# QCD phase structure at finite baryon density from heavy-ion collisions

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# **Strongly interacting matter**

• Theory of strong interactions: *Quantum Chromodynamics* (QCD)

$$\mathcal{L} = \sum_{q=u,d,s,...} \bar{q} \left[ i \gamma^{\mu} (\partial_{\mu} - i g A^{a}_{\mu} \lambda_{a}) - m_{q} 
ight] q - rac{1}{4} G^{a}_{\mu
u} G^{\mu
u}_{a}$$



- Basic degrees of freedom: quarks and gluons that carry color charge
- At smaller energies confined into baryons (qqq) and mesons  $(q\bar{q})$

### Scales

- Length: 1 femtometer =  $10^{-15}$  m
- Temperature: 100 MeV $/k_B = 10^{12}$  K

### Where is it relevant?

- Early Universe
- Astrophysics: Neutron star (mergers)

Studied in laboratory with heavy-ion collisions





### **Non-perturbative methods**

**First-principle tool: Lattice QCD** 

### Ab-initio calculation of hadron masses





BMW Collaboration, Science 322, 1224 (2008)

Remarkable agreement of QCD with the experiment

### **QCD** transition from lattice **QCD**





- Analytic crossover at vanishing net baryon density at  $k_B T_{pc} \approx 155$  MeV a first-principle result [Y. Aoki et al., Nature 443, 675 (2006)]
- Finite densities inaccessible due to **sign problem**, but many effective theories predict first-order phase transition and the **QCD critical point**

### First-principle constraints on the QCD critical point

Indirect lattice QCD methods offer glimpse into small  $\mu_B/T$ 

• Taylor expansion around  $\mu_B/T=0$ 

$$\frac{p(T,\mu_B)}{T^4} = \frac{p(T,0)}{T^4} + \frac{\chi_2^B(T,0)}{2!}(\mu_B/T)^2 + \frac{\chi_4^B(T,0)}{4!}(\mu_B/T)^4 + \dots$$

No hints for the critical point at T > 135 MeV Critical point  $\mu_B/T < 3$  disfavored

- Relativistic virial expansion in fugacities via analytic continuation from imaginary  $\mu_B/T$ 

$$\frac{p(T,\mu_B)}{T^4} = \sum_{k=0}^{\infty} p_k(T) \cosh\left(\frac{k\,\mu_B}{T}\right)$$

Expansion sees singularity in the complex plane, Im  $[\mu_B/T] = \pi$ Critical point at  $\mu_B/T < \pi$  disfavored

#### Critical point, if it exists, likely located beyond the reach of lattice methods



[HotQCD Collaboration, PRD 95, 054504 (2017)]



[V.V., Steinheimer, Philipsen, Stoecker, PRD 97, 114030 (2018)]

### Relativistic heavy-ion collisions – "Little Bangs"



**Control** parameters

- Collision energy  $\sqrt{s_{NN}} = 2.4 5020 \text{ GeV}$
- Size of the collision region

### Measurements

• Final hadron abundances and momentum distributions

### **QCD** phase diagram with heavy-ion collisions



STAR event display

Thousands of particles created in relativistic heavy-ion collisions

Apply concepts of statistical mechanics

### Particle production in heavy-ion collisions

Ideal gas law (E. Clapeyron, 1834)  $P_i V = N_i k_B T$ 

is the simplest model of particle production



$$N_i = \frac{d_i V}{2\pi^2} \int dk \ k^2 \left[ 1 \pm \exp\left(\frac{\sqrt{k^2 + m_i^2} - \mu_i}{T}\right) \right]^{-1}$$

Bose-Einstein & Fermi-Dirac, 1924-1926



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ALICE collaboration (Quark Matter 2018)

# Hadron resonance gas (HRG) model

• **HRG model:** free gas of known hadrons and resonances



 $p(T, \mu_B) = T \phi_M(T) + 2 T \phi_B(T) \cosh(\mu_B/T)$ mesons $\phi_{M(B)}(T) = \sum_{i \in M(B)} \frac{d_i}{2\pi^2} \int dk \, k^2 \exp\left(-\frac{\sqrt{m_i^2 + k^2}}{T}\right)$ 

- Hadronic interactions dominated by resonance formation\*
- Leading order in relativistic virial expansion
- Matches well with lattice QCD below  $T_{pc}$
- Non-resonant interactions incorporated in extended descriptions

#### HRG model and heavy-ion collisions:

Basis for the thermal model of particle production



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All bells and whistles implemented in open source codes, e.g. Thermal-FIST [VV, Stoecker, Comput. Phys. Commun. 244, 295 (2019)]

\* Dashen, Ma, Bernstein, "S-matrix formulation of statistical mechanics", Phys. Rev. (1969); Prakash, Venugopalan, Nucl. Phys. A (1992)

# Mapping heavy-ion collisions onto the QCD phase diagram



 $\sqrt{S_{NN}}$ 

 $\mu_B \nearrow$ 

For  $p_T$  differential observables (spectra, flow, ...) use relativistic hydrodynamics

## **QCD** critical point



Figure from Bzdak et al., Phys. Rept. '20

What is the nature of the quark-hadron transition at finite baryon density?

Is there a QCD phase transition and critical point? Where?

Tackle these questions with heavy-ion collisions

### **Critical point and fluctuations**

Density fluctuations at macroscopic length scales

Critical opalescence



Unfortunately, we cannot do this in heavy-ion collisions

### **Event-by-event fluctuations and statistical mechanics**

Consider a fluctuating number N

Cumulants:  $G_N(t) = \ln \langle e^{tN} \rangle = \sum_{n=1}^{\infty} \kappa_n \frac{t^n}{n!}$ variance  $\kappa_2 = \langle (\Delta N)^2 \rangle = \sigma^2$ 

skewness

 $\kappa_3 = \langle (\Delta N)^3 
angle$ 

kurtosis

 $\kappa_4 = \langle (\Delta N)^4 \rangle - 3 \langle (\Delta N^2) \rangle^2$ 

width
asymmetry

peak shape





### Statistical mechanics:

Grand partition function

$$ln Z^{
m gce}(T,V,\mu) = ln \left[\sum_{N} e^{\mu N} Z^{
m ce}(T,V,N)
ight],$$

$$\kappa_n \propto rac{\partial^n (\ln Z^{
m gce})}{\partial (\mu_N)^n}$$

### Cumulants measure chemical potential derivatives of the (QCD) equation of state

### **Applications**

• (QCD) critical point – large critical fluctuations of baryon (proton) number



M. Stephanov, Phys. Rev. Lett. (2011)

• Test of (lattice) QCD at  $\mu_B \approx 0$ 



Correlation length  $\xi \to \infty$  diverges at the critical point

$$\kappa_2\sim\xi^2$$
,  $\kappa_3\sim\xi^{4.5}$ ,  $\kappa_4\sim\xi^7$ 

Looking for non-monotonic dependence of  $\kappa_4$  vs  $\sqrt{s_{NN}}$ 

• Freeze-out from fluctuations



Borsanyi et al. PRL 113, 052301 (2014); Bazavov et al. PRL 109, 192302 (2012) 14

### Example: Liquid-gas transition with van der Waals equation



VV, Anchishkin, Gorenstein, Poberezhnyuk, Phys. Rev. C 92, 054901 (2015)

### Measuring cumulants in heavy-ion collisions



Cumulants are extensive,  $\kappa_n \sim V$ , use ratios to cancel out the volume

$$\frac{\kappa_2}{\langle N \rangle}$$
,  $\frac{\kappa_3}{\kappa_2}$ ,  $\frac{\kappa_4}{\kappa_2}$ 

### **Experimental measurements**



Reduced errors (better statistics) to come soon from beam energy scan II program

Can we learn more from the more accurate data available for  $\kappa_2$  and  $\kappa_3$ ?

### **Theory vs experiment: Challenges for fluctuations**

### Theory



 $\ensuremath{\mathbb{C}}$  Lattice QCD@BNL

- Coordinate space
- In contact with the heat bath
- Conserved charges
- Uniform
- Fixed volume

### **Experiment**



STAR event display

- Momentum space
- Expanding in vacuum
- Non-conserved particle numbers
- Inhomogenous
- Fluctuating volume

# When are the measured fluctuations grand-canonical?

- Consider event-by-event fluctuations of particle number in acceptance  $\Delta Y_{accept}$  around midrapidity
- Scales
  - $\Delta Y_{accept}$  acceptance
  - $\Delta Y_{total}$  full space
  - $\Delta Y_{corr}$  rapidity correlation length (thermal smearing)
  - $\Delta Y_{kick}$  diffusion in the hadronic phase
- **GCE** applies if  $\Delta Y_{total} \gg \Delta Y_{accept} \gg \Delta Y_{kick}, \Delta Y_{corr}$
- In practice  $\Delta Y_{total} \gg \Delta Y_{accept}$  and  $\Delta Y_{accept} \gg \Delta Y_{corr}$  are not simultaneously satisfied
  - Corrections from baryon conservation are large [Bzdak et al., PRC '13]
  - $\Delta Y_{corr} \sim 1 \sim \Delta Y_{accept}$  [Ling, Stephanov, PRC '16]

### Need dynamical description



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V. Koch, arXiv:0810.2520

### Dynamical approaches to the QCD critical point search

- 1. Dynamical model calculations of critical fluctuations
  - Fluctuating hydrodynamics
  - Equation of state with tunable critical point [P. Parotto et al, Phys. Rev. C 101, 034901 (2020)]
  - Predict CP signatures dependent on its location

Under development within the Beam Energy Scan Theory (BEST) Collaboration

**BEST** [X. An et al., Nucl. Phys. A 1017, 122343 (2022)]

- 2. Deviations from precision calculations of the non-critical baseline
  - Include essential non-critical contributions to (net-)proton number cumulants
  - Exact baryon conservation + hadronic interactions (hard core repulsion)
  - Based on realistic hydrodynamic simulations tuned to bulk data

[VV, C. Shen, V. Koch, Phys. Rev. C 105, 014904 (2022)]

### **Excluded volume effect**

Incorporate repulsive baryon (nucleon) hard core via excluded volume VV, M.I. Gorenstein, H. Stoecker, Phys. Rev. Lett. 118, 182301 (2017)

Amounts to a van der Waals correction for baryons in the HRG model

 $V \rightarrow V - bN$ 



 $\leftarrow 2r \rightarrow$ 





Figure from Ishii et al., PRL '07

• Net baryon kurtosis suppressed as in lattice QCD

$$\frac{\chi_4^B}{\chi_2^B} \simeq 1 - \frac{12b\phi_B(T)}{\Phi_B(T)} + O(b^2)$$

• Reproduces virial coefficients of baryon interaction from lattice QCD

Excluded volume from lattice QCD: b

$$b \approx 1 \text{ fm}^3$$



VV, A. Pasztor, S. Katz, Z. Fodor, H. Stoecker, Phys. Lett. B 755, 71 (2017) 21

### **RHIC-BES:** Hydrodynamic description in non-critical scenario

- Collision geometry based 3D initial state
  - Constrained to net proton distributions [Shen, Alzhrani, Phys. Rev. C '20]
- Viscous hydrodynamics evolution MUSIC-3.0
  - Energy-momentum and baryon number conservation
  - Crossover equation of state based on lattice QCD [Monnai, Schenke, Shen, Phys. Rev. C '19; also Noronha-Hostler, Parotto, Ratti, Stafford, Phys. Rev. C '19]
- Cooper-Frye particlization at  $\epsilon_{sw} = 0.26 \text{ GeV}/\text{fm}^3$

$$\omega_p rac{dN_j}{d^3 p} = \int_{\sigma(x)} d\sigma_\mu(x) \, p^\mu \, rac{d_j \, \lambda_j^{\mathsf{ev}}(x)}{(2\pi)^3} \, \exp\left[rac{\mu_j(x) - u^\mu(x) p_\mu}{T(x)}
ight].$$

- Particlization respects QCD-based baryon number distribution
  - Incorporated via baryon excluded volume b = 1 fm<sup>3</sup>
     [VV, V. Koch, Phys. Rev. C 103, 044903 (2021)]
- Incorporates exact global baryon conservation via a novel method [VV, Phys. Rev. C 105, 014903 (2022)]







### Net proton cumulant ratios



- Data at  $\sqrt{s_{NN}} \ge 20$  GeV consistent with non-critical physics (baryon conservation and repulsion)
- Effect from baryon conservation is larger than from repulsion
- Excess of skewness in data at  $\sqrt{s_{NN}} < 20$  GeV hint of attractive interactions? Critical point?

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### **Correlation Functions**

• Analyze genuine multi-particle correlations via factorial cumulants  $\hat{C}_n$ [Bzdak, Koch, Strodthoff, Phys. Rev. C '17]

$$\hat{C}_1 = \kappa_1, \qquad \hat{C}_3 = 2\kappa_1 - 3\kappa_2 + \kappa_3, \\ \hat{C}_2 = -\kappa_1 + \kappa_2, \quad \hat{C}_4 = -6\kappa_1 + 11\kappa_2 - 6\kappa_3 + \kappa_4$$

- Three- and four-particle correlations are small without a CP
  - Multi-particle correlations expected near the critical point [Ling, Stephanov, PRC '15]

- Signals from the data at  $\sqrt{s_{NN}} \le 20$  GeV
  - Excess of two-proton correlations
  - Possibility of significant 4-proton correlations
  - Critical point?





# Lower energies $\sqrt{s_{NN}} \le 7.7$ GeV



Naïvely, could indicate QCD critical point near HADES freeze-out at  $T \approx 70$  MeV,  $\mu_B \approx 875$  MeV

Some effective QCD approaches do predict the critical point close to that region, e.g. holography [Critelli, Noronha, Noronha-Hostler, Portillo, Ratti, Rougemont, Phys. Rev. D 96, 096026 (2017)] [Grefa et al., Phys. Rev. D 104, 034002 (2021)]

### **QCD** phase structure: What we learned so far



- Data at high energies ( $\sqrt{s_{NN}} \ge 20$  GeV) consistent with "non-critical" physics
- Disfavors critical point at  $\mu_B/T < 2-3$ , consistent with what we know from lattice QCD
- Interesting physics at high densities probed by future experiments, neutron stars & their mergers

### **Outlook: Equation of state for heavy ions and neutron stars**



P. Senger (GSI)



L. Weih, L. Rezzolla (Frankfurt)

# **Summary**

- Strongly interacting QCD matter under extreme conditions
  - Undergoes a transition to quark-gluon plasma at trillion kelvin degrees
  - Behaves like a fluid described by hydrodynamics
  - Phase structure at finite baryon density still largely unknown, no phase transition at high energies
- Fluctuations are a powerful tool to explore the QCD phase diagram
  - Heavy-ion data are described quantitatively at  $\sqrt{s_{NN}} \ge 20$  GeV ( $\mu_B/T < 3$ ) without critical point
  - Possible critical point signals at  $\sqrt{s_{NN}} < 14.5$  GeV
  - More evidence at lower energies to come from future experiments and connections to neutron star phenomena

### Thanks for your attention!

# Backup slides

### **Acceptance dependence of two-particle correlations**

- Changing  $y_{max}$  slope at  $\sqrt{s_{NN}} \le 14.5$  GeV? ٠
- Volume fluctuations? [Skokov, Friman, Redlich, PRC '13] •
  - $C_2/C_1 + = C_1 * v_2$ •
  - Can improve low energies but spoil high energies? •
- Exact electric charge conservation? •
  - Worsens the agreement at  $\sqrt{s_{NN}} \le 14.5$ , higher energies ٠ virtually unaffected

⊢ <sub>15</sub>/

das

μ (MeV)

- **Attractive interactions?** 
  - Could work if baryon repulsion turns • into attraction in the high- $\mu_B$  regime (MeV) 20
  - **Critical point?** ٠



### **Outlook: baryon cumulants from protons**

- net baryon  $\neq$  net proton
- Baryon cumulants can be reconstructed from proton cumulants via binomial (un)folding based on isospin randomization [Kitazawa, Asakawa, Phys. Rev. C 85 (2012) 021901]
  - Amounts to an additional "efficiency correction" and requires the use of joint factorial moments, only experiment can do it model-independently









- Net protons described within errors but not sensitive to the equation of state for the present experimental acceptance
- Large effect from resonance decays for lighter particles + conservation of electric charge/strangeness
- Future measurements will require larger acceptance







### Effects of baryon annihilation and local conservation

O. Savchuk, V.V., V. Koch, J. Steinheimer, H. Stoecker, arXiv:2106.08239

Baryon annihilation  $B\overline{B} \rightarrow n\pi$  in afterburners (UrQMD, SMASH) suppresses baryon yields



- ALICE data requires local baryon conservation across  $\Delta y \sim \pm 1.5$  with UrQMD annihilations (no regenerations) or global conservation ( $\Delta y \sim \Delta y_{tot}$ ) without annihilations
- Local conservation and  $B\overline{B}$  annihilation can be constrained from data through the combined analysis of  $\kappa_2[p-\overline{p}]$  and  $\kappa_2[p+\overline{p}]$

# Thermodynamic analysis of HADES data

#### VV, Koch, in preparation

- Single freeze-out scenario: Emission from Siemens-Rasmussen hypersurface with Hubblelike flow
  - $\rightarrow$  Pion and proton spectra o.k. [S. Harabasz et al., PRC 102, 054903 (2020)]
- Uniform  $T \approx 70$  MeV,  $\mu_B \approx 875$  MeV across the fireball [A. Motornenko et al., PLB 822, 136703 (2021)]

### • Fluctuations:

- Same as before but incorporate additional binomial filtering to account for protons bound in light nuclei
- Uniform fireball  $\rightarrow$  Final proton cumulants are linear combinations of baryon susceptibilities  $\chi^B_n$





• In the grand-canonical limit (no baryon conservation) the data are described well with

$$\frac{\chi_2^B}{\chi_1^B} = 9.35 \pm 0.40, \qquad \frac{\chi_3^B}{\chi_2^B} = -39.6 \pm 7.2, \qquad \frac{\chi_4^B}{\chi_2^B} = 1130 \pm 488$$

- Could be indicative of a critical point near the HADES freeze-out at  $T \approx 70$  MeV,  $\mu_B \approx 875$  MeV
- However, the results are challenging to describe with baryon conservation included

