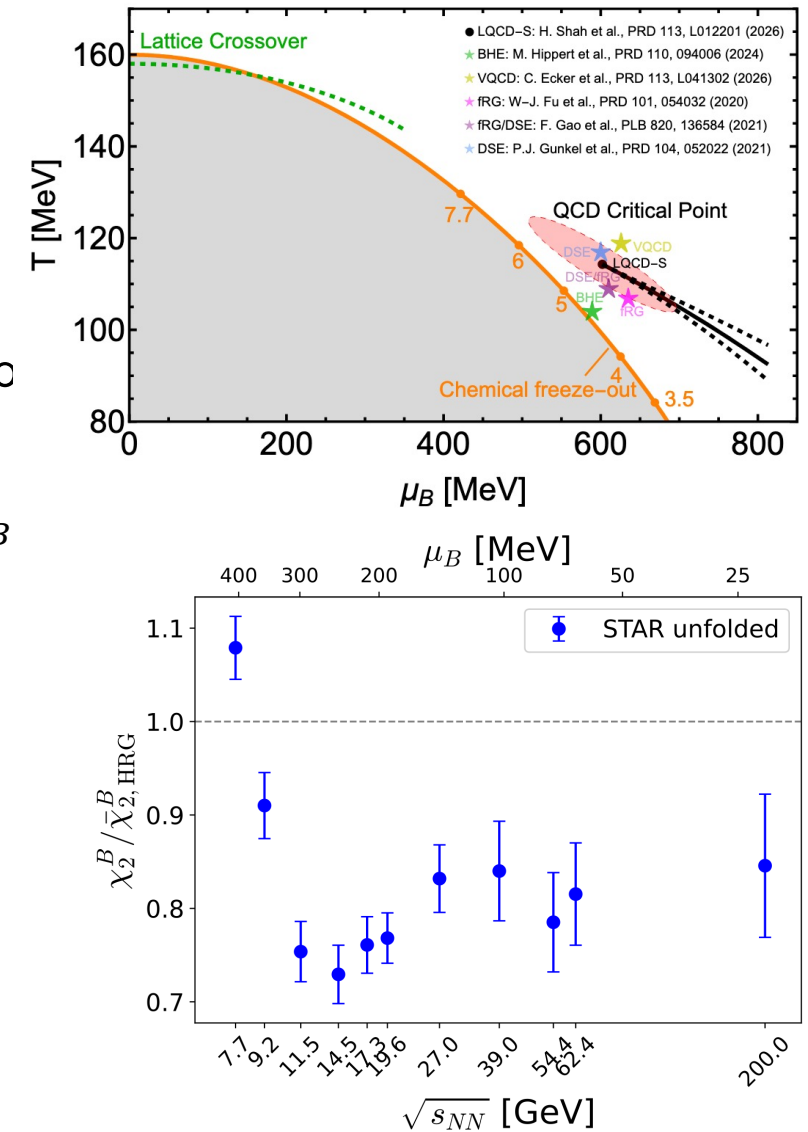
- Sensitive to the hadron list
- Sensitive to assumed freeze-out line
- Lattice EoS (4D-TEoS) is an extrapolation (no truncation error included)

Summary

- HRG is well approximated by a gas of confining strings
 - EV interactions may extend the description to $T \sim 200$ MeV
- Constant entropy contours: CP estimate in the 4D phase diagram
 - Strangeness neutrality shifts the CP from ~ 600 MeV to ~ 680 MeV
- Hydrodynamics-based inference with fluctuations in heavy-ion collisions
 - Non-critical physics describe the proton data at $\sqrt{s_{NN}} \geq 20$ GeV
 - χ_2^B consistent with LQCD at $\mu_B < 300$ MeV, enhancement at larger μ_B
 - Too large errors to infer high-order susceptibilities from data

Outlook:

- 4D EoS with a critical point from constant entropy contours
- Quantitative predictions for fluctuations in heavy-ion collisions
- Bayesian inference framework for HIC fluctuations data



Thanks for your attention